

Progress in the Nucleon EDM Calculations in Lattice QCD

Sergey N. Syritsyn,
*Stony Brook University & RIKEN / BNL Research Center
together with LHP and RBC collaborations*



Stony Brook
University

BROOKHAVEN
NATIONAL LABORATORY



HOKUSAI
GREAT WAVE

Argonne
NATIONAL LABORATORY

*13th Conference on the Intersections of
Particle and Nuclear Physics
Palm Springs, CA, May 29–Jun 3, 2018*

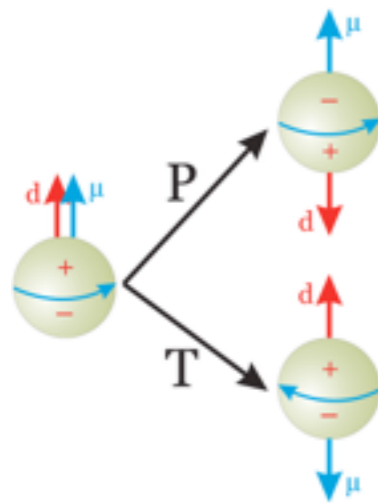
Outline

- θ_{QCD} -induced nucleon EDM:
 - *Previous calculations*
 - *New challenges*
 - *Improved techniques on a lattice*
 - *Outlook for θ_{QCD} -nEDM*
- Quark chromo-EDM-induced nucleon EDM
 - *Preliminary results at the physical point*

CP Violation: Electric Dipole Moments

EDMs are the most sensitive probes of CPv:

- Prerequisite for Baryogenesis
- Evidence for SM Extensions
- (θ_{QCD} in particular) Strong CP problem



$$\vec{d}_N = d_N \frac{\vec{S}}{S} \quad \mathcal{H} = -\vec{d}_N \cdot \vec{E}$$

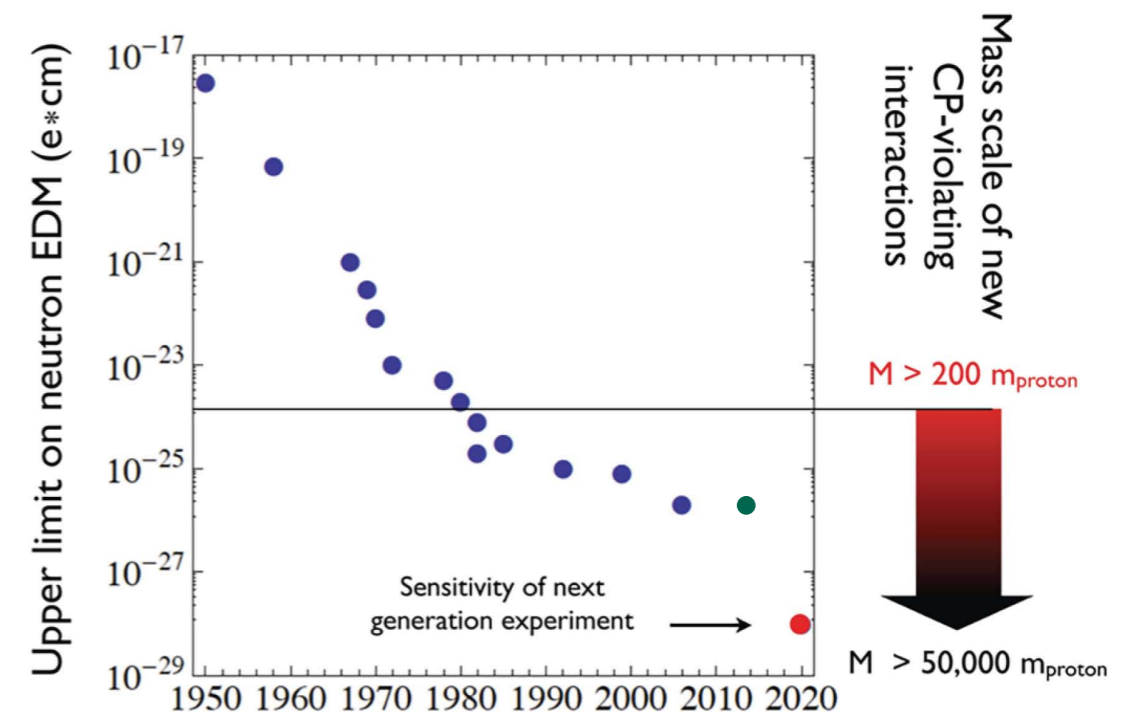
$$\text{OR} \quad \mathcal{L}_{int} = eA_{\mu}^{\text{em}} \mathcal{V}^{\mu} \quad (\text{P,T-even})$$

$$+ eA_{\mu}^{\text{em}} \mathcal{A}^{\mu} \quad (\text{P,T-odd})$$

Experimental Outlook: Neutron EDM

	$10^{-28} e \text{ cm}$
CURRENT LIMIT	<300
Spallation Source @ORNL	< 5
Ultracold Neutrons @LANL	~30
PSI EDM	<50 (I), <5 (II)
ILL PNPI	<10
Munich FRMII	< 5
RCMP TRIUMF	<50 (I), <5 (II)
JPARC	< 5
Standard Model (CKM)	< 0.001

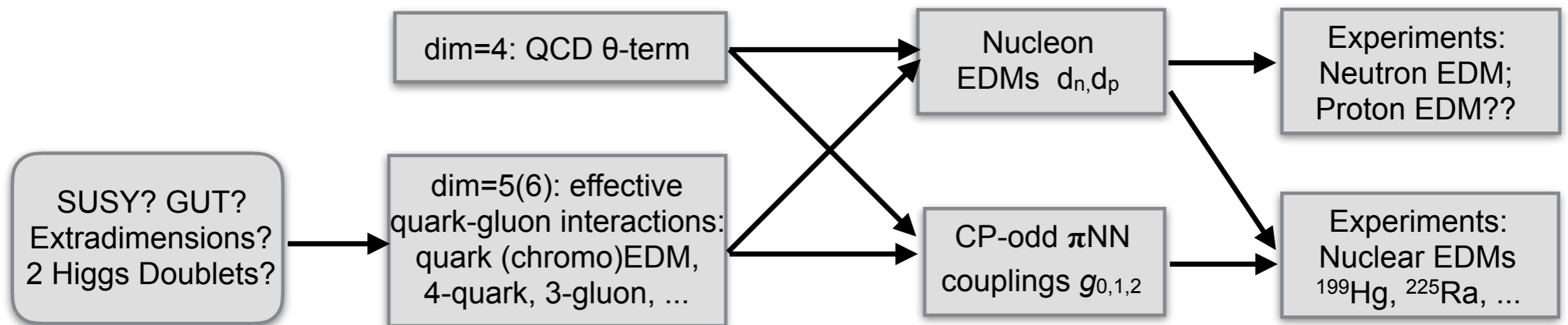
[B.Filippone's talk, KITP 2016]



nEDM sensitivity :

- 1–2 years : next best limit
- 3–4 years : x10 improvement
- 7-10 years : x100 improvement

Nucleon EDMs: a Window into New Physics



- Effective quark-gluon CPv interactions organized by dimension

$$\mathcal{L}_{eff} = \sum_i \frac{c_i}{[\Lambda_{(i)}]^{d_i-4}} \mathcal{O}_i^{[d_i]} \quad [\text{J.Engel, M. Ramsey-Musolf, U. van Kolck, Prog.Part.Nucl.Phys. 71 (2013), pp. 21-74}]$$

$$d=4 : \theta_{QCD}$$

$$d=5(6) : \text{quark EDM, quark-gluon chromo EDM}$$

$$d=6 : \text{4-fermion CPv, 3-gluon (Weinberg)}$$

$$d_{n,p}$$

$$F_3^{n,p}(Q^2)$$

- $d_{n,p} = d_{n,p}^\theta \theta_{QCD} + d_{n,p}^{cEDM} c_{cEDM} + \dots$

$$c_i \iff d_{n,p} ?$$

lattice QCD calculations are needed to relate to constrain $\theta_{QCD}, c_{cEDM}, \dots$

CP-odd Nucleon Structure on a Lattice

CP-broken vacuum on a lattice:

- Linear response to CP-odd interaction (e.g., QCD θ -term)

$$\langle \mathcal{O} \dots \rangle_{CP} = \langle \mathcal{O} \dots \rangle_{CP-even} - i\theta \langle Q \cdot \mathcal{O} \dots \rangle_{CP-even} + O(\theta^2)$$

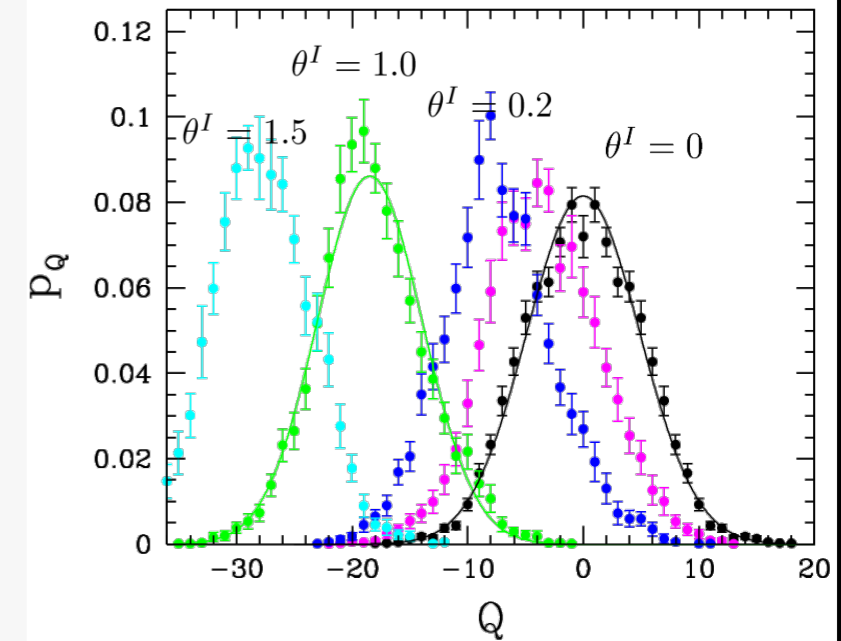
[S. Aoki et al (2005); F. Berruto et al (2005); A.Shindler et al (2015) ;
C. Alexandrou et al (2015) ; E. Shintani et al (2016)]

- Simulation with dynamical (imaginary) θ'_{QCD}

$$\langle \mathcal{O} \dots \rangle_{\theta} \sim \int \mathcal{D}U e^{-S - \theta' Q} (\mathcal{O} \dots)$$

[R.Horsley et al (2008) ; F.K.Guo et al (2015)]

new gauge ensembles \Rightarrow better sampling of $Q \neq 0$ sectors



Extraction of d_N

- Nucleon spectrum in the bg. electric field

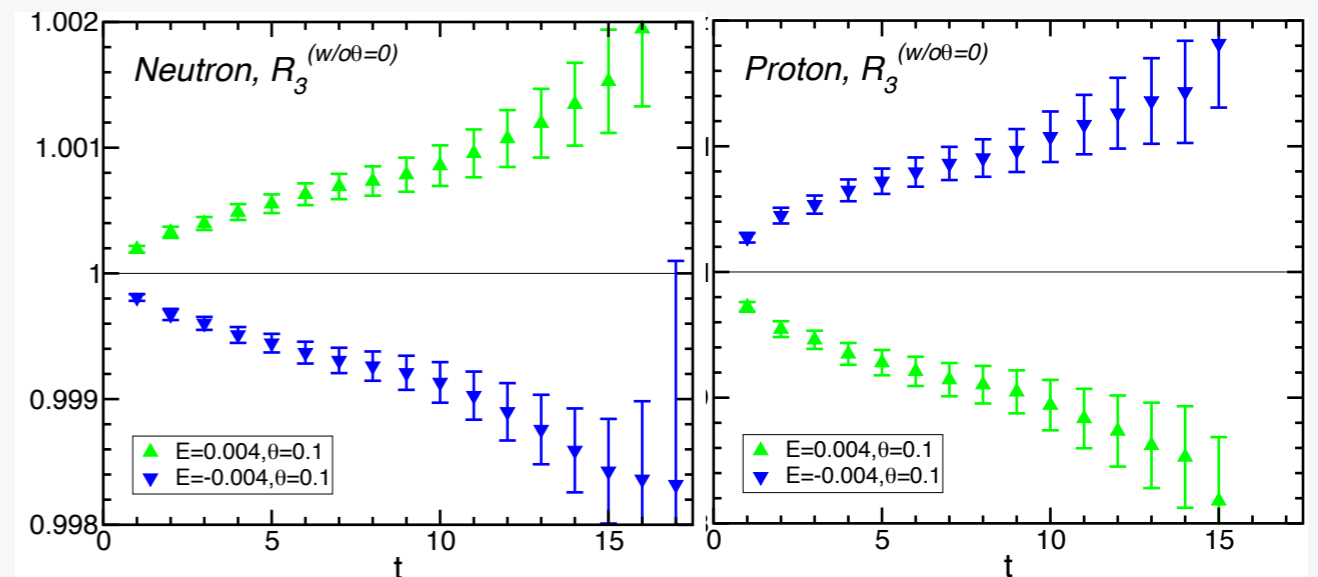
[S.Aoki et al '89 ; E.Shintani et al '06;
E.Shintani et al, PRD75, 034507(2007)]

$$\langle N(t) \bar{N}(0) \rangle_{\theta, \vec{E}} \sim e^{-(E \pm \vec{d}_N \cdot \vec{E})t}$$

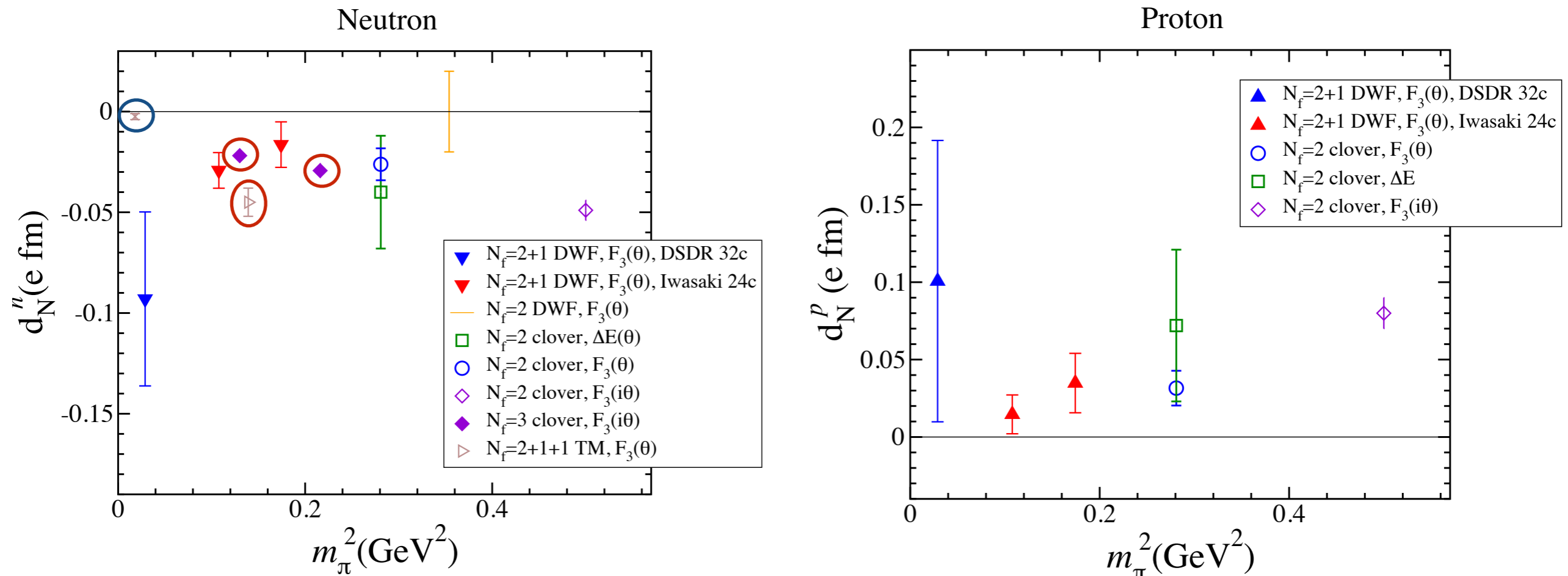
- P,T-odd Form Factor $d_N = F_3(0)/2m_N$

[E.Shintani et al '05, '15 ; F.Berruto et al '05 ;
A.Shindler et al '15 ; C.Alexandrou et al '15]

Require extrapolation $F_3(Q^2 \rightarrow 0)$



θ_{QCD} -induced Nucleon EDM



Summary of nEDM from LQCD circa 2015

[E.Shintani, T.Blum, T.Izubuchi, A.Soni, PRD93, 094503(2015)]

- Phenomenology: $|d_n| \approx \theta_{QCD} \times (0.4 \dots 2.5) \cdot 10^{-3} \text{ e fm}$
 - Lattice [Guo et al 2015]: $|d_n| \approx \theta_{QCD} \times (4 \cdot 10^{-3} \text{ e fm})$
- \implies tighter constraint on θ_{QCD} ?

Unfortunately, there was a problem...

Nucleon "Parity Mixing"

CPv interaction induces a chiral phase in fermion fields:

$$\begin{aligned}
 \langle \text{vac} | N | p, \sigma \rangle_{\mathcal{CP}} &= e^{i\alpha\gamma_5} u_{p,\sigma} = \tilde{u}_{p,\sigma} \\
 \downarrow & \qquad \qquad \qquad \downarrow \\
 u [u^T C \gamma_5 d] & \qquad \qquad \qquad (\not{\partial} + m_N e^{-2i\alpha\gamma_5}) \tilde{u}_p = 0 \\
 \text{(P-even lattice nucleon field)} & \qquad \qquad \qquad \sum_{\sigma} \tilde{u}_{p,\sigma} \bar{\tilde{u}}_{p,\sigma} \sim (-i\not{p}_{\mathcal{E}} + m_N e^{2i\alpha\gamma_5})
 \end{aligned}$$

Vector current M.E. has to be defined with positive-parity spinors to define $F_{2,3}$
 [SNS, S.Aoki, et al (2017) arXiv:1701.07792]

$$\langle N_{p'} | \bar{q} \gamma^\mu q | N_p \rangle_{\mathcal{CP}} = \bar{u}_{p'} \left[F_1 \gamma^\mu + \underbrace{(F_2 + iF_3 \gamma_5)}_{\Gamma_{\mathcal{E}}^\mu} \frac{i\sigma^{\mu\nu} (p' - p)_\nu}{2m_N} \right] u_p \qquad \begin{aligned} \gamma_4 u &= +u \\ \bar{u} \gamma_4 &= +\bar{u} \end{aligned}$$

... otherwise, $F_{2,3}$ mix under chiral rotation and lead to fake EDM/EDFF signal

$$\begin{aligned}
 e^{i\alpha\gamma_5} \Gamma^\mu e^{i\alpha\gamma_5} &\leftrightarrow \Gamma^\mu & \text{"}F_3\text{"} &\approx [F_3]_{\text{true}} - 2\alpha [F_2]_{\text{true}} \\
 e^{2i\alpha} (\text{"}F_2\text{"} + i\text{"}F_3\text{"}) &= (F_2 + iF_3)_{\text{true}} & \text{"}d_{n,p}\text{"} &\approx [d_{n,p}]_{\text{true}} - 2\alpha \frac{\kappa_{n,p}}{2m_N}
 \end{aligned}$$

The same issue is addressed correctly in EFT (ChPT) calculations

Nucleon "Parity Mixing" (2)

With proper definition of $F_{2,3}$ [SNS, S.Aoki, *et al* (2017) arXiv:1701.07792]

- coupling of E,B to spin in the forward limit

$$\langle H_{\text{int}} \rangle = eA_\mu \langle J^\mu \rangle = -\frac{eG_M(0)}{2m_N} \vec{\Sigma} \cdot \vec{H} - \frac{eF_3(0)}{2m_N} \vec{\Sigma} \cdot \vec{E}$$

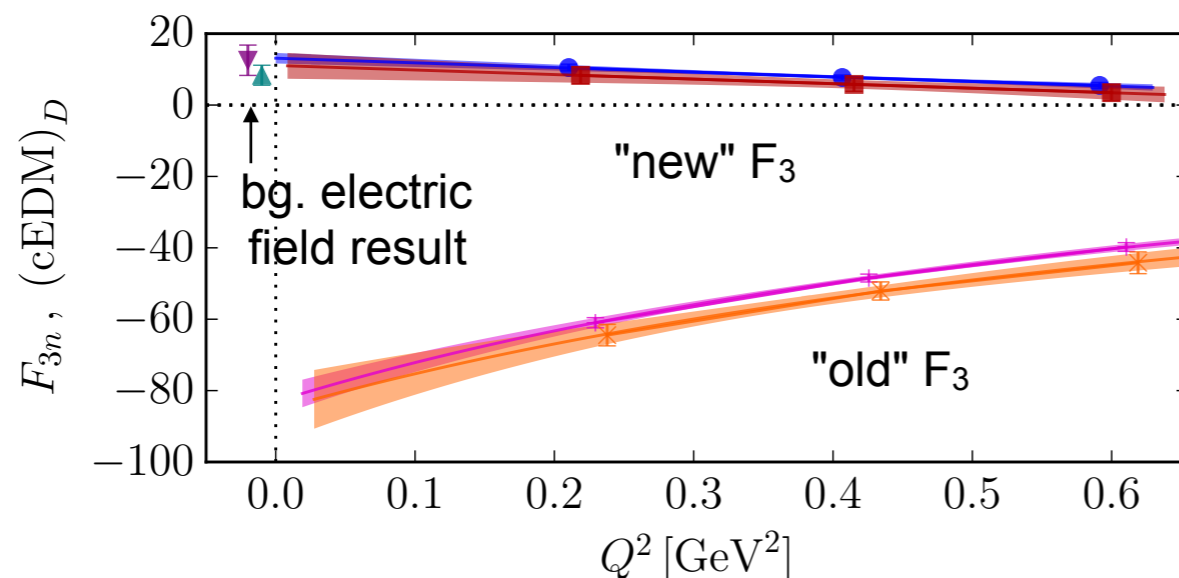
- poles of the Dirac operator in bg. electric & magnetic fields

$$\mathcal{L}_N = \bar{N} \left[i\not{\partial} - me^{-2i\alpha\gamma_5} - Q\gamma_\mu A^\mu - (\tilde{\kappa} + i\tilde{\zeta}\gamma_5) \frac{1}{2} F_{\mu\nu} \frac{\sigma^{\mu\nu}}{2m_N} \right] N$$

$$\curvearrowright E_N(\vec{p}=0) - m_N = -\frac{\kappa}{2m_N} \vec{\Sigma} \cdot \vec{H} - \frac{\zeta}{2m_N} \vec{\Sigma} \cdot \vec{E} + O(\kappa^2, \zeta^2)$$

with $\kappa + i\zeta = e^{2i\alpha\gamma_5} (\tilde{\kappa} + i\tilde{\zeta})$

- Numerical test: compare EDFF with mass shift in uniform bg. electric field



d-cEDM induces large mixing
 $\alpha_D \approx 30(0.2)$

Large F_{2n} contribution to " F_{3n} "
 $"F_{3n}^D" = [F_{3n}^D]_{\text{true}} - 2\alpha_D F_{2n}$

Recent Lattice Results on θ_{QCD} -induced nEDM

Correction to previous results:

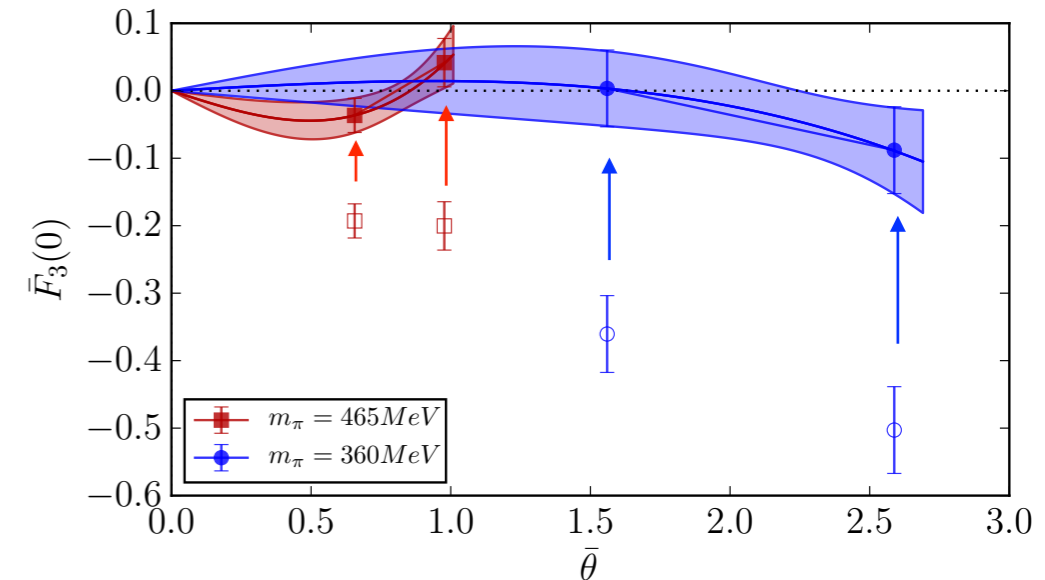
$$[F_3]_{\text{true}} = "F_3" + 2\alpha F_2$$

- [F. Guo *et al* (QCDSF), PRL115:062001 (2015)]
dynamical calculations with finite imag. θ' angle

- [C.Alexandrou *et al* (ETMC), PRD93:074503 (2016)]
 $d_n = -0.045(06) \text{ e fm } (\sim 7.5\sigma) \rightarrow +0.008(6) \text{ e fm } (1.3\sigma)$

- *Uniform bg. electric field method is not affected by "parity mixing"*

*Precision in Ref. [E.Shintani *et al*, D78:014503 (2008)] is insufficient for comparison*



	m_π [MeV]	m_N [GeV]	F_2	α	\tilde{F}_3	F_3	
[ETMC 2016]	n	373	1.216(4)	$-1.50(16)^a$	$-0.217(18)$	$-0.555(74)$	0.094(74)
[Shintani <i>et al</i> 2005]	n	530	1.334(8)	$-0.560(40)$	$-0.247(17)^b$	$-0.325(68)$	$-0.048(68)$
	p	530	1.334(8)	0.399(37)	$-0.247(17)^b$	0.284(81)	0.087(81)
[Berruto <i>et al</i> 2006]	n	690	1.575(9)	$-1.715(46)$	$-0.070(20)$	$-1.39(1.52)$	$-1.15(1.52)$
	n	605	1.470(9)	$-1.698(68)$	$-0.160(20)$	0.60(2.98)	1.14(2.98)
[Guo <i>et al</i> 2015]	n	465	1.246(7)	$-1.491(22)^c$	$-0.079(27)^d$	$-0.375(48)$	$-0.130(76)^d$
	n	360	1.138(13)	$-1.473(37)^c$	$-0.092(14)^d$	$-0.248(29)$	$0.020(58)^d$

After removing spurious contributions,

- *no lattice signal for θ_{QCD} -induced nEDM $\Leftrightarrow d_N$ is very small*
- *no conflict with phenomenology values or m_q scaling*

θ -Term Noise Reduction for EDM

Lattice signal for θ -nEDM $d_N \sim \langle Q \cdot (N(x) J_\mu \bar{N}(0)) \rangle_{CP-even}$

Top. charge Q is global $Q \sim \int_{V_4} (G\tilde{G})$ with $\langle |Q|^2 \rangle \sim V_4$

\Rightarrow Variance of correlator $\sim V_4$

Constrain Q integral to the relevant volume

- constrain Q in time, $|t_Q - t_J| \leq \Delta t$

[E.Shintani, T.Blum, T.Izubuchi, A.Soni, PRD93, 094503(2015)]

- "cluster decomposition" [K.-F.Liu et al, 1705.06358]:

constrain Q in 4-d around "sink" within $|r| < R$

Proper account of nucleon parity mixing is critical for correct determination of F_3

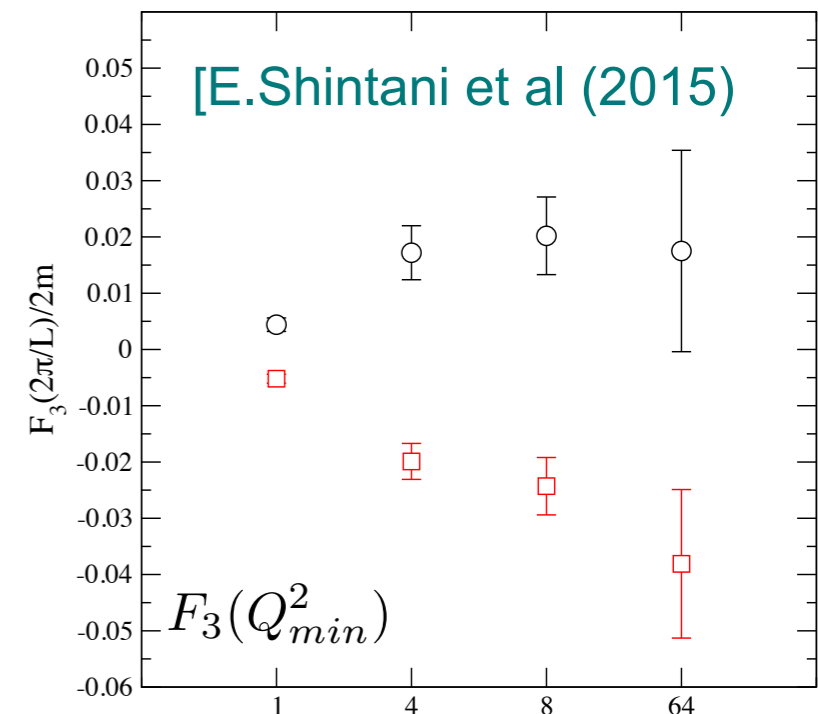
\Rightarrow nucleon states must "settle" in the new vacuum

$$N^{(+)} \rightarrow \tilde{N}^{(+)} \approx N^{(+)} + i\alpha N^{(-)}$$

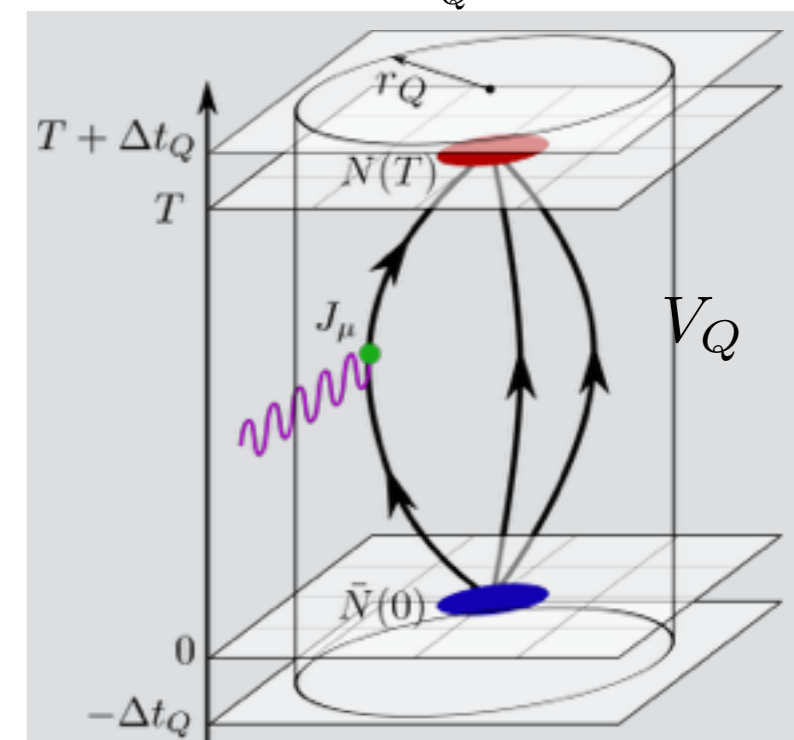
$$N^{(-)} \rightarrow \tilde{N}^{(-)} \approx N^{(-)} - i\alpha N^{(+)}$$

\Rightarrow treat time differently from space:

4d "cylinder" $V_Q : |\vec{z}| < r_Q, -\Delta t_Q < z_0 < T + \Delta t_Q$

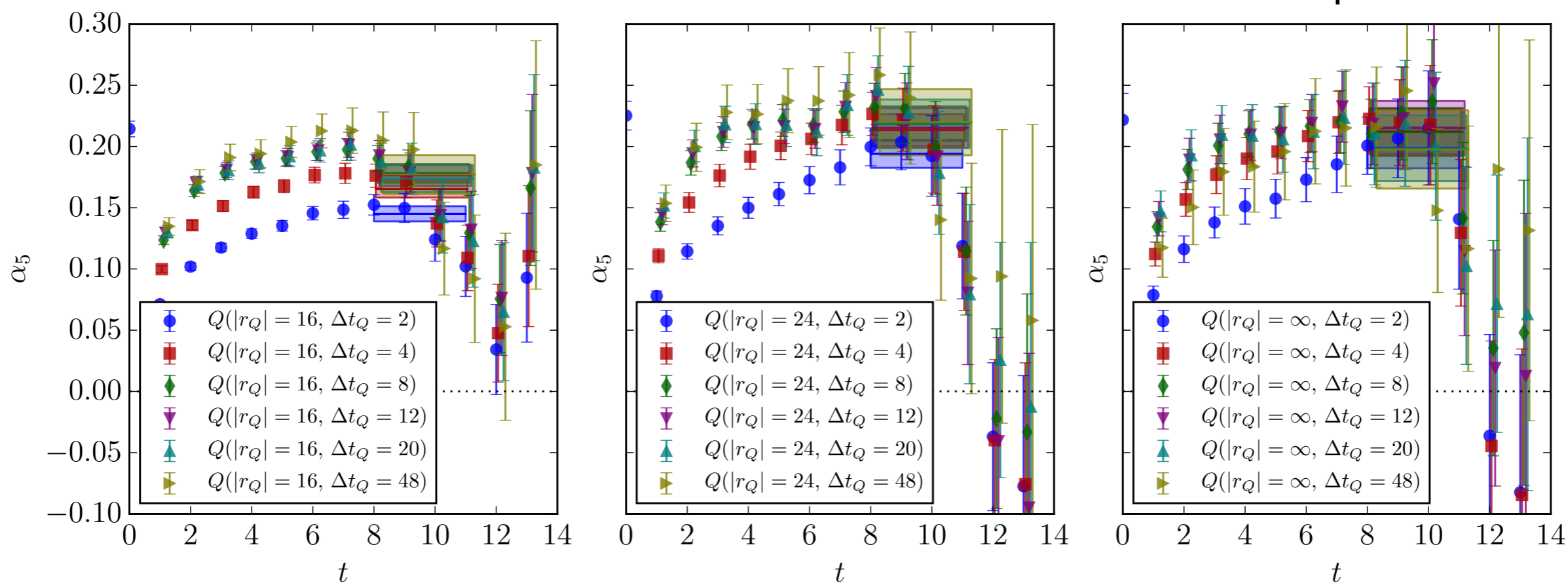


$$Q \approx \int_{V_Q} d^4 z q(z)$$



Noise Reduction: θ -induced Parity-mixing

PRELIMINARY 48c96 mpi=140MeV

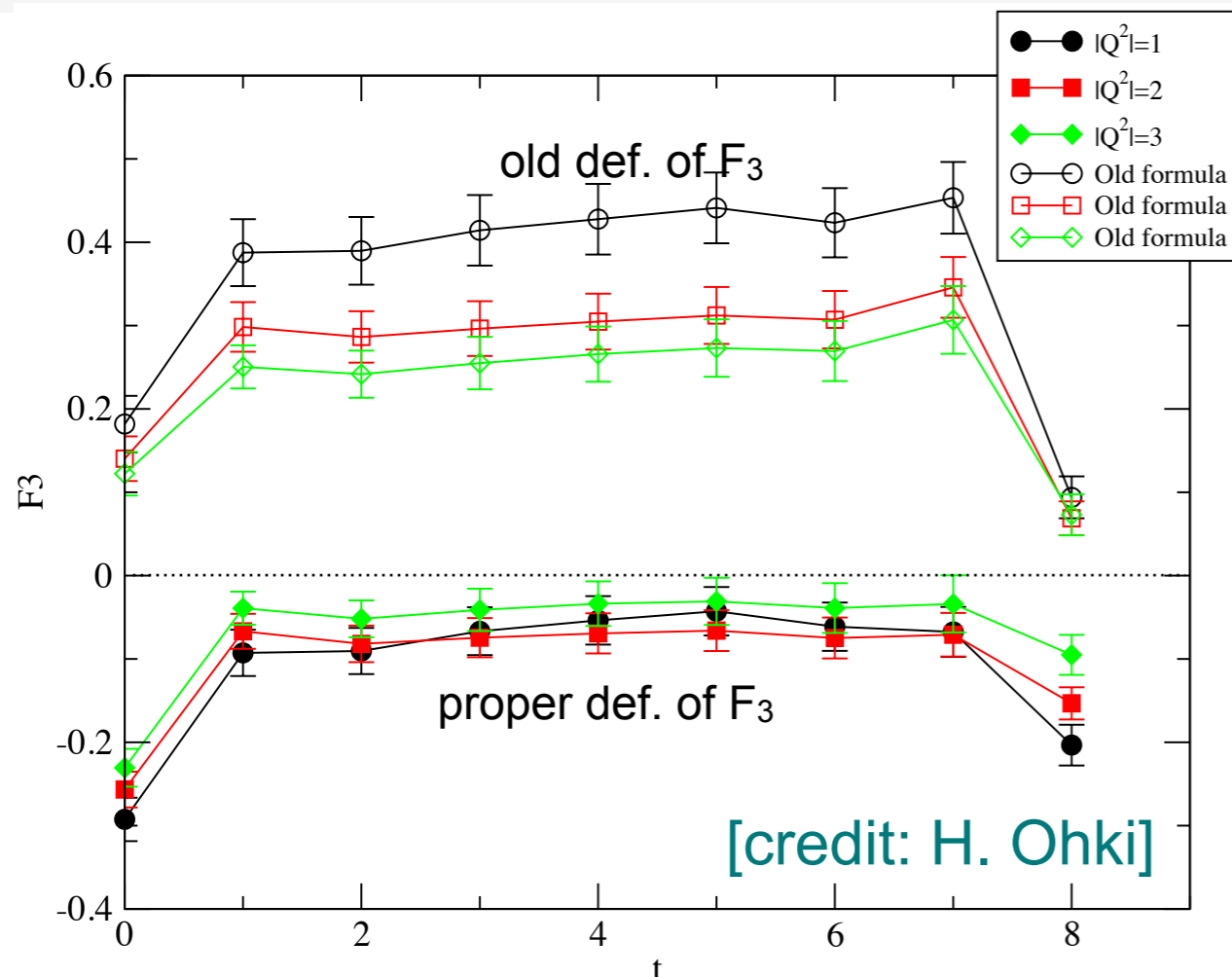


Parity-mixing angle from constrained Q sum

Reassuring results for noise reduction at the physical point

- required time region is small, $\Delta t_Q \gtrsim 8a \approx 1.2 \text{ fm}$
- spatial region must be large, $r_Q \gtrsim 20a \approx 2.3 \text{ fm}$

θ -nEDM Feasible at the Physical Point?



Preliminary Results with $m_\pi=330$ MeV

● Q sampled with $\Delta t_Q = 4a$, $r_Q = \infty$

Best guess for neutron EDM d_n : extrapolation in $m_q \sim (m_\pi)^2$

● chiral fermions, $m_\pi=330$ MeV

\implies phys.point $|F_3(0)| \approx 0.020$, $|d_n| \approx 0.002$ e fm

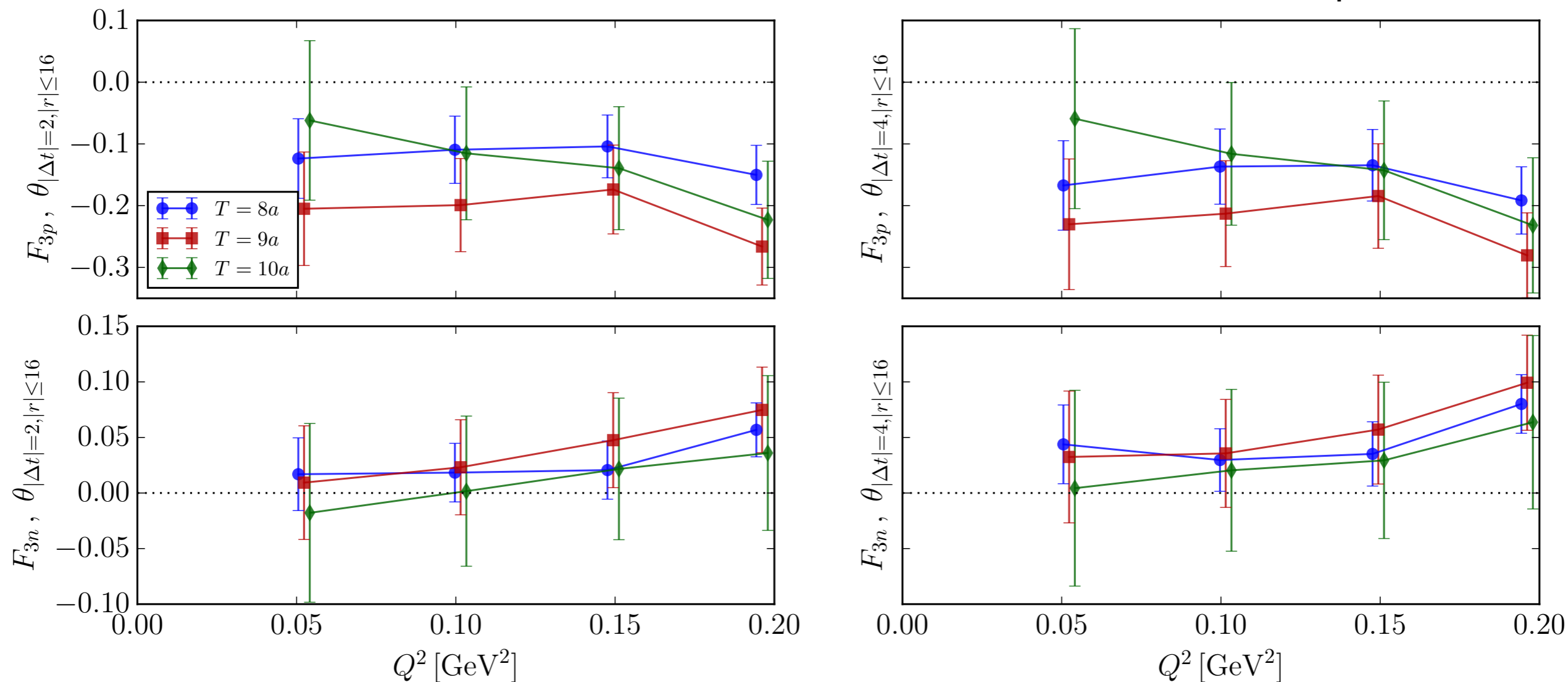
● Wilson fermions, $m_\pi=360$ MeV [Guo et al 2015]

\implies phys.point $|F_3(0)| \leq 0.012$, $|d_n| \leq 0.001$ e fm

$$|F_{3n}^{\text{phys}}(0)| \sim O(10^{-2}) \theta, \quad |d_n| \sim O(10^{-3}) \text{ e fm } \theta$$

Noise Reduction: θ -induced EDFF F3

PRELIMINARY 48c96 mpi=140MeV



- EDFF F₃ from constrained Q sum: *the most aggressive Q cuts*
- 33k lattice samples, ~ 30 M core-hours on Argonne BlueGene/Q
- connected diagrams only
- result compatible with zero, $|F_{3n}| \leq 0.05$

Need to constrain $|F_{3n}| \approx 0.01..0.02$: θ -nEDM remains difficult at the physical point...

Outlook for θ -nEDM

Resort to simpler calculations

- heavier pion masses + EFT for extrapolations
- quenched calculations (see e.g. recent [J.Dragos et al,1711.04730])

Physical point calculations of θ -nEDM will be necessary to renormalize effects from other CPv sources of higher-dim. [T.Bhattacharya et al (2015)]

New lattice simulations at the physical point with dynamical θ -term

- coarse ($a=0.2$ fm) physical-point lattice \implies reduced cost due to lattice volume
- chiral lattice fermions allow independent $a \rightarrow 0$, $m_q \rightarrow 0$ limits
- enhance d_N signal with $\langle Q \rangle \neq 0$ – more critical at light quark masses

\implies 2018 ALCC award for 50 M BG/Q core-hours

Ensembles with dynamical θ -term will be also useful for CPv πN coupling

Another Source of CPv: Quark Chromo-EDM

$$\mathcal{L}_{\text{cEDM}} = \sum_{q=u,d} \frac{\tilde{\delta}_q}{2} \bar{q} [G_{\mu\nu} \sigma^{\mu\nu} \gamma_5] q$$

- $O(a^{-2})$ mixing with dim-3 pseudoscalar density
 \Rightarrow need non-perturbative subtractions

- Non-chiral (e.g. Wilson) fermions have a $O(a)$ clover term ("chromo-magnetic DM")

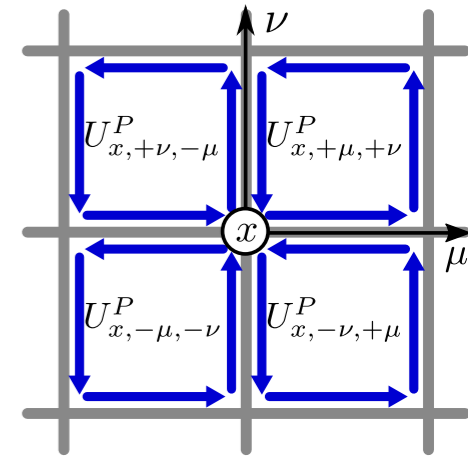
$$\mathcal{L}^{\text{clover}} = a \frac{c}{4} \bar{q} [G_{\mu\nu} \sigma^{\mu\nu}] q$$

Condensate realignment in presence of CPv $q \rightarrow e^{i\gamma_5 \Omega} q$
 assuring $\langle \text{vac} | \mathcal{L}_m + \mathcal{L}_{\text{CP}} | \pi^a \rangle = 0$

mixes (chromo)EDM and (chromo)MDM:

$$\delta \mathcal{L}_{\text{cEDM}} = \delta(\bar{q} [\tilde{D}_q G_{\mu\nu} \sigma^{\mu\nu} \gamma_5] q) = \bar{q} [\{\Omega, \tilde{D}_q\} G_{\mu\nu} \sigma^{\mu\nu}] q \sim \delta \mathcal{L}_{\text{cMDM}}$$

\Rightarrow *Chirally-symmetric actions avoid these cMDM contributions*

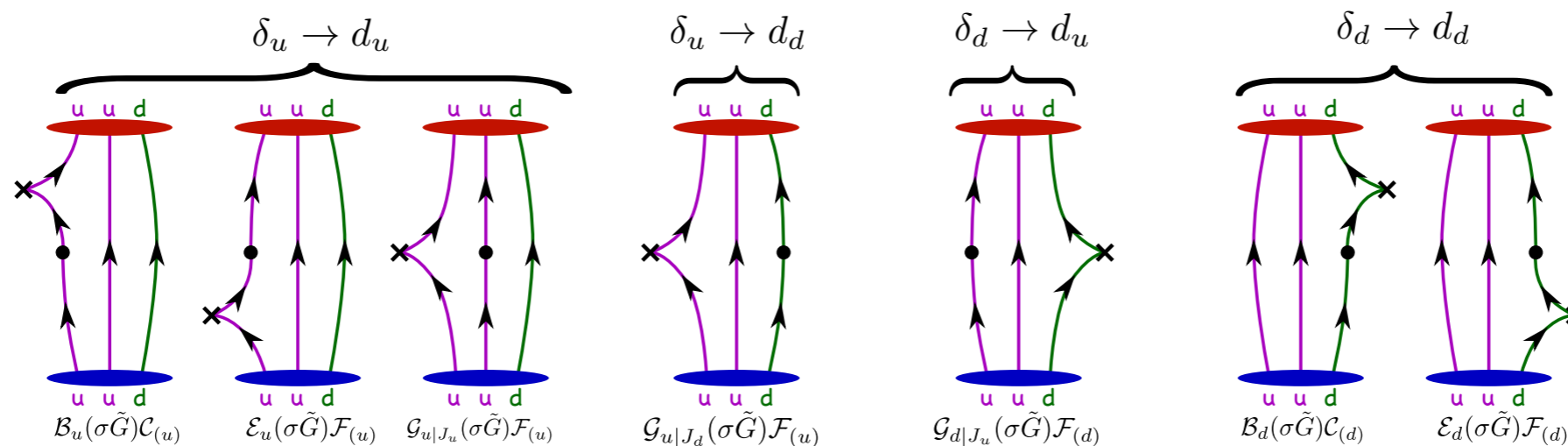


Quark-Gluon EDM: Insertions of dim-5 Operators

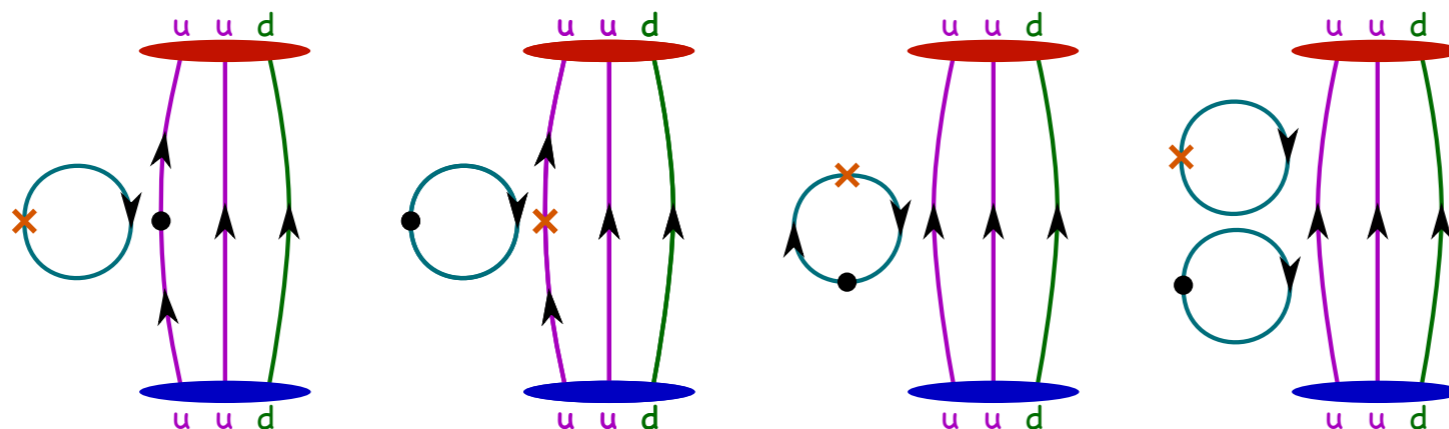
$$\mathcal{L}^{(5)} = \sum_q \tilde{d}_q \bar{q}(G \cdot \sigma)\gamma_5 q \quad \begin{matrix} \nearrow \\ \searrow \end{matrix} \begin{matrix} \langle N(y) \bar{N}(0) \int d^4x \bar{q}(G \cdot \sigma)\gamma_5 q \rangle \\ \langle N(y) [\bar{\psi}\gamma^\mu\psi]_z \bar{N}(0) \int d^4x \bar{q}(G \cdot \sigma)\gamma_5 q \rangle \end{matrix}$$

First calculations : [T.Bhattacharya et al(LANL, LATTICE'15,'16)]

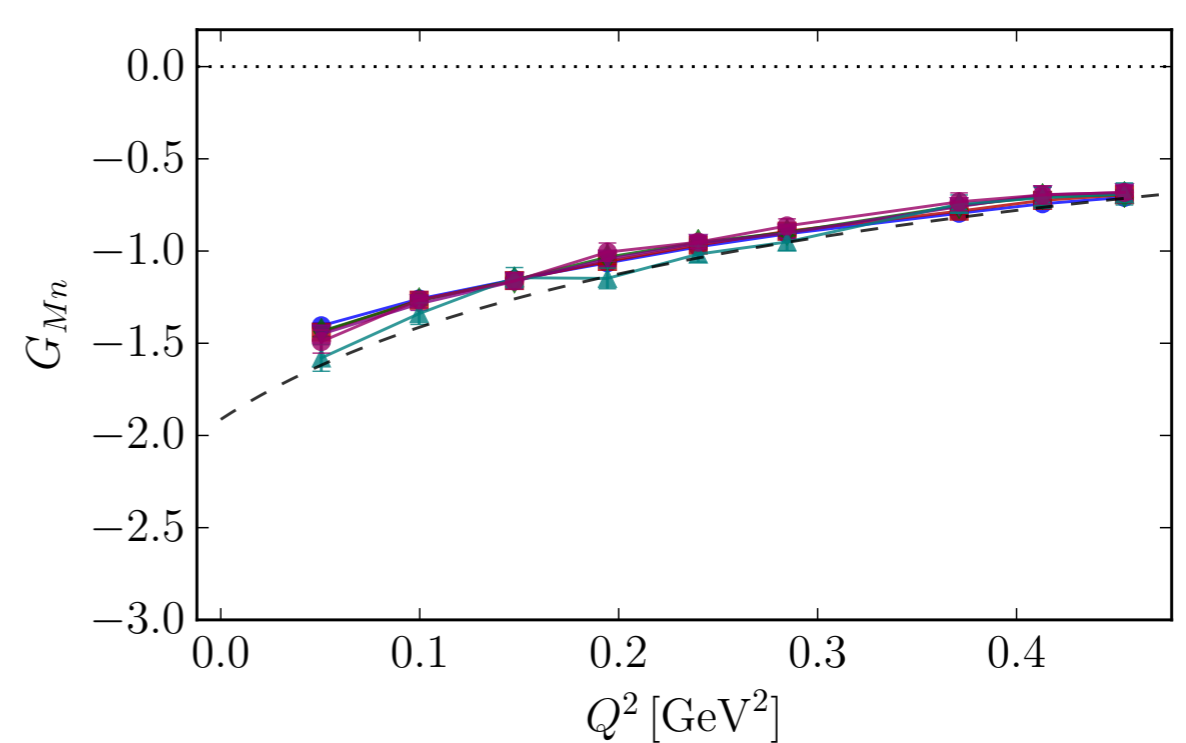
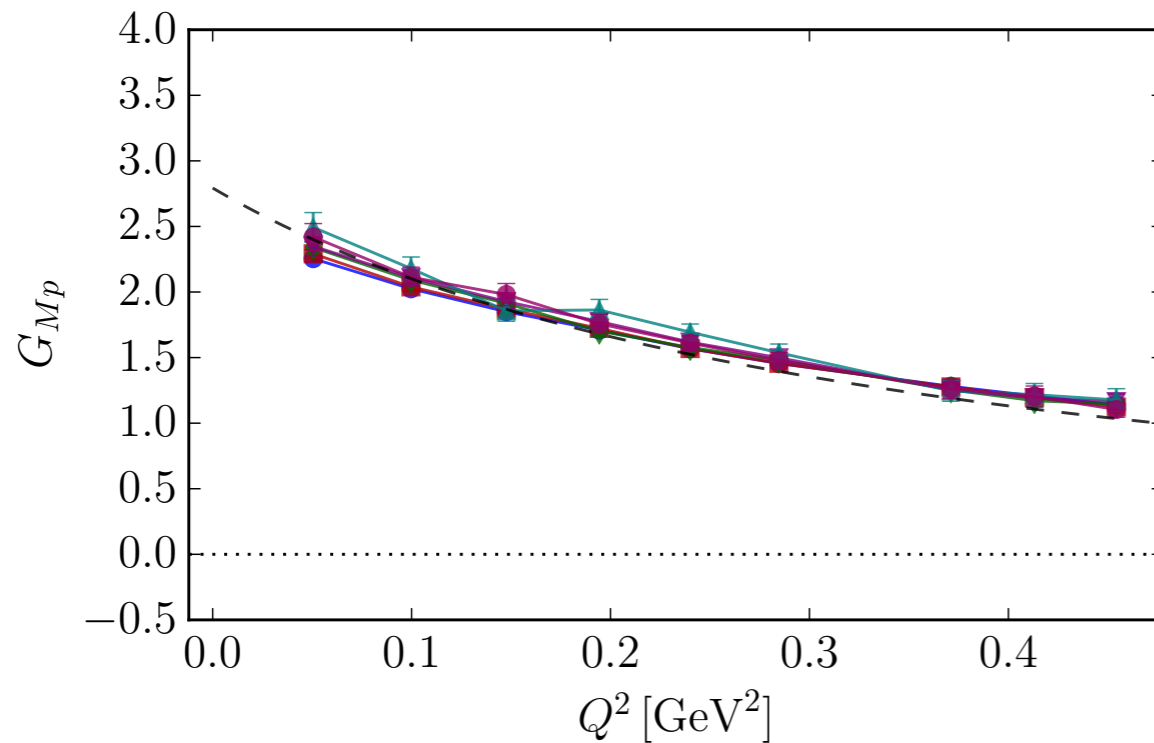
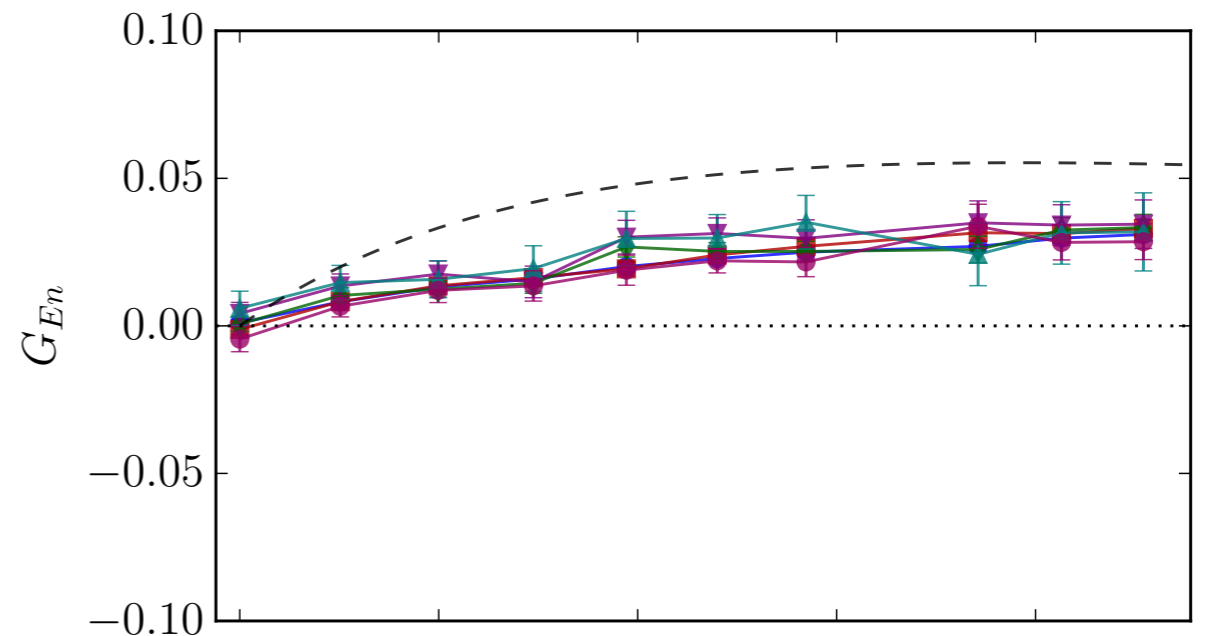
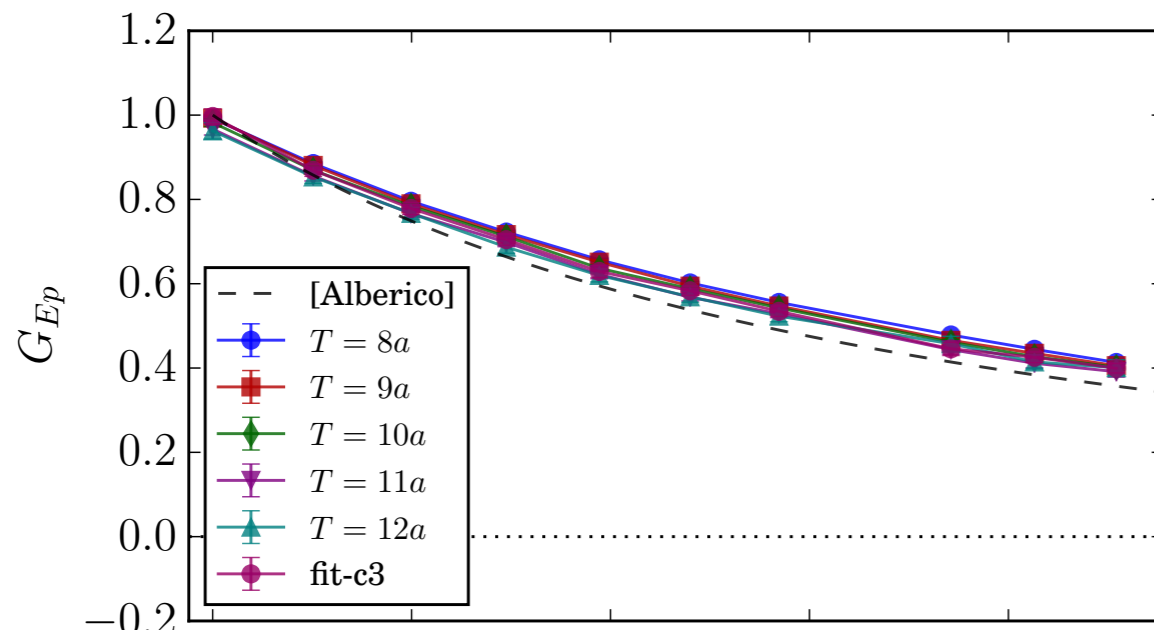
- This work: Only quark-connected insertions



- In future: Single- and double-disconnected diagrams (contribute to isosinglet cEDM, mix with θ -term)



Nucleon Sachs Form Factors

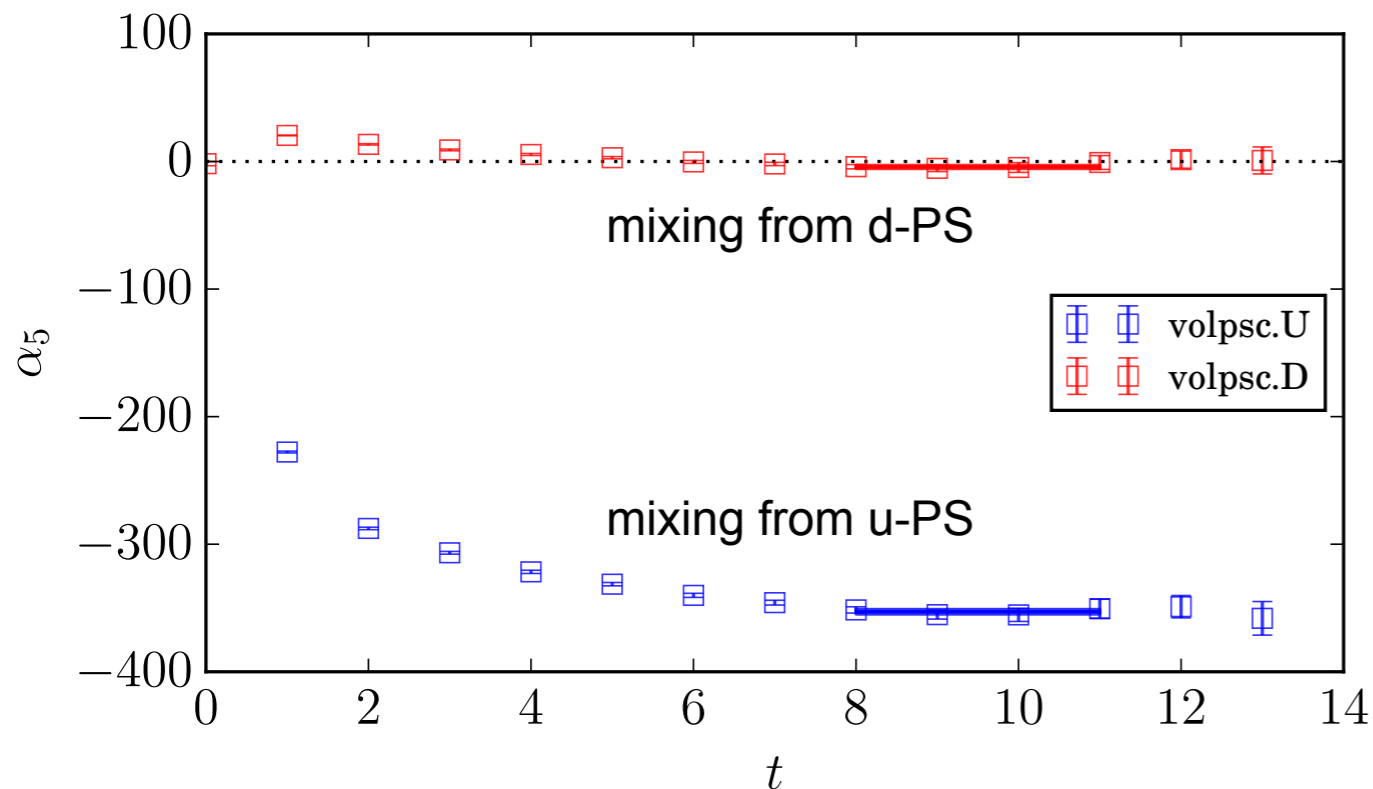
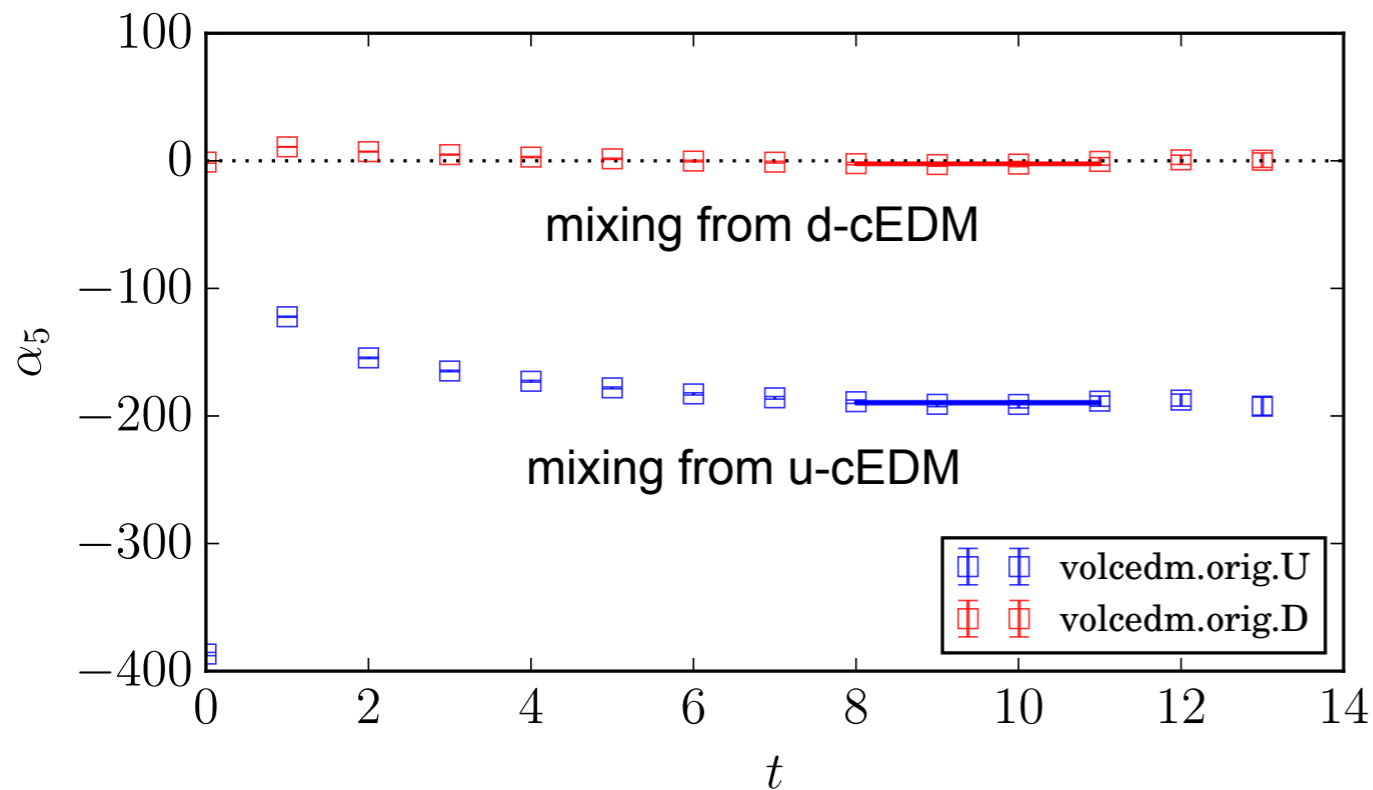


$$G_E = F_1 - \frac{Q^2}{4m_N^2} F_2$$

$$G_M = F_1 + F_2$$

- $(5.5 \text{ fm})^3 \times (11 \text{ fm})$ box, $m_\pi = 140 \text{ MeV}$
- connected-only contractions

Parity Mixing (Proton)



$$N_\delta = \epsilon^{abc} u_\delta^a (u^{aT} \mathcal{C} \gamma_5 d^c)$$

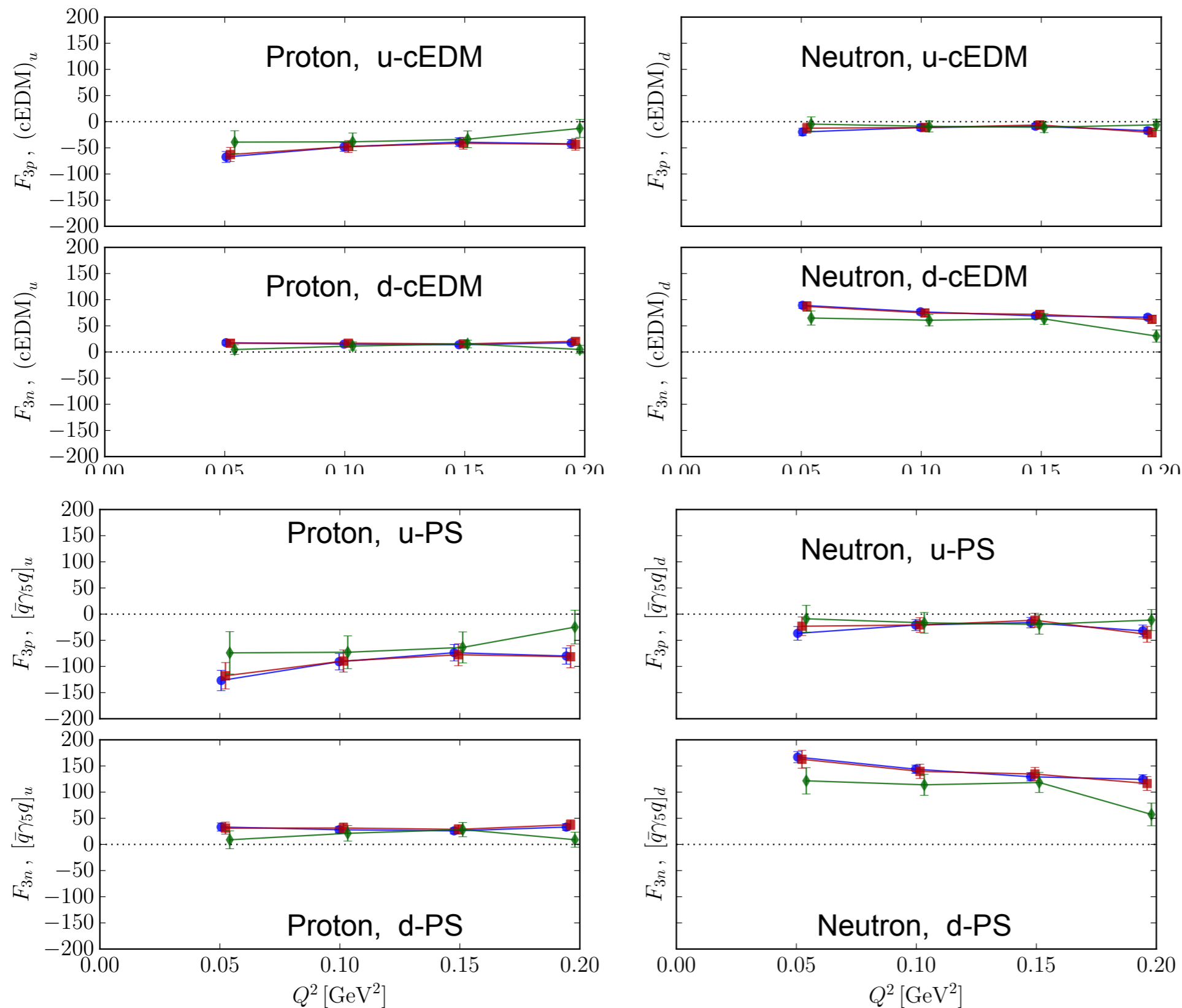
$$\langle N(t) \bar{N}(0) \rangle_{\mathcal{CP}} = \frac{-i\not{p} + m_N e^{2i\alpha_5 \gamma_5}}{2m_N} e^{-E_N t}$$

$$\hat{\alpha}_5 = \frac{\alpha_5}{\tilde{d}} = -\frac{\text{ReTr}[T^+ \gamma_5 \cdot C_{2pt}^{\overline{\mathcal{CP}}}(t)]}{\text{ReTr}[T^+ \cdot C_{2pt}^{\mathcal{CP}}(t)]}, \quad t \rightarrow \infty$$

(flavors labeled for the proton)

similarity effect on nucleon likely due to mixing between cEDM and PS

Proton & Neutron EDFF Form Factors (bare)



- (5.5 fm)³x(11 fm) box
- $m_\pi=140$ MeV
- connected-only
- no renormalization

Nucleon EDM : Summary

- Previously reported lattice results for θ_{QCD} -induced nEDM contain spurious contributions from mixing with the anomalous mag.moment
- Corrected θ_{QCD} -nEDM lattice values are small, consistent with zero
 - Disagreement with phenomenology/EFT is eliminated
 - Much higher lattice statistics are required to constrain θ_{QCD}
- Based on *preliminary* analysis at a heavier pion mass (330 MeV), at the physical point expect $|d_n| \approx (1..2) \times 10^{-3}$ e fm
 - Even with variance-reduction techniques, O(300) M core*hours may be required
- Promising results for *quark* cEDM-induced EDFF
 - Renormalization & mixing subtractions are underway*