

Searches for Dark Matter mediators with the ATLAS Detector

Peter McNamara

On behalf of the ATLAS Collaboration



CoEPP

ARC Centre of Excellence for
Particle Physics at the Terascale

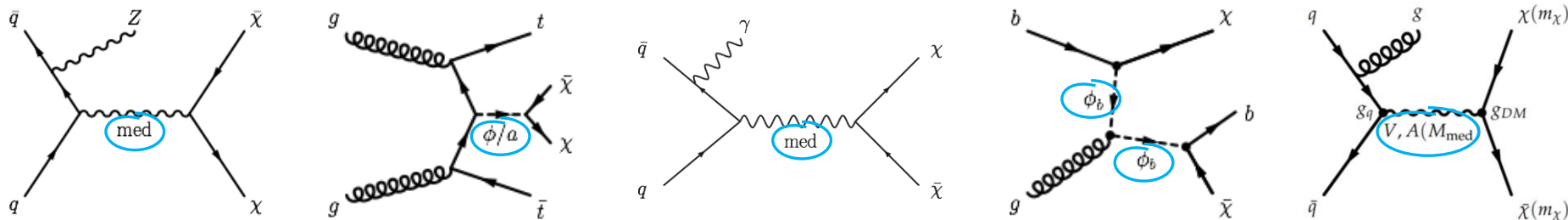


THE UNIVERSITY OF
MELBOURNE

Dark Matter searches at the LHC use simplified models

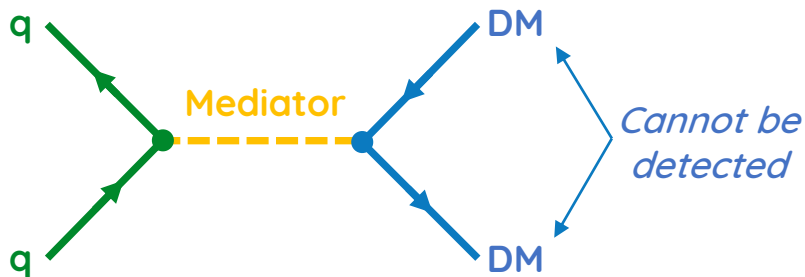
- Production of Dark matter goes via a mediator

See Young-Kee Kim's talk

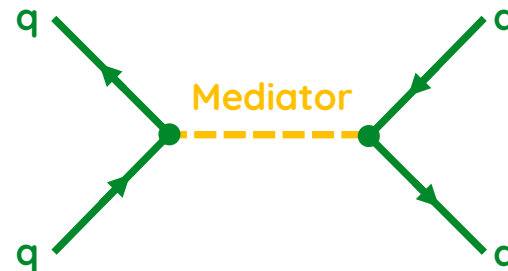


Resonance searches are complimentary to these direct DM searches

- Look for mediators decaying into quarks (dijet resonance)



Mediators produced by quarks ...
(which decay to Dark Matter)



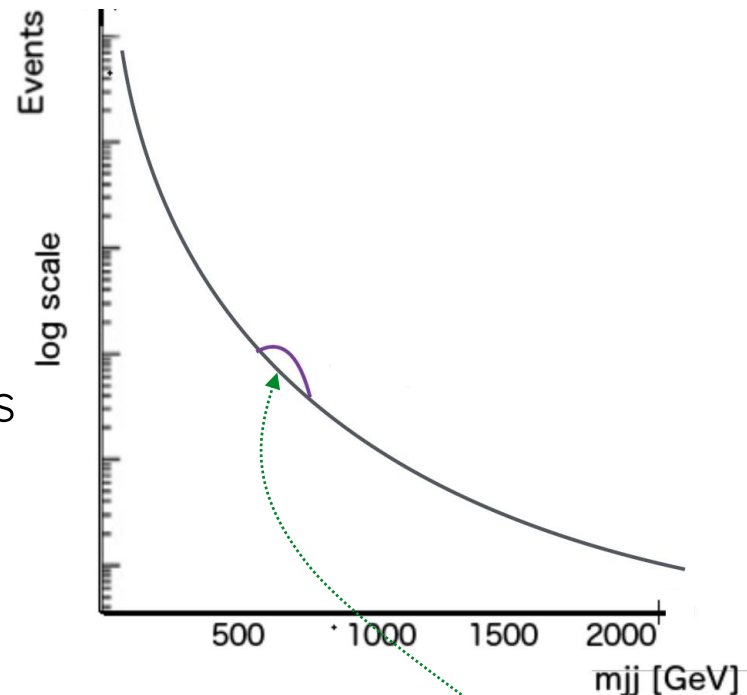
also decay to quarks
(producing jets)

QCD predicts smoothly falling dijet invariant mass (m_{jj})

- New mediator decaying to jets will introduce a bump

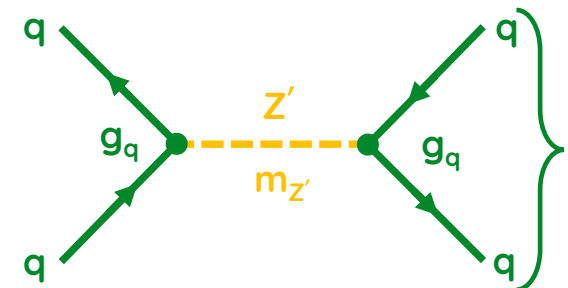
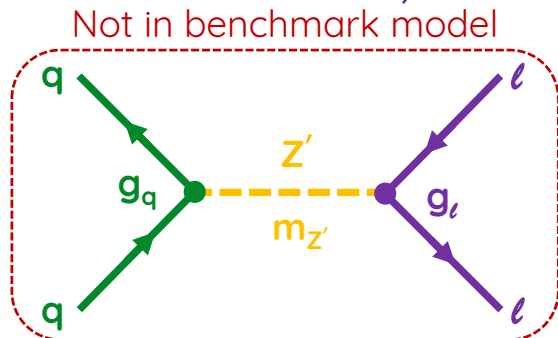
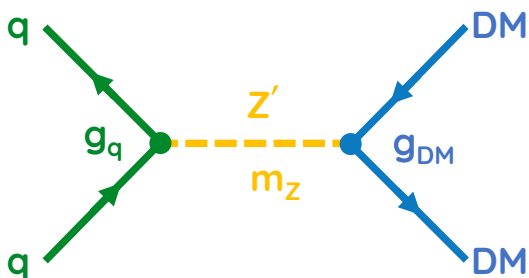
Benchmark Leptophobic Z' model parameters

- Axial-vector (or vector) mediator (Z')
- Coupling to quarks - g_q (universal)
- Coupling to leptons - $g_\ell = 0$ (leptophobic)
- Coupling to dark matter - g_{DM}
- Mediator mass - $m_{Z'}$
- Dark matter mass - m_{DM}



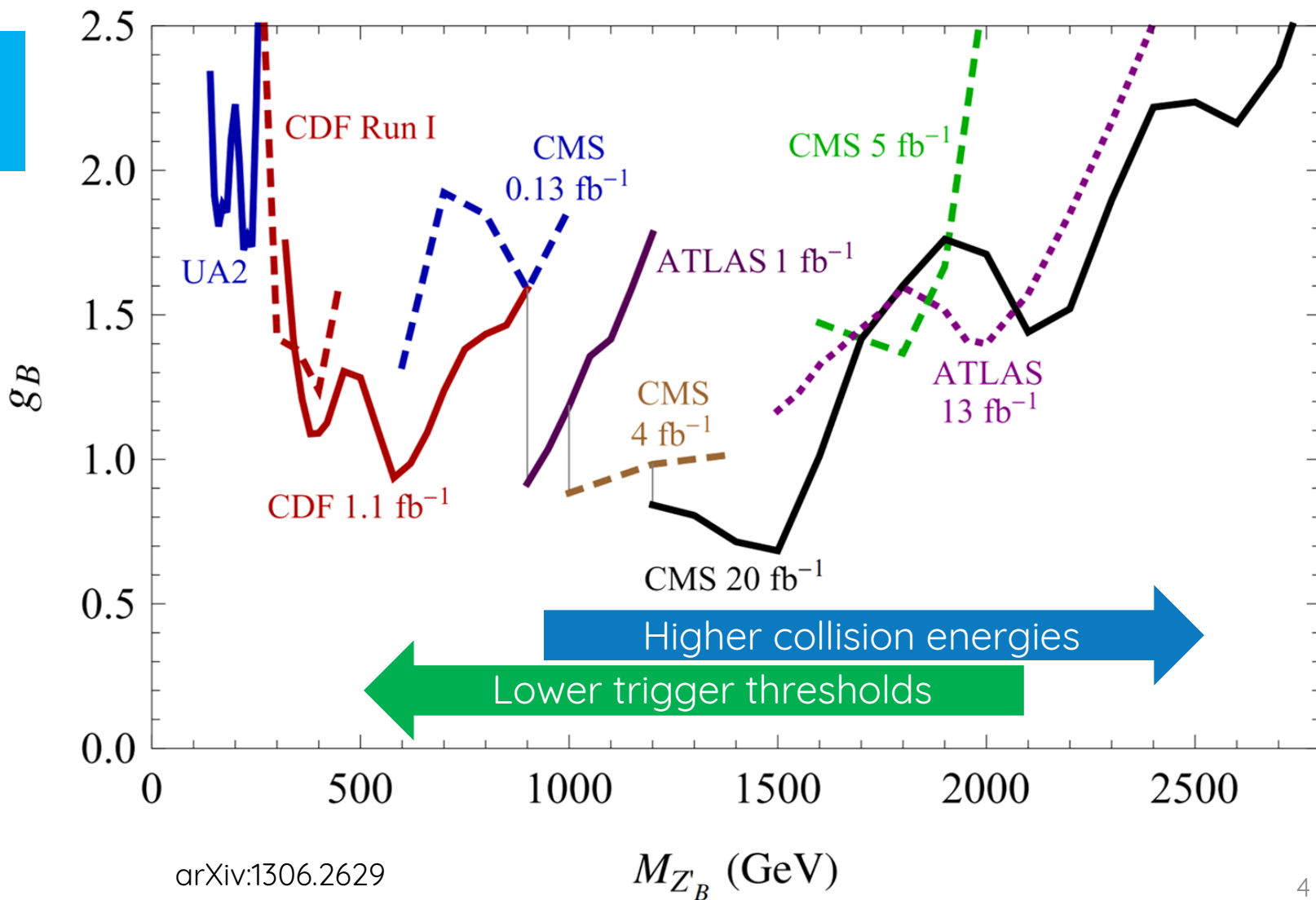
See Sébastien Rettie's talk for dilepton resonance searches

Dijet Resonance



Limits are dependent on a number of factors

More Events Recorded



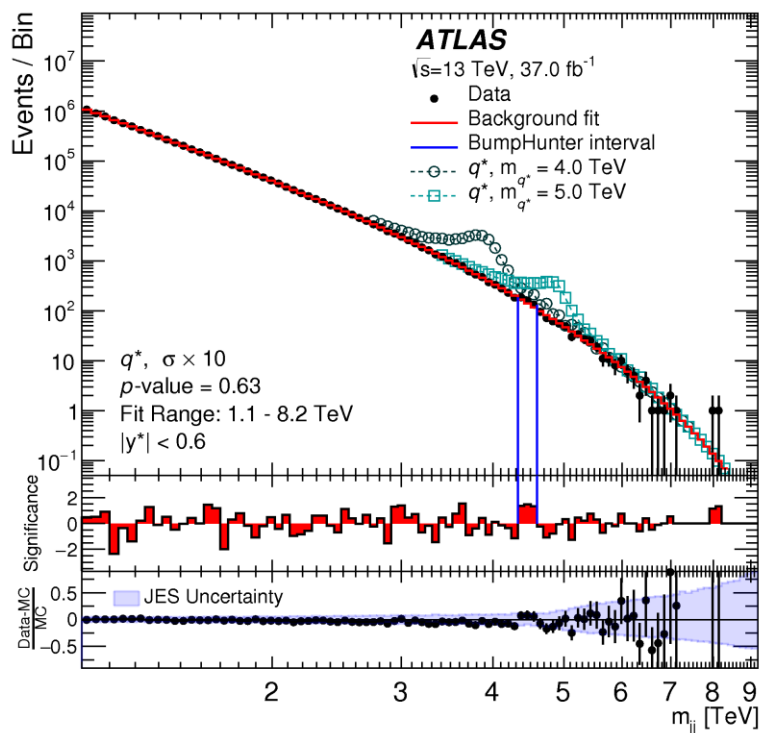
Selections are applied to

- Ensure trigger and selection efficiency
- Reduce backgrounds (y^* rapidity difference)

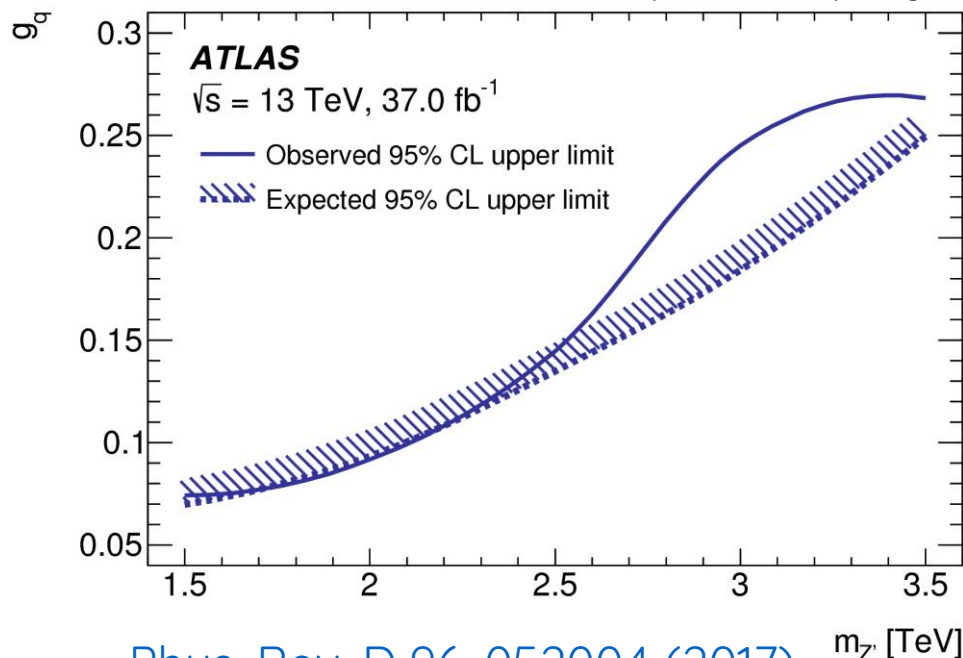
Estimate background from sliding window fit using $f(z) = p_1(1 - z)^{p_2} z^{p_3}$. $z = m_{jj} / \sqrt{s}$

Search for excess with bumpHunter algorithm

Selections	
Lead Jet p_T	> 440 GeV
Second Jet p_T	> 60 GeV
m_{jj}	> 1.1 TeV
$ y^* = 0.5 y_1 - y_2 $	< 0.6

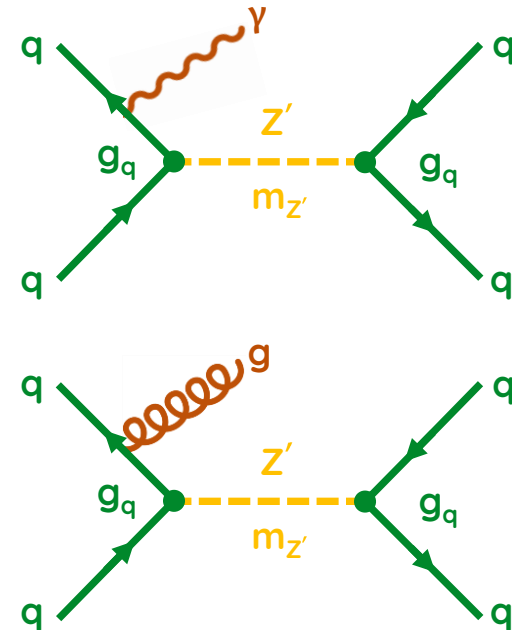
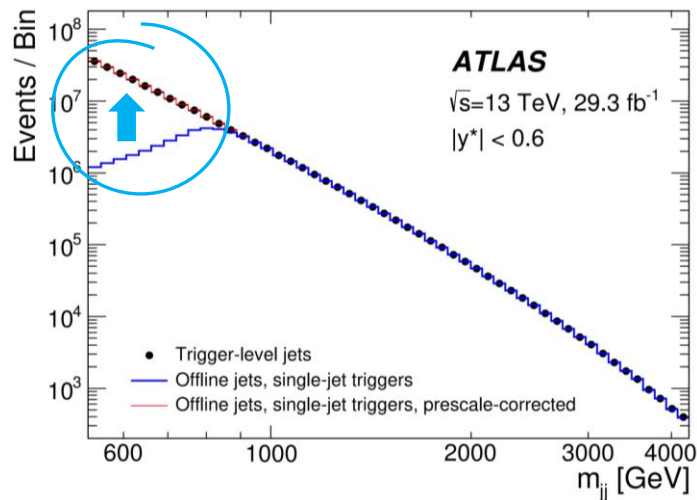


Limit on axial vector Z' quark coupling



[Phys. Rev. D 96, 052004 \(2017\)](#)

How do you access lower masses?



Trigger Level Analysis (TLA)

Lower the trigger threshold by reducing amount of data saved per event, keep only trigger level objects

Initial State Radiation (ISR) Selection

Examine a boosted signature by requiring an ISR photon or jet to pass the trigger threshold

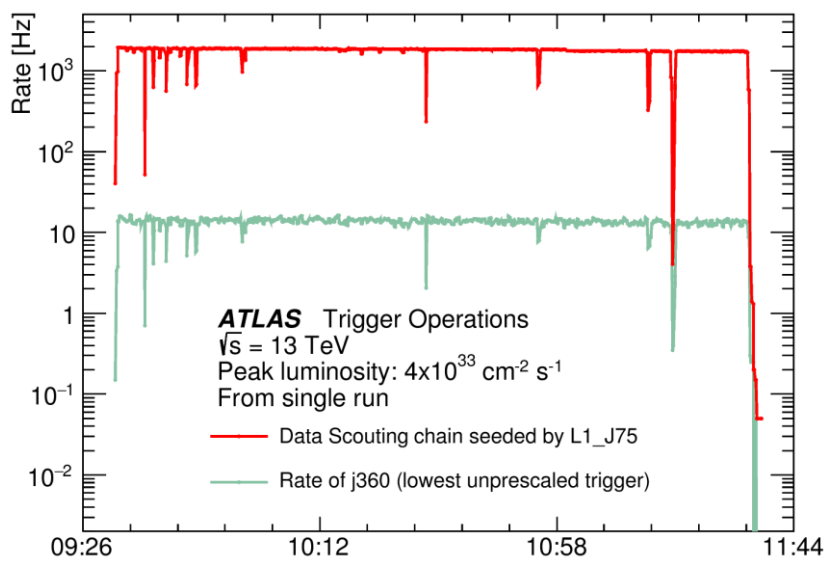
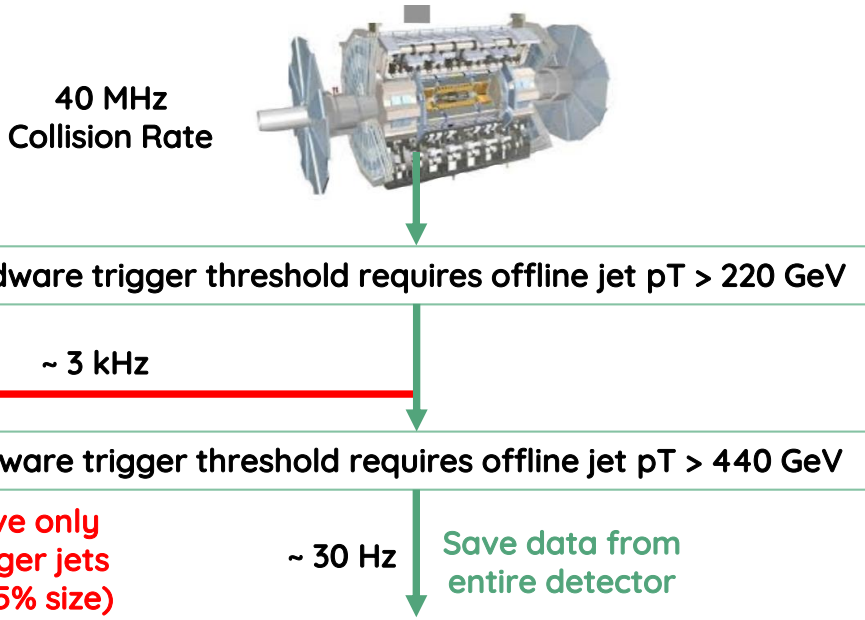
First stage hardware trigger identifies jets using coarse calorimeter information

Second stage software trigger then reconstructs jets using same algorithm as offline jets (anti kt)

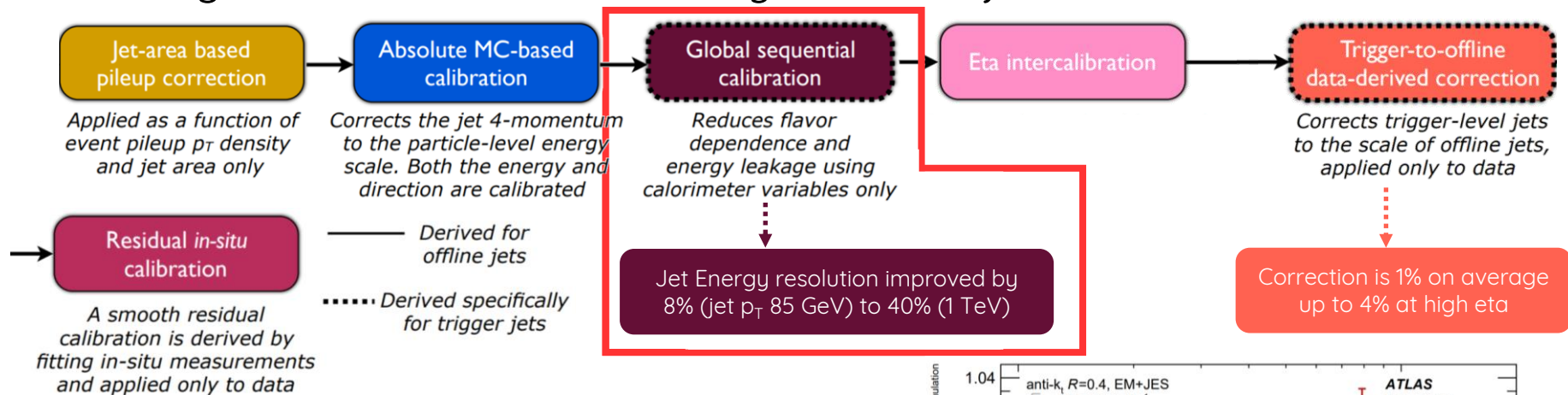
Limited bandwidth so high threshold required to save entire event

Save only trigger jets

- Small size
- Higher rate →
- Similar bandwidth

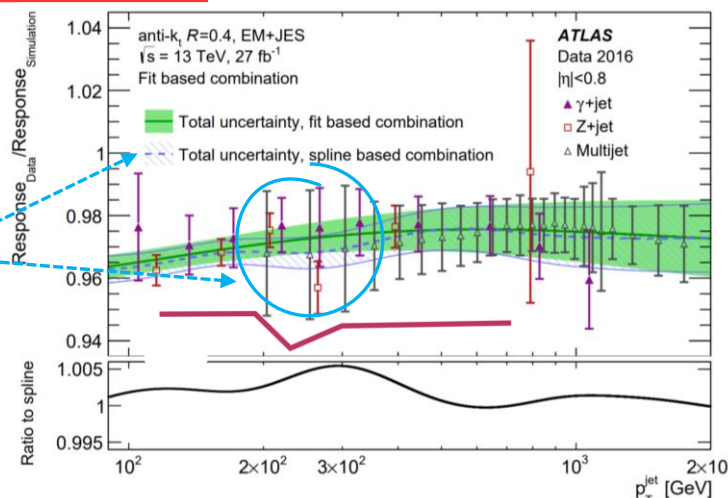
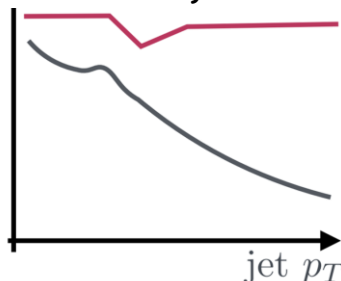


Jet energies are calibrated similarly to offline jets ([Eur. Phys. J. C 76 \(2016\) 581](#))



Calibrate using momentum balance of offline jets against well calibrated objects

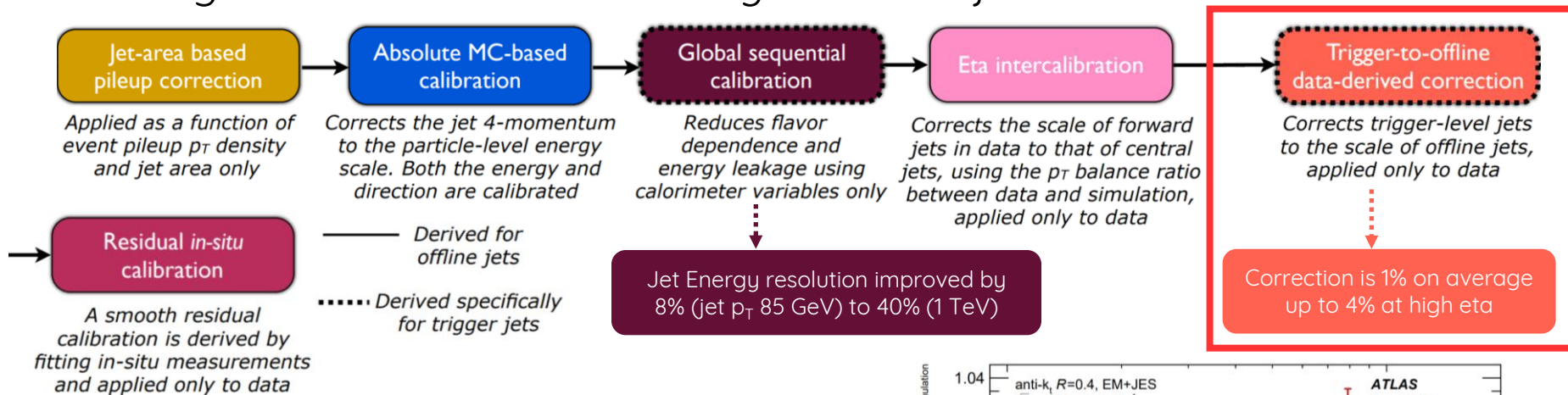
- Needs to be smooth as can introduce bumps
- Use polynomial fit in $\log(p_T)$ rather than spline



After calibration energy of trigger and offline jets agree within 0.05%

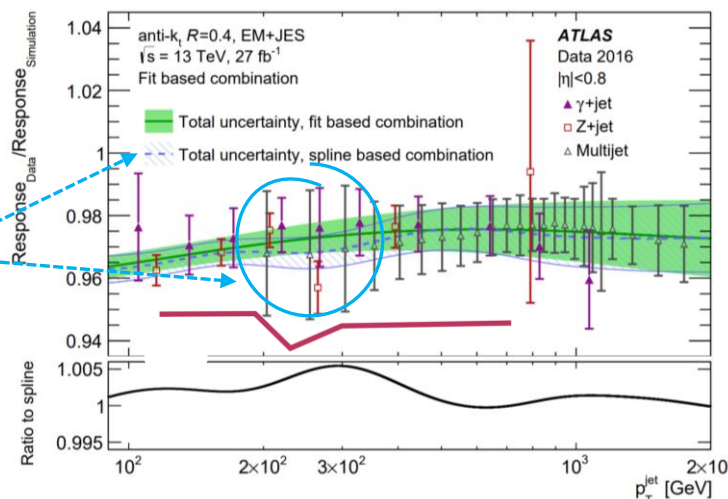
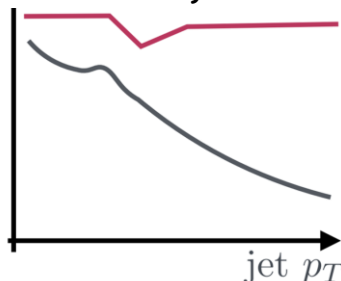
Jet energy scale uncertainty is 1-2%, similar to that of offline jets

Jet energies are calibrated similarly to offline jets ([Eur. Phys. J. C 76 \(2016\) 581](#))



Calibrate using momentum balance of offline jets against well calibrated objects

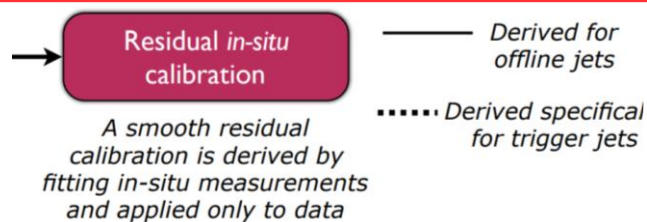
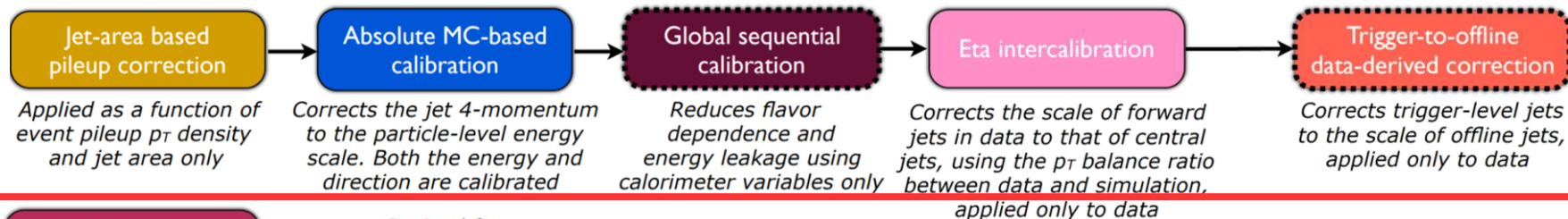
- Needs to be smooth as can introduce bumps
- Use polynomial fit in $\log(p_T)$ rather than spline



After calibration energy of trigger and offline jets agree within 0.05%

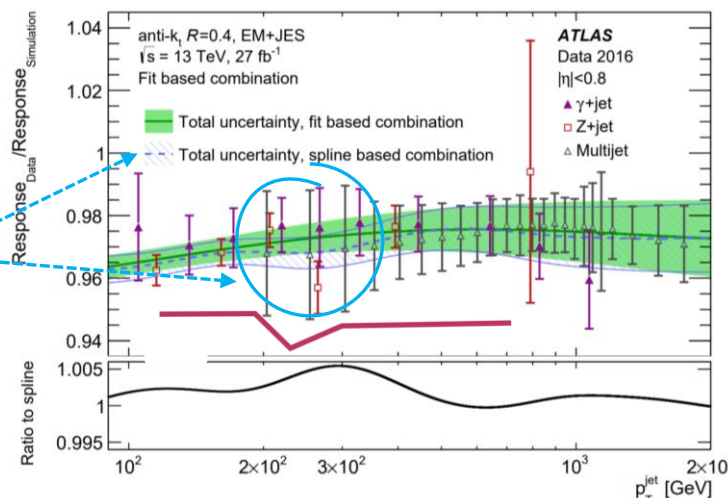
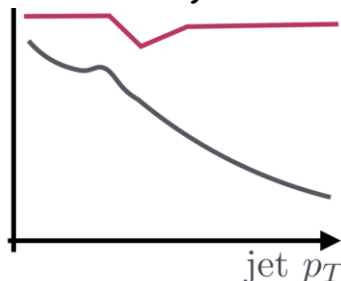
Jet energy scale uncertainty is 1-2%, similar to that of offline jets

Jet energies are calibrated similarly to offline jets ([Eur. Phys. J. C 76 \(2016\) 581](#))



Calibrate using momentum balance of offline jets against well calibrated objects

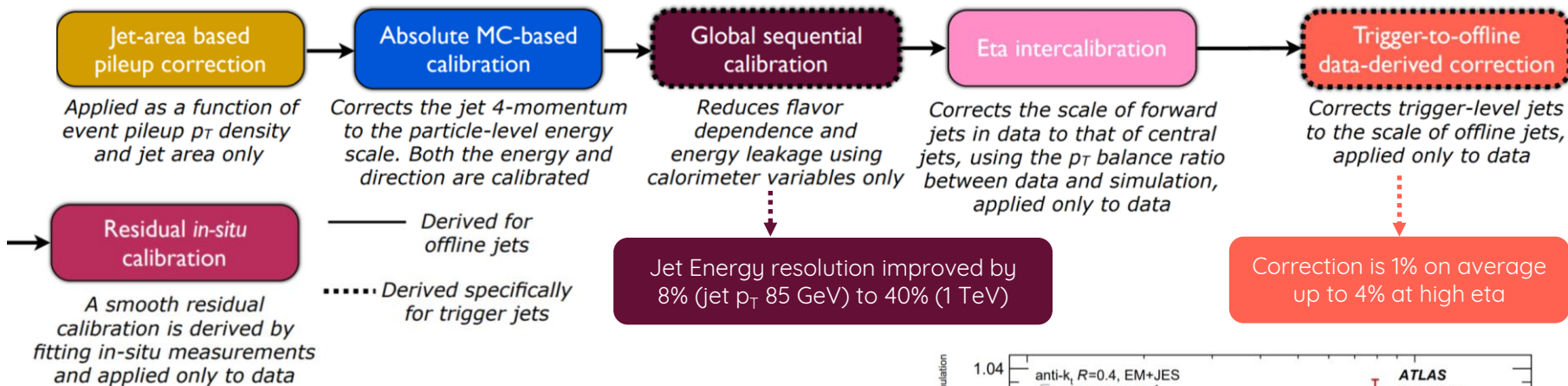
- Needs to be smooth as can introduce bumps
- Use polynomial fit in $\log(p_T)$ rather than spline



After calibration energy of trigger and offline jets agree within 0.05%

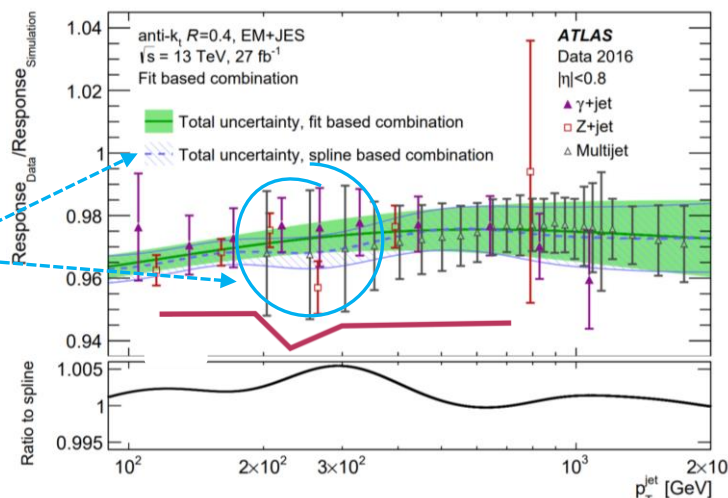
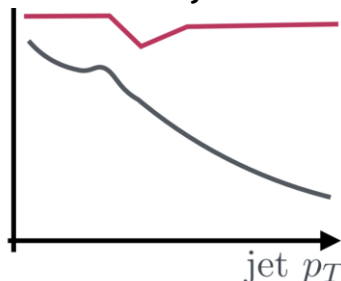
Jet energy scale uncertainty is 1-2%, similar to that of offline jets

Jet energies are calibrated similarly to offline jets ([Eur. Phys. J. C 76 \(2016\) 581](#))



Calibrate using momentum balance of offline jets against well calibrated objects

- Needs to be smooth as can introduce bumps
- Use polynomial fit in $\log(p_T)$ rather than spline



After calibration energy of trigger and offline jets agree within 0.05%

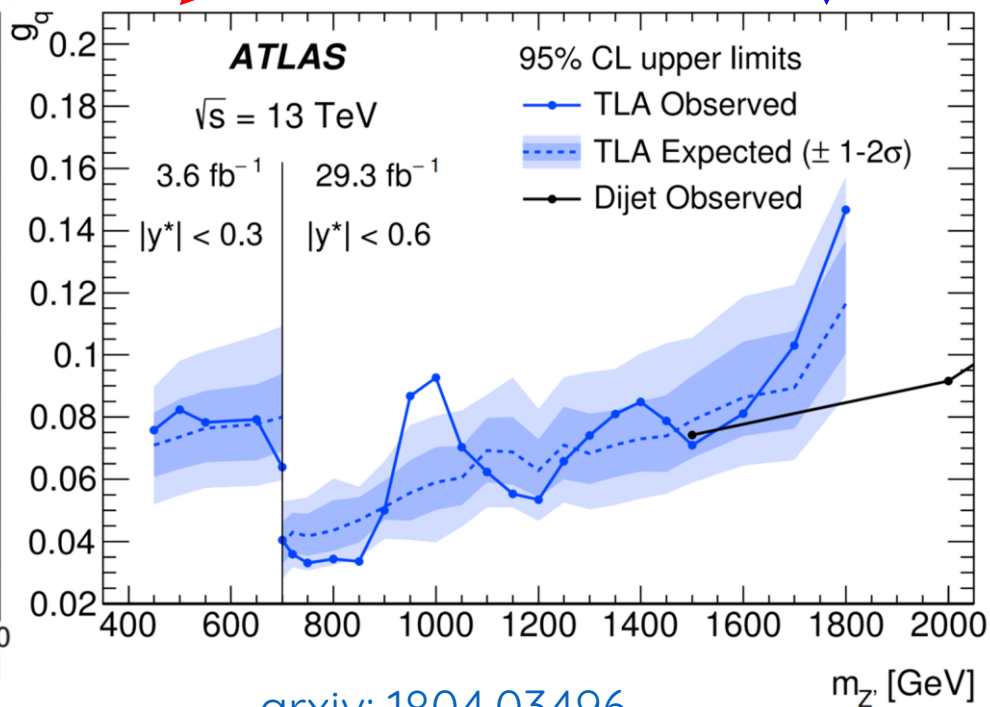
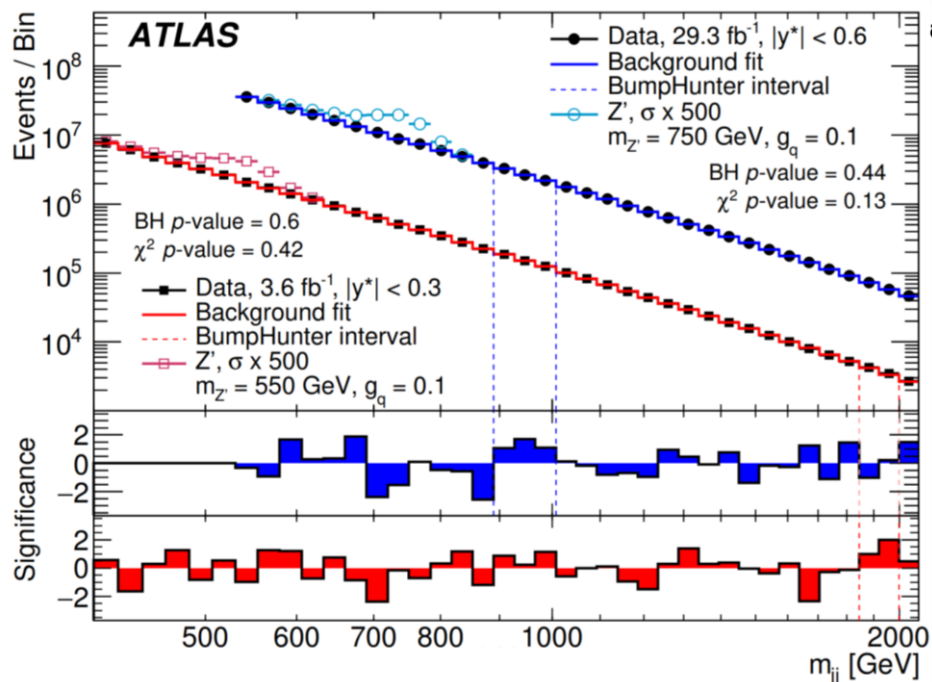
Jet energy scale uncertainty is 1-2%, similar to that of offline jets

Two selections based on trigger threshold target different mass ranges

Fit background and look for bumps

Set Bayesian limit

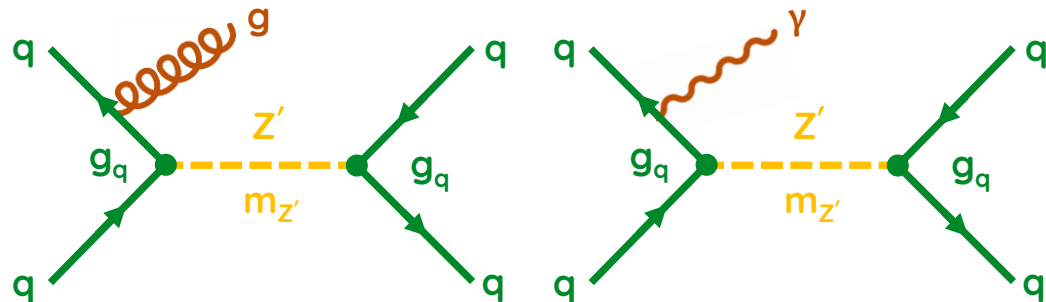
Selections	HW 75 GeV Trigger (3.6 fb ⁻¹)	HW 100 GeV Trigger (29.3fb ⁻¹)
Lead Jet pT	> 185 GeV	> 220 GeV
Second Jet pT	> 85 GeV	> 85 GeV
m_{jj}	> 450 GeV	> 700 GeV
$ y^* = 0.5 y_1 - y_2 $	< 0.3	< 0.6



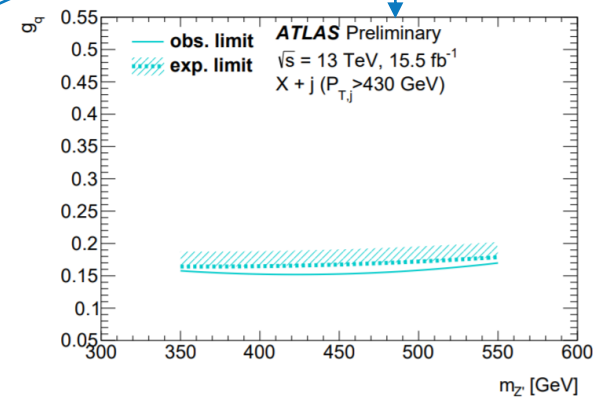
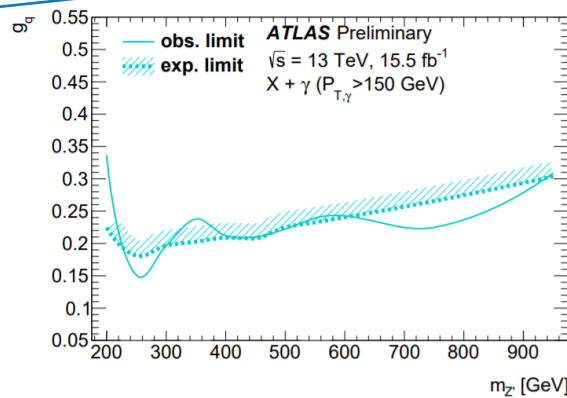
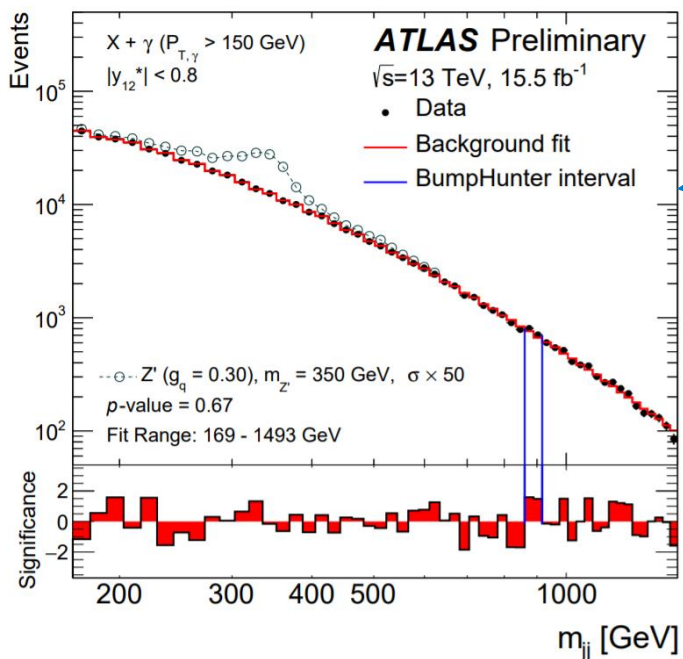
[arxiv:1804.03496](https://arxiv.org/abs/1804.03496)

Add energetic ISR photon or jet to signature to allow better sensitivity to light resonances

- satisfies the trigger threshold
- allows lower masses to be examined
- reduced production rates



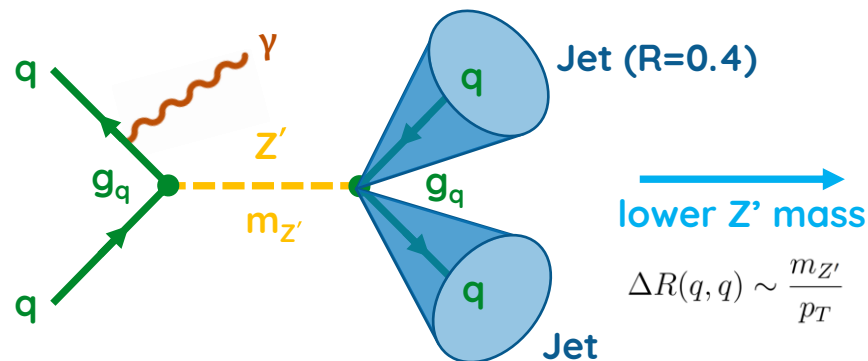
Selections	ISR Photon	ISR Jet
ISR Object	> 150 GeV	> 430 GeV
First (Second) Jet p_T	> 25 GeV	
Second (Third) Jet p_T	> 25 GeV	
$ y^* = 0.5 y_1 - y_2 $	< 0.8	< 0.6



ISR selection has allowed masses down to 200 GeV to be examined

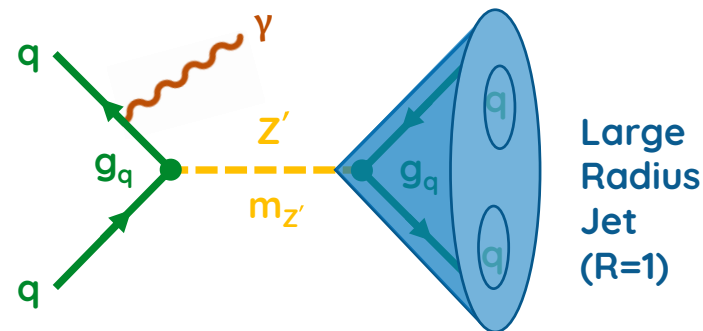
- Targets regime where jets from the mediator are separated
- To go lower examine a larger single jet instead

ISR dijet search



Less boosted by ISR, separate jets

Boosted ISR dijet search



More boosted by ISR, jets are more collimated
Reconstruct as single large radius jet

To reduce pileup and soft radiation effects, large radius jets are trimmed

- Reclustered with smaller radius using kt algorithm
- Removed if smaller clusters if carry < 5% of total jet p_T ([Eur. Phys. J. C 76 \(2016\) 154](https://arxiv.org/abs/1508.04099))

Use jet substructure to reduce background and select the signal

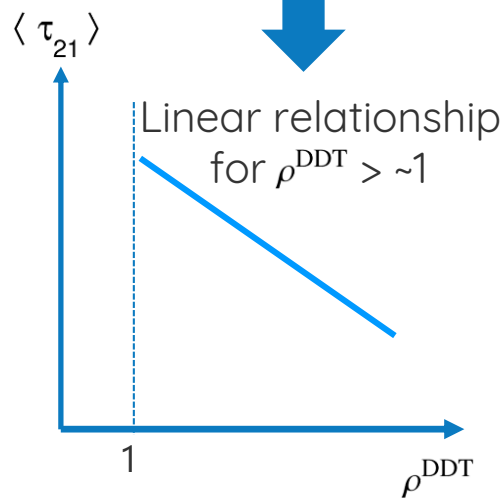
Signal has two particles in the jet vs one for the background

- τ_N is a measure of jet's compatibility with having N subjets
- $\tau_{21} = \tau_2 / \tau_1$ discriminates between jets due to one particle and two

However τ_{21} is correlated with the large radius jet's mass (m_j)

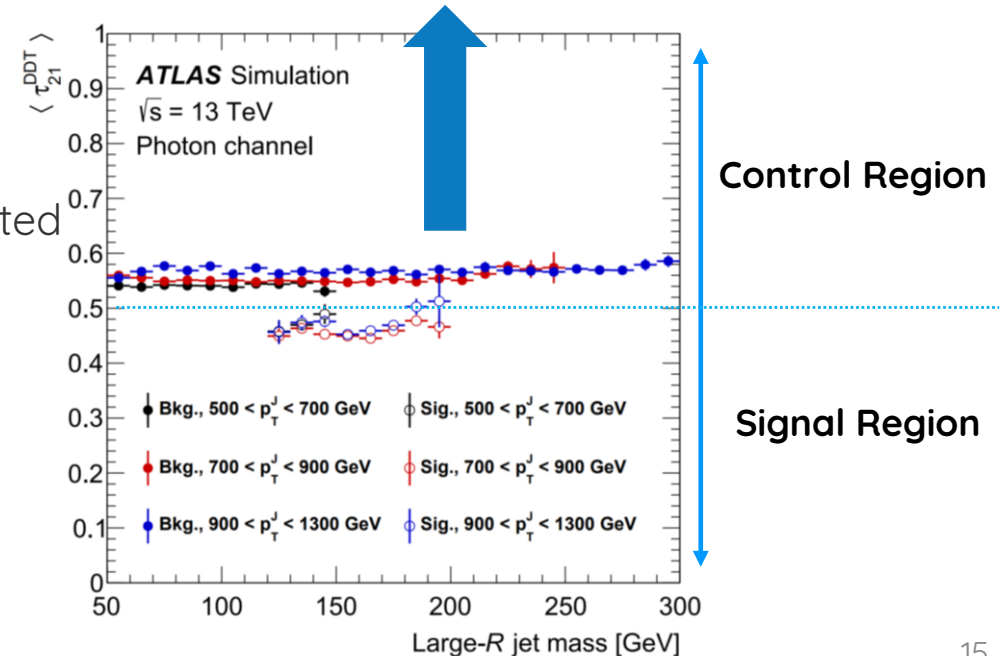
Use designated decorrelated tagger method [JHEP 05 \(2016\) 156](#)

$$\rho^{\text{DDT}} \equiv \log \left(\frac{m_j^2}{p_T^J \times \mu} \right)$$

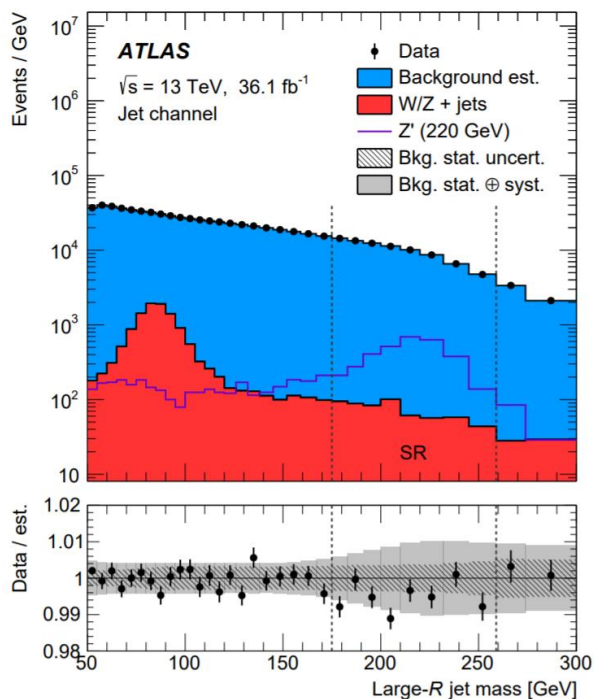


Define τ_{21}^{DDT} a linearly corrected version of τ_{21}

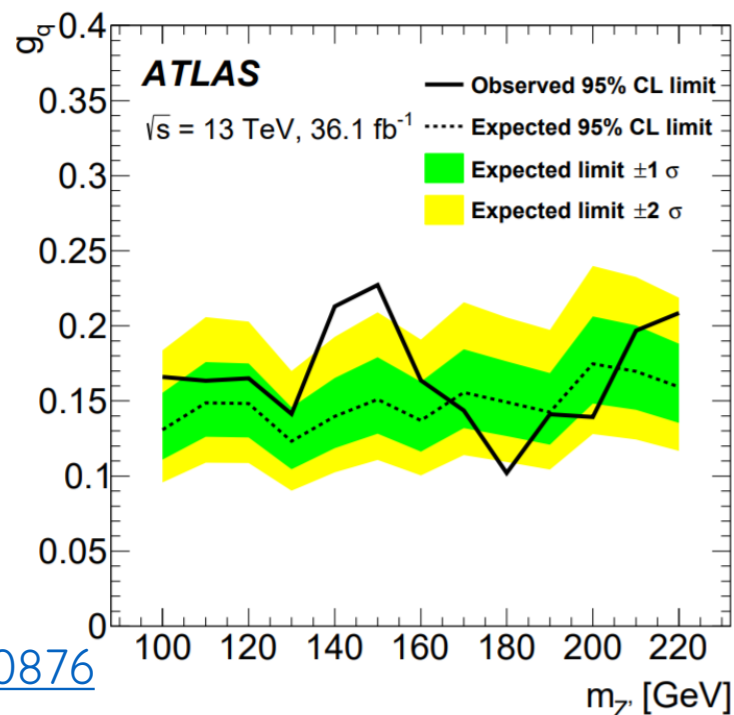
Now independent of jet mass



Signal Region Selections	ISR Photon	ISR Jet
ISR Object pT	> 155 GeV	> 420 GeV
Large Radius Jet pT	> 200 GeV	> 450 GeV
$\Delta\phi(\text{Large Jet, photon/jet})$	> $\pi/2$	
Large jet momentum	> 2 x the jet mass	
τ_{21}^{DDT}	< 0.5	
ρ^{DDT}	> 1.5	



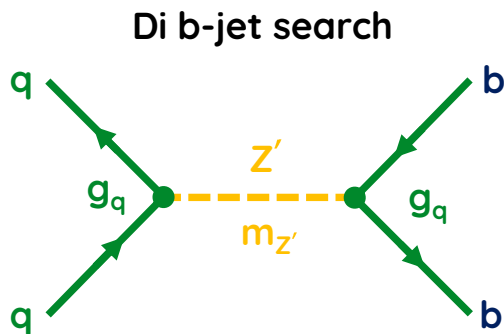
Probes Z' masses down to 100 GeV



[arxiv: 1801.0876](https://arxiv.org/abs/1801.0876)

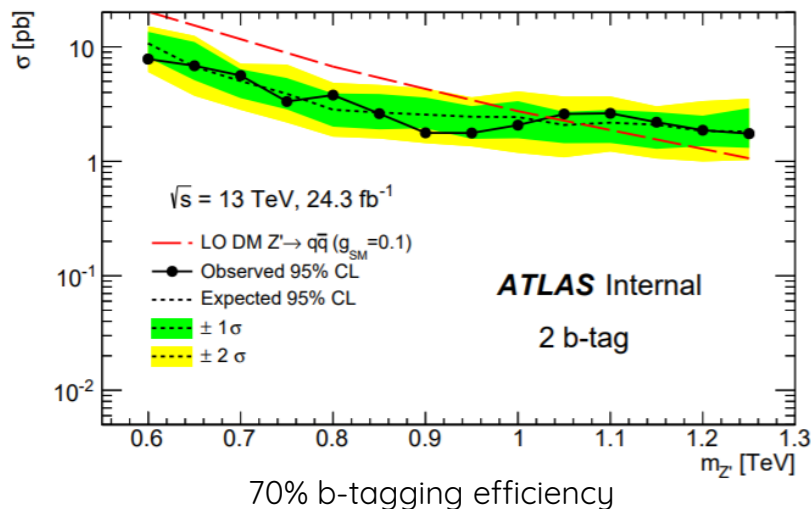
Specific searches for resonances with pairs bottom or top quarks in the final dijet resonance

See Siyuan Sun's talk for details



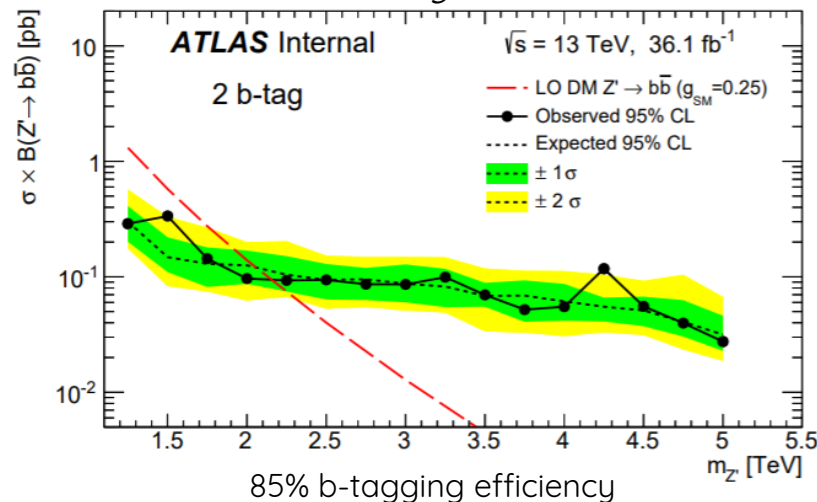
For lower mediator masses

Excludes masses up to 1.02 TeV ($g_q=0.1$)



For higher mediator masses

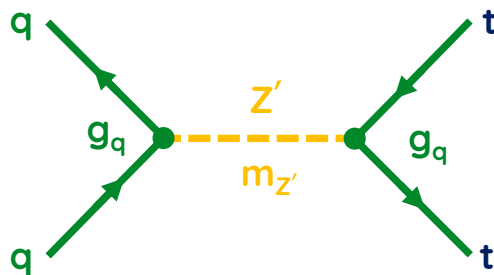
Excludes masses up to 2.1 TeV ($g_q=0.25$)
assumes only $Z' \rightarrow bb$



Specific searches for resonances with pairs bottom or top quarks in the final dijet resonance

See Siyuan Sun's talk for details

Search for top quark pairs



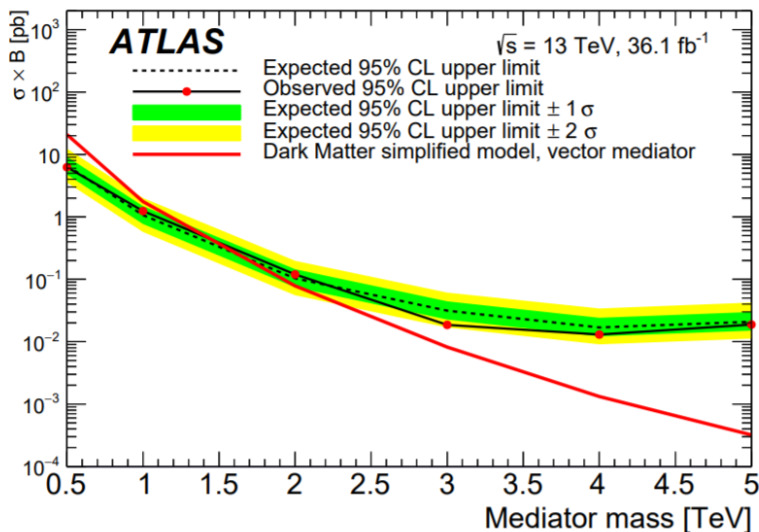
Vector Mediator

$$g_q = 0.25$$

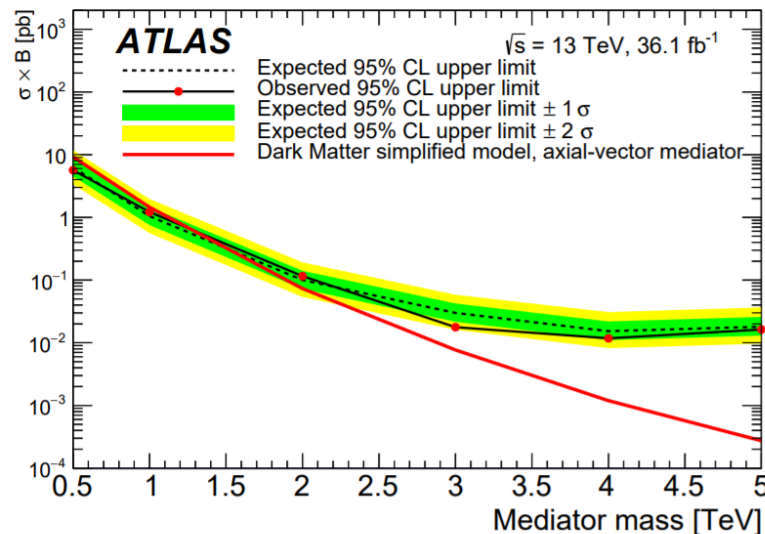
$$g_{DM} = 1$$

$$g_l = 0$$

Axial vector Mediator

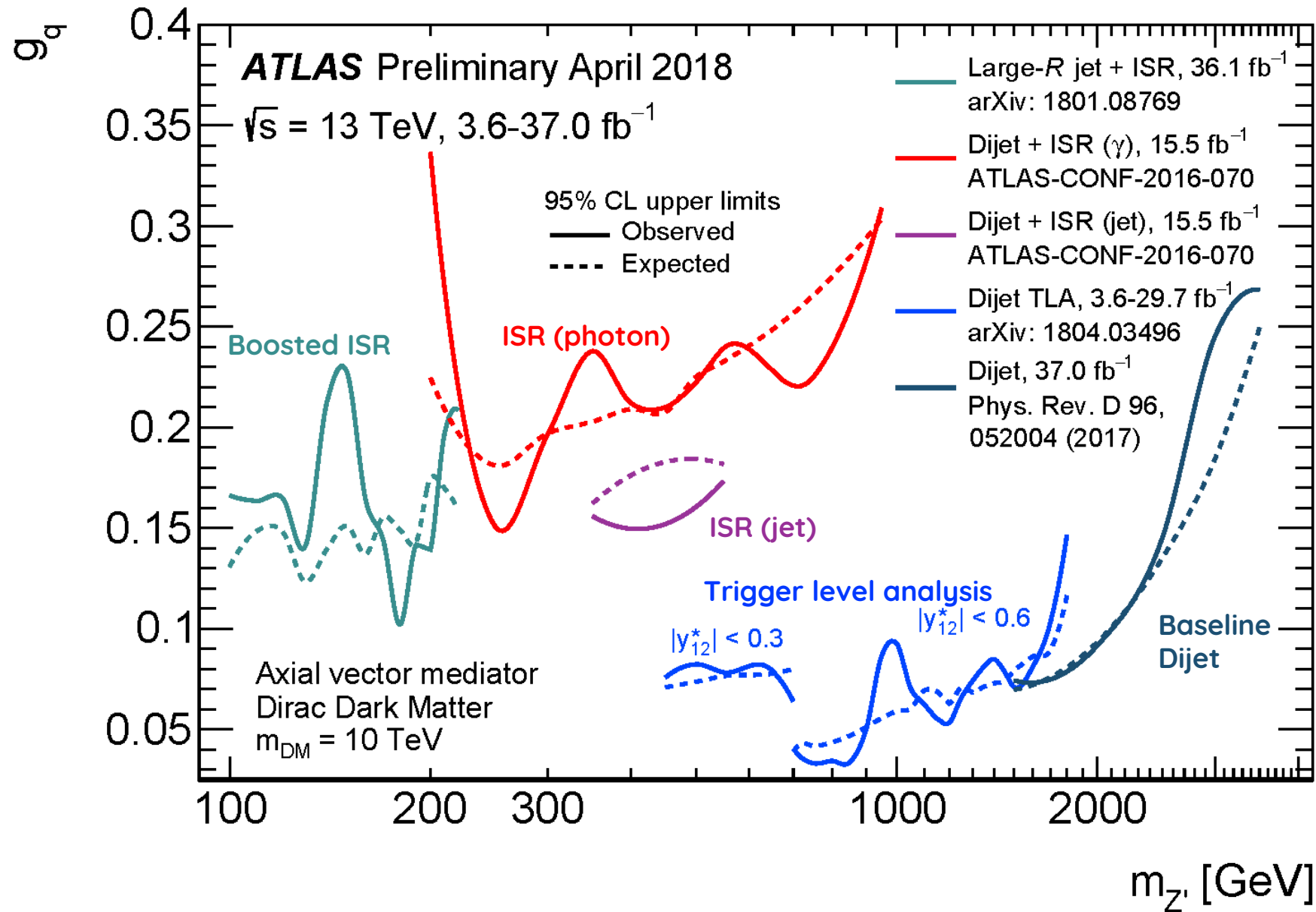


Excludes masses up to 1.4 TeV



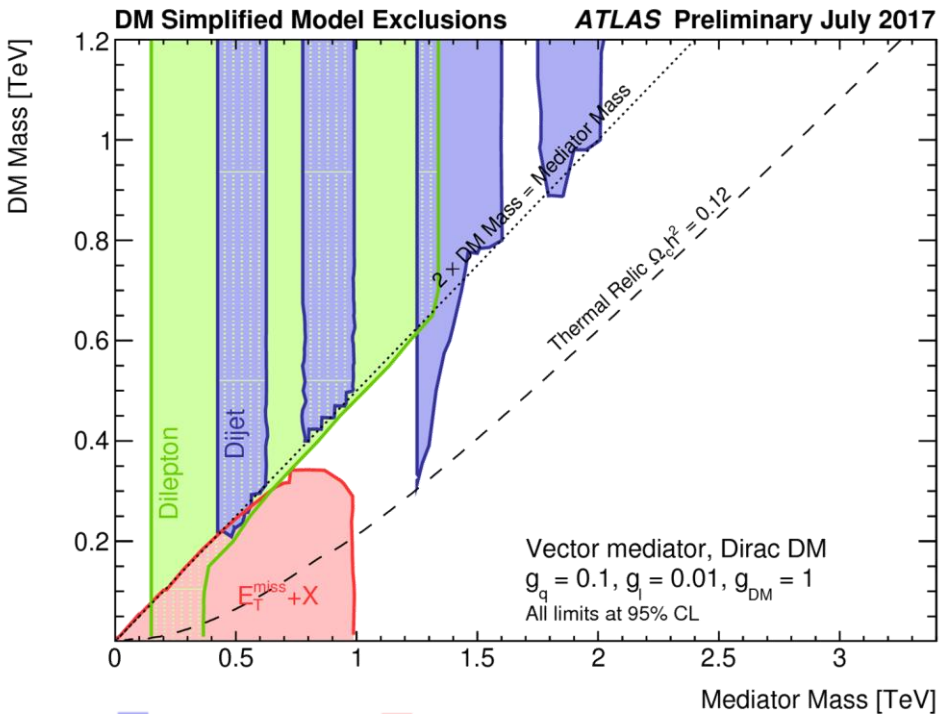
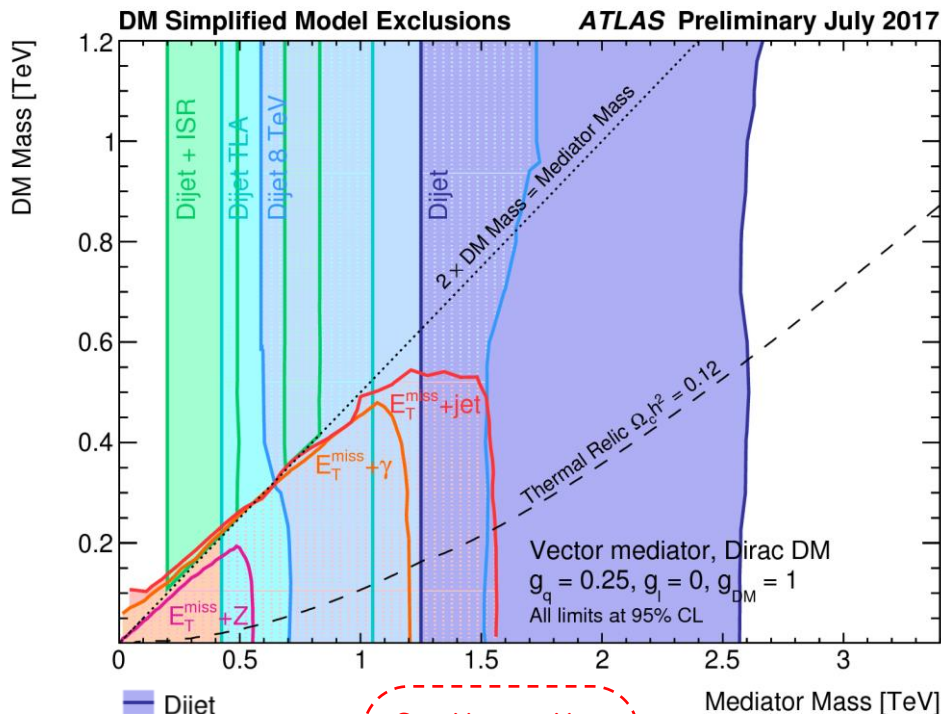
Excludes masses up to 1.2 TeV

[arxiv: 1804.10823](https://arxiv.org/abs/1804.10823)



$$g_q = 0.25 \quad g_{DM} = 1 \quad g_l = 0$$

$$g_q = 0.1 \quad g_{DM} = 1 \quad g_l = 0.01$$



- **Dijet**
 $\sqrt{s} = 13 \text{ TeV}, 37.0 \text{ fb}^{-1}$
arXiv:1703.09127 [hep-ex]
- **Dijet 8 TeV**
 $\sqrt{s} = 8 \text{ TeV}, 20.3 \text{ fb}^{-1}$
Phys. Rev. D. 91 052007 (2015)
- **Dijet TLA**
 $\sqrt{s} = 13 \text{ TeV}, 3.4 \text{ fb}^{-1}$
ATLAS-CONF-2016-030
- **Dijet + ISR**
 $\sqrt{s} = 13 \text{ TeV}, 15.5 \text{ fb}^{-1}$
ATLAS-CONF-2016-070

- See Young-Kee Kim's talk**

 - $E_T^{\text{miss}} + \gamma$
 $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$
Eur. Phys. J. C 77 (2017) 393
 - $E_T^{\text{miss}} + \text{jet}$
 $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$
ATLAS-CONF-2017-060
 - $E_T^{\text{miss}} + Z$
 $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$
ATLAS-CONF-2017-040

- **Dijet**
Dijet 8 TeV $\sqrt{s} = 8 \text{ TeV}, 20.3 \text{ fb}^{-1}$
Phys. Rev. D. 91 052007 (2015)
Dijet $\sqrt{s} = 13 \text{ TeV}, 37.0 \text{ fb}^{-1}$
arXiv:1703.09127 [hep-ex]
Dijet TLA $\sqrt{s} = 13 \text{ TeV}, 3.4 \text{ fb}^{-1}$
ATLAS-CONF-2016-030
Dijet + ISR $\sqrt{s} = 13 \text{ TeV}, 15.5 \text{ fb}^{-1}$
ATLAS-CONF-2016-070

- $E_T^{\text{miss}} + X$
 $E_T^{\text{miss}} + \gamma$ $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$
Eur. Phys. J. C 77 (2017) 393
 $E_T^{\text{miss}} + \text{jet}$ $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$
ATLAS-CONF-2017-060
- **Dilepton**
 $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$
CERN-EP-2017-119

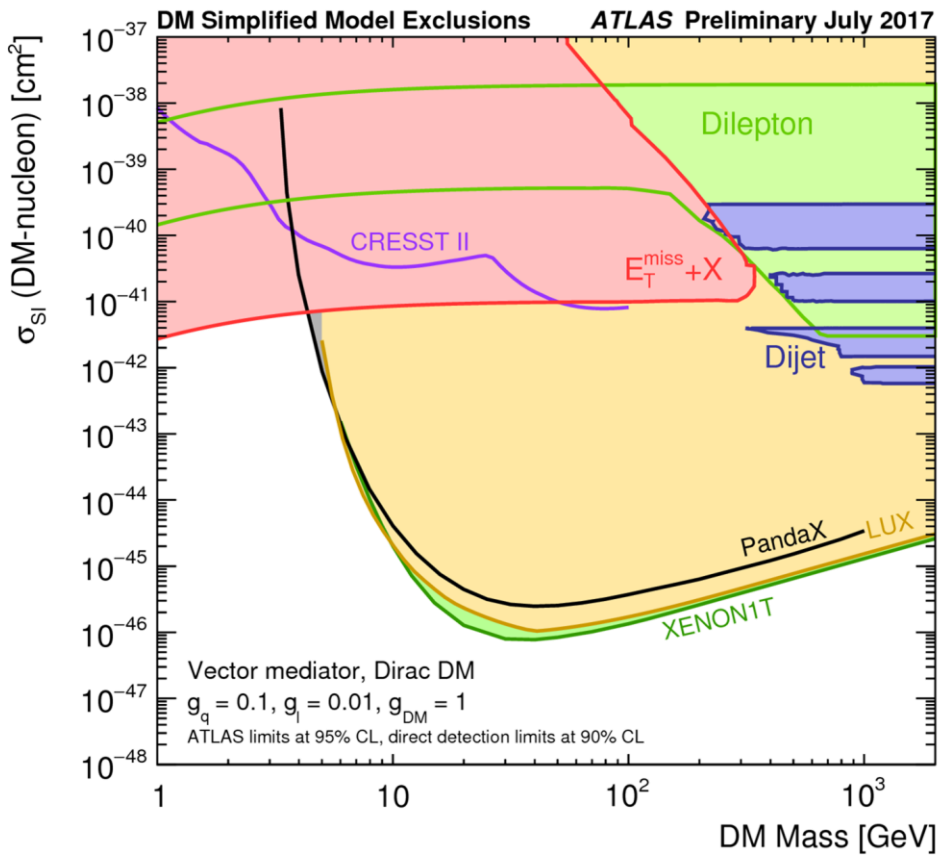
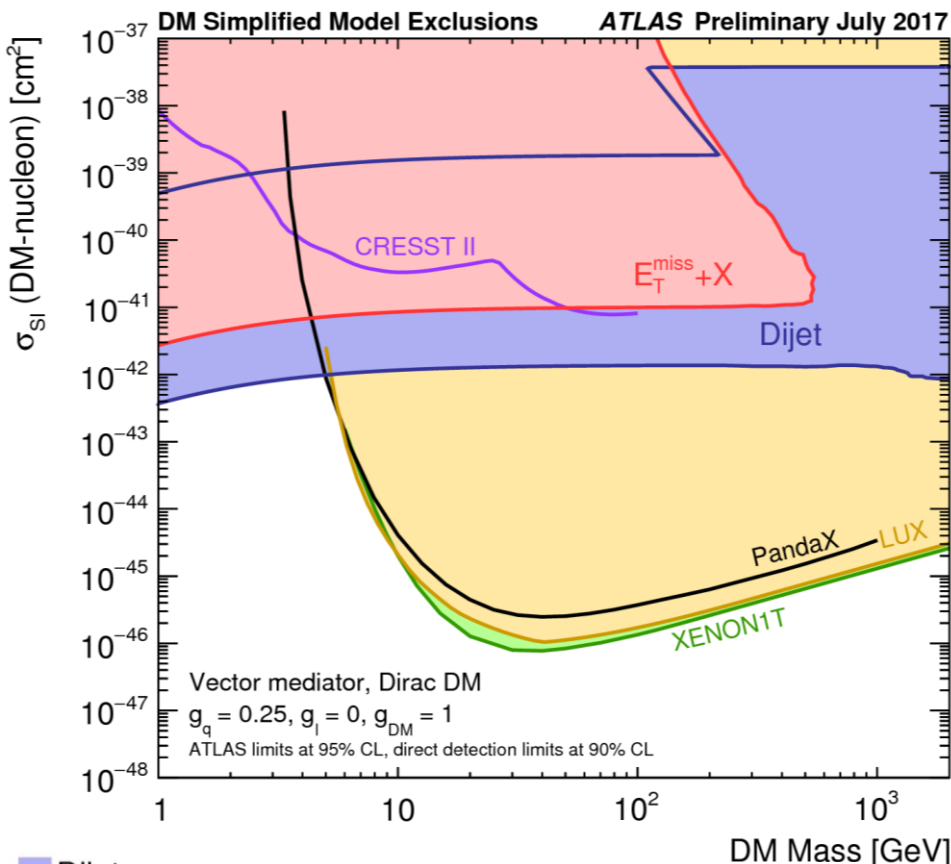
See Young-Kee Kim's talk

See Sébastien Rettie's talk

Not updated for latest TLA and ISR results

$g_q = 0.25$ $g_{DM} = 1$ $g_l = 0$

$g_q = 0.1$ $g_{DM} = 1$ $g_l = 0.01$



- **Dijet**
 Dijet 8 TeV $\sqrt{s} = 8$ TeV, 20.3 fb⁻¹
 Phys. Rev. D. 91 052007 (2015)
 Dijet $\sqrt{s} = 13$ TeV, 37.0 fb⁻¹
 arXiv:1703.09127 [hep-ex]
 Dijet TLA $\sqrt{s} = 13$ TeV, 3.4 fb⁻¹
 ATLAS-CONF-2016-030
 Dijet + ISR $\sqrt{s} = 13$ TeV, 15.5 fb⁻¹
 ATLAS-CONF-2016-070
- **$E_T^{\text{miss}} + X$**
 $E_T^{\text{miss}} + \gamma$ $\sqrt{s} = 13$ TeV, 36.1 fb⁻¹
 Eur. Phys. J. C 77 (2017) 393
 $E_T^{\text{miss}} + \text{jet}$ $\sqrt{s} = 13$ TeV, 36.1 fb⁻¹
 ATLAS-CONF-2017-060
 $E_T^{\text{miss}} + Z$ $\sqrt{s} = 13$ TeV, 36.1 fb⁻¹
 ATLAS-CONF-2017-040
- **CRESST II**
 arXiv:1509.01515v1
- **XENON1T**
 arXiv:1705.06655v2
- **PandaX**
 arXiv:1607.07400
- **LUX**
 arXiv:1608.07648; arXiv:1602.03489

*Not updated
for latest TLA
and ISR
results*

- **Dijet**
 Dijet 8 TeV $\sqrt{s} = 8$ TeV, 20.3 fb⁻¹
 Phys. Rev. D. 91 052007 (2015)
 Dijet $\sqrt{s} = 13$ TeV, 37.0 fb⁻¹
 arXiv:1703.09127 [hep-ex]
 Dijet TLA $\sqrt{s} = 13$ TeV, 3.4 fb⁻¹
 ATLAS-CONF-2016-030
 Dijet + ISR $\sqrt{s} = 13$ TeV, 15.5 fb⁻¹
 ATLAS-CONF-2016-070
- **Dilepton**
 $\sqrt{s} = 13$ TeV, 36.1 fb⁻¹
 CERN-EP-2017-119
- **CRESST II**
 arXiv:1509.01515v1
- **XENON1T**
 arXiv:1705.06655v2
- **PandaX**
 arXiv:1607.07400
- **LUX**
 arXiv:1608.07648; arXiv:1602.03489

ATLAS has an ***extensive program of searches for dark matter mediators*** in run 2

- Complementarity between dedicated searches for mediators and dark matter candidates

Searches for mediators with dijets in run 2 are examining:

- ***Lower couplings*** thanks to a large number of recorded events
- ***Higher masses*** thanks to higher collision energies
- ***Lower masses*** thanks to new experimental techniques
 - ***Trigger level analysis*** extending to ***lower masses and low couplings***
 - ***ISR search*** extends search to ***lower masses***
 - ***Boosted ISR*** further extends search to ***even lower masses***

Searches ***cover all mediator masses between 100 GeV and 3.5 TeV***

Looking forward to the full run 2 results

A 3D visualization of the ATLAS detector, showing the complex structure of the calorimeters and tracking chambers. A central event is shown with a green starburst of tracks radiating from a point, indicating a particle collision. The detector is rendered in shades of blue and yellow, with a dark background.

Additional Slides



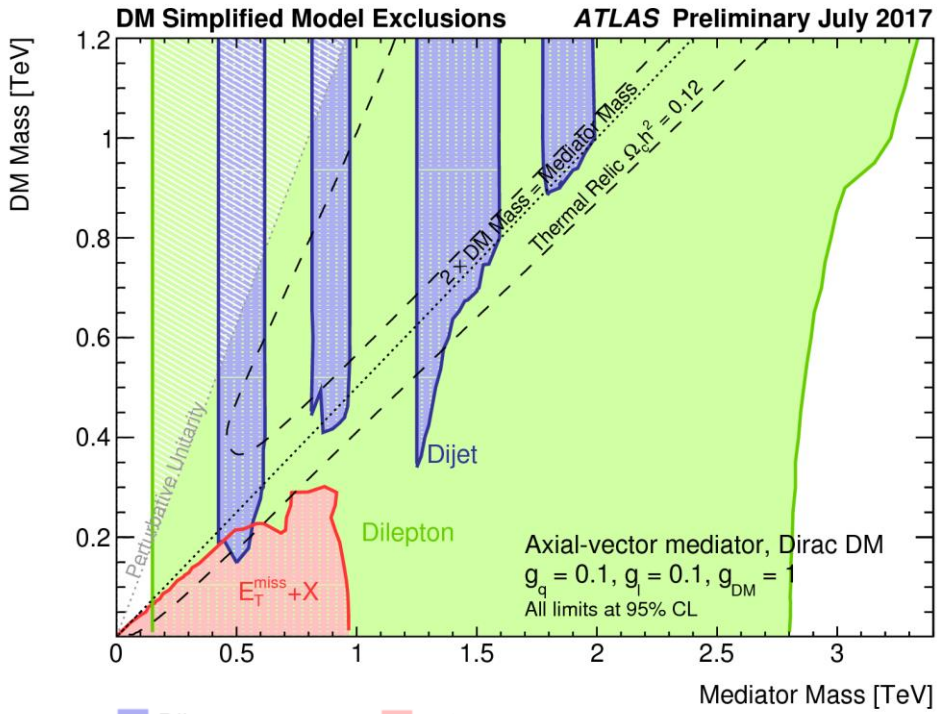
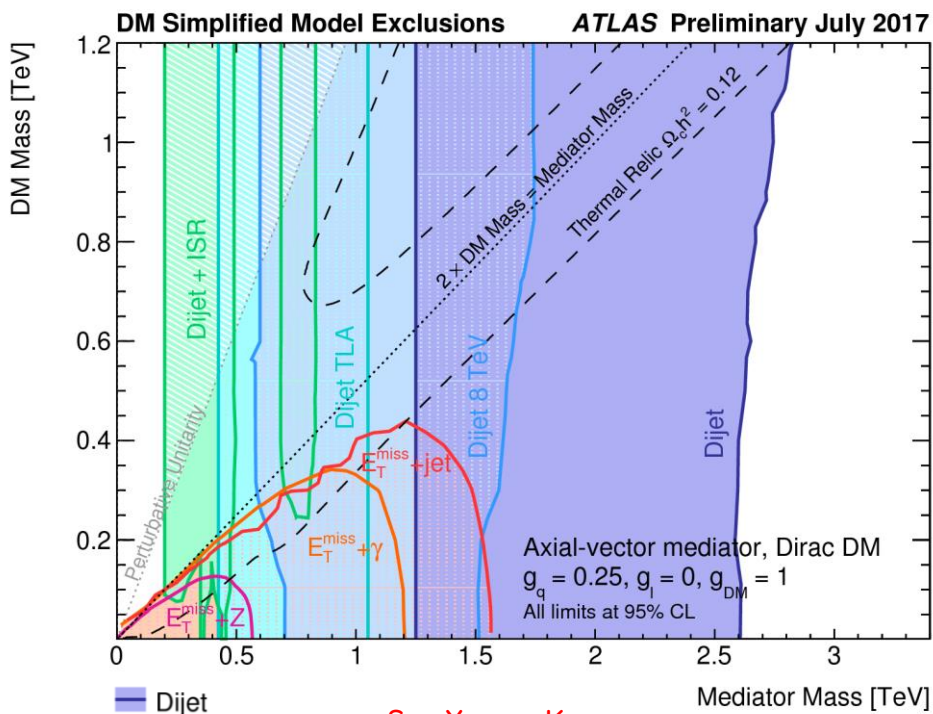
Run: 305777

Event: 4144227629

2016-08-08 08:51:15 CEST

$g_q = 0.25$ $g_{DM} = 1$ $g_l = 0$

$g_q = 0.25$ $g_{DM} = 1$ $g_l = 0.1$



- **Dijet**
 $\sqrt{s} = 13 \text{ TeV}, 37.0 \text{ fb}^{-1}$
 arXiv:1703.09127 [hep-ex]
- **Dijet 8 TeV**
 $\sqrt{s} = 8 \text{ TeV}, 20.3 \text{ fb}^{-1}$
 Phys. Rev. D. 91 052007 (2015)
- **Dijet TLA**
 $\sqrt{s} = 13 \text{ TeV}, 3.4 \text{ fb}^{-1}$
 ATLAS-CONF-2016-030
- **Dijet + ISR**
 $\sqrt{s} = 13 \text{ TeV}, 15.5 \text{ fb}^{-1}$
 ATLAS-CONF-2016-070
- $E_T^{\text{miss}} + \gamma$
 $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$
 Eur. Phys. J. C 77 (2017) 393
- $E_T^{\text{miss}} + \text{jet}$
 $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$
 ATLAS-CONF-2017-060
- $E_T^{\text{miss}} + Z$
 $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$
 ATLAS-CONF-2017-040

- **Dijet**
 Dijet 8 TeV $\sqrt{s} = 8 \text{ TeV}, 20.3 \text{ fb}^{-1}$
 Phys. Rev. D. 91 052007 (2015)
 Dijet $\sqrt{s} = 13 \text{ TeV}, 37.0 \text{ fb}^{-1}$
 arXiv:1703.09127 [hep-ex]
 Dijet TLA $\sqrt{s} = 13 \text{ TeV}, 3.4 \text{ fb}^{-1}$
 ATLAS-CONF-2016-030
 Dijet + ISR $\sqrt{s} = 13 \text{ TeV}, 15.5 \text{ fb}^{-1}$
 ATLAS-CONF-2016-070
- $E_T^{\text{miss}} + X$
 $E_T^{\text{miss}} + \gamma$ $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$
 Eur. Phys. J. C 77 (2017) 393
 $E_T^{\text{miss}} + \text{jet}$ $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$
 ATLAS-CONF-2017-060
- **Dilepton**
 $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$
 CERN-EP-2017-119

See Young-Kee Kim's talk

See Young-Kee Kim's talk

See Sébastien Rettie's talk

Not updated for latest TLA and ISR results

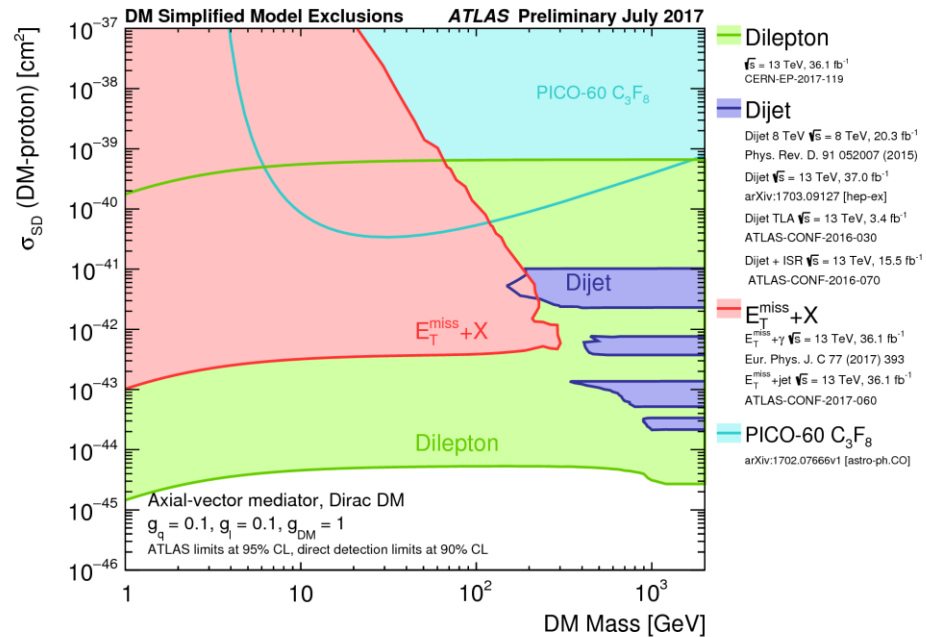
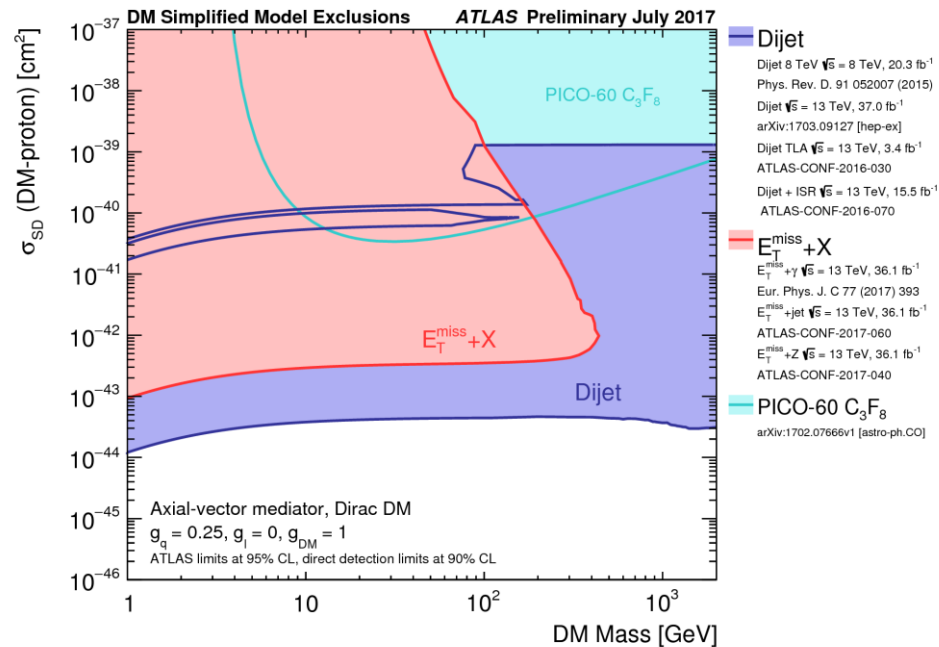
$$\sigma_{\text{SI}} \simeq 6.9 \times 10^{-41} \text{ cm}^2 \cdot \left(\frac{g_q g_{\text{DM}}}{0.25} \right)^2 \left(\frac{1 \text{ TeV}}{M_{\text{med}}} \right)^4 \left(\frac{\mu_{n\chi}}{1 \text{ GeV}} \right)^2 \quad \text{Vector}$$

$$\sigma^{\text{SD}} \simeq 2.4 \times 10^{-42} \text{ cm}^2 \cdot \left(\frac{g_q g_{\text{DM}}}{0.25} \right)^2 \left(\frac{1 \text{ TeV}}{M_{\text{med}}} \right)^4 \left(\frac{\mu_{n\chi}}{1 \text{ GeV}} \right)^2 \quad \text{Axial vector}$$

$$\mu_{n\chi} = m_n m_{\text{DM}} / (m_n + m_{\text{DM}})$$

$$g_q = 0.25 \quad g_{DM} = 1 \quad g_l = 0$$

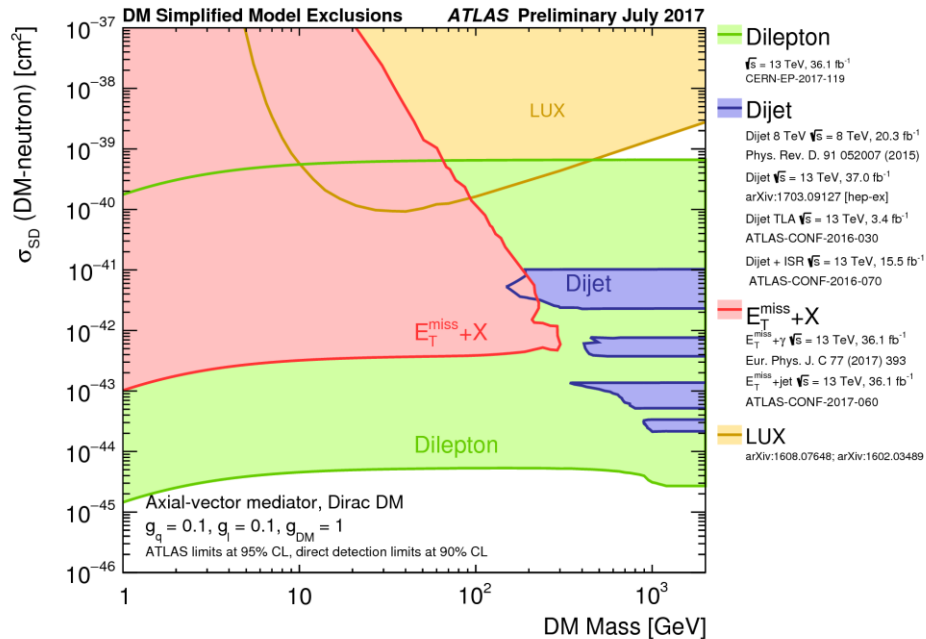
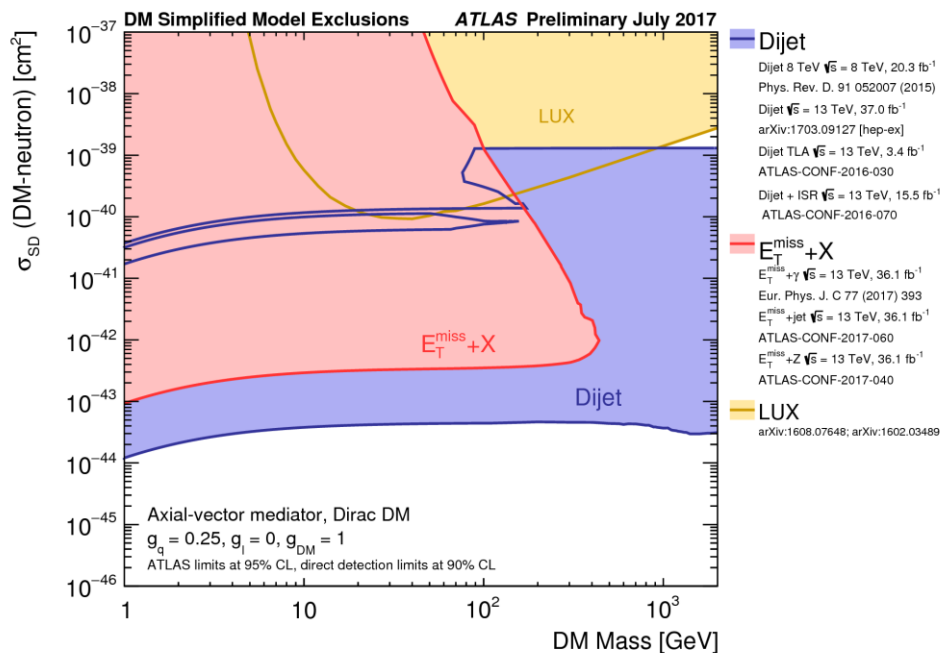
$$g_q = 0.1 \quad g_{DM} = 1 \quad g_l = 0.1$$



Not updated for latest TLA and ISR results

$$g_q = 0.25 \quad g_{DM} = 1 \quad g_l = 0$$

$$g_q = 0.1 \quad g_{DM} = 1 \quad g_l = 0.1$$



Not updated for latest TLA and ISR results

Uses (at trigger level)

- Energy fraction in Electromagnetic Calorimeter
- Energy fraction in Hadronic Calorimeter
- Minimum number of Calorimeter cells which contain 90% of the jet energy

No track based variables available

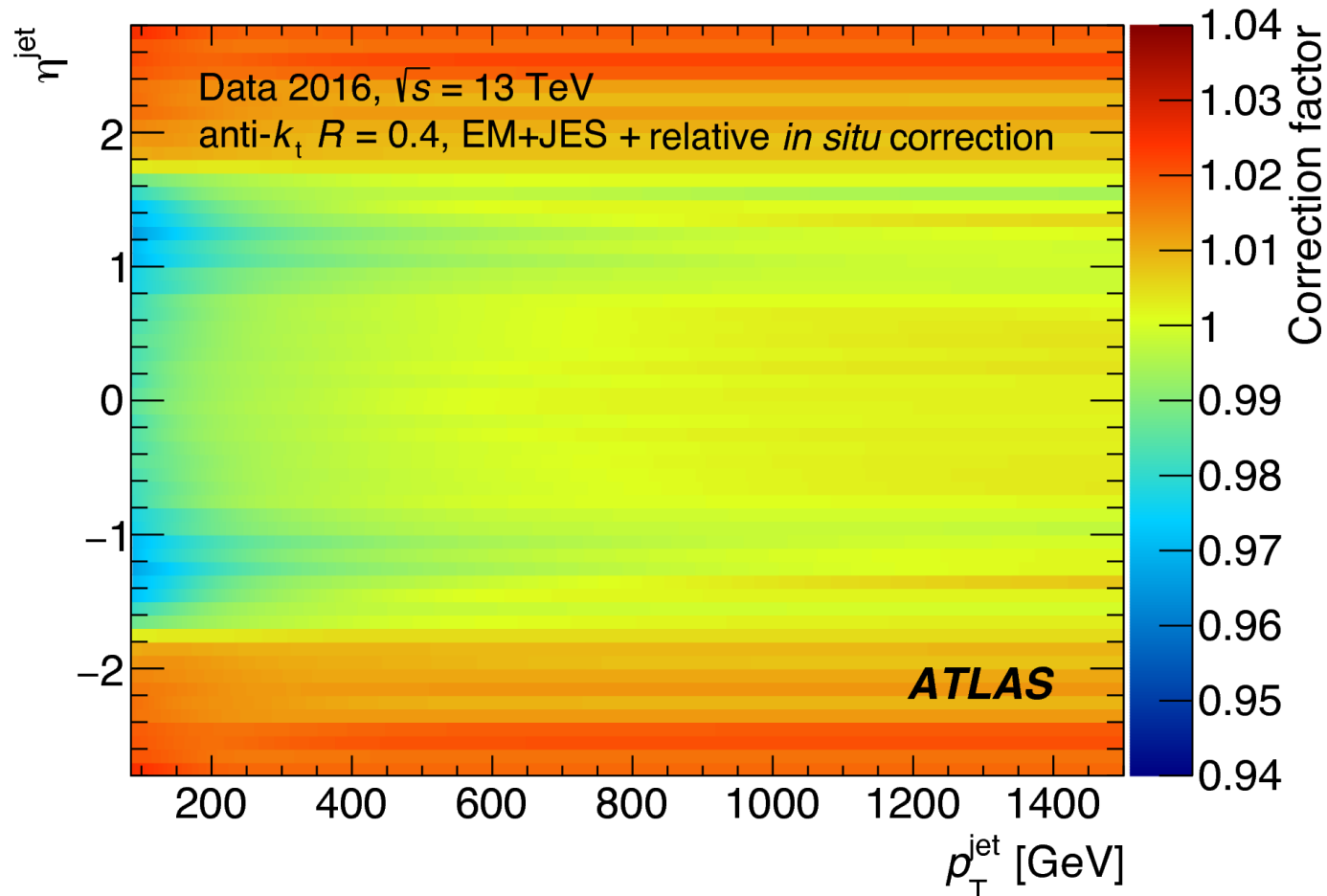
→ Global sequential
calibration

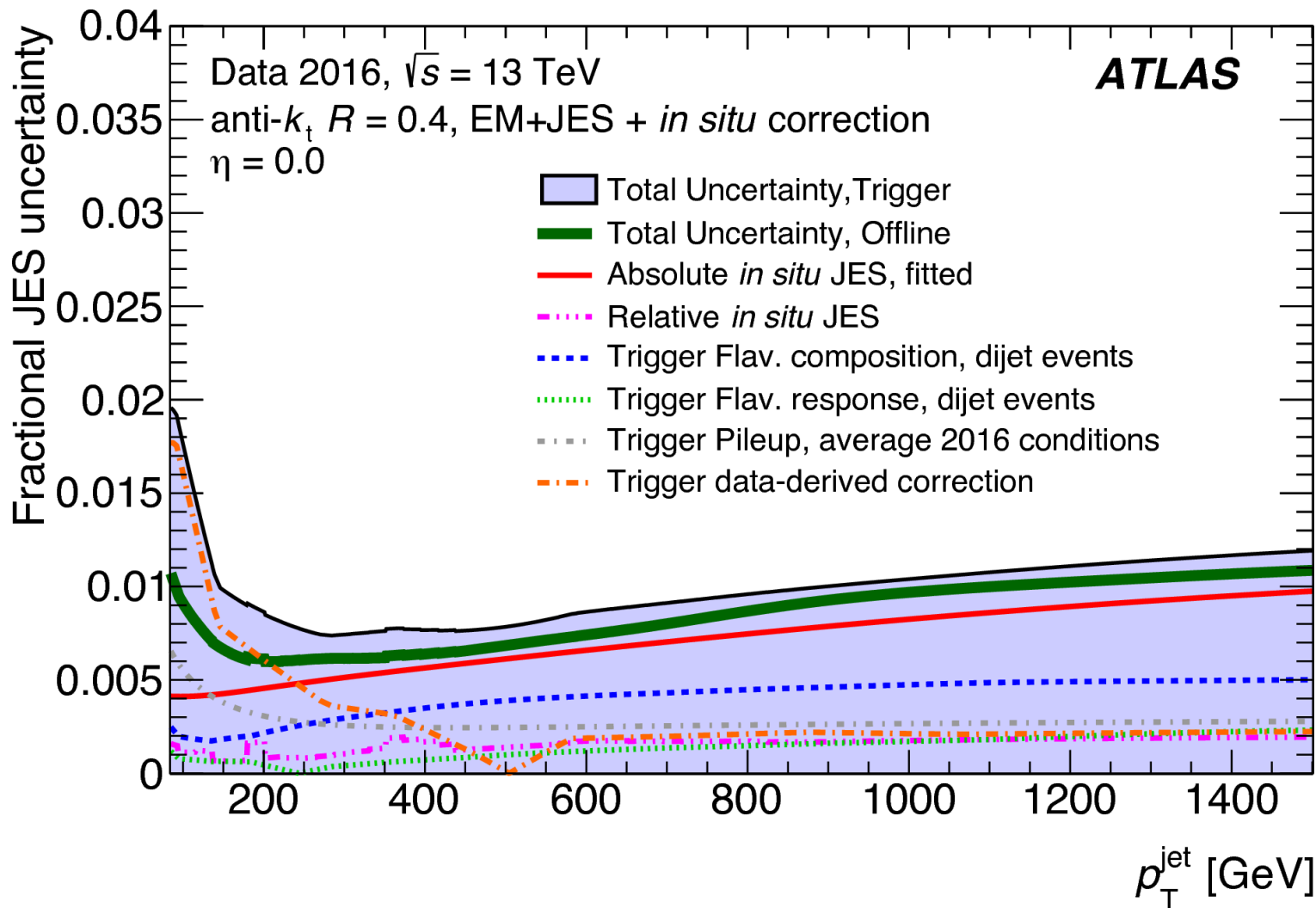
*Reduces flavor
dependence and
energy leakage using
calorimeter variables only*

Correction factor applied to trigger-level jets using the p_T response between trigger-level and offline jets split by η and p_T of the jet

Trigger-to-offline
data-derived correction

Corrects trigger-level jets
to the scale of offline jets,
applied only to data





Fit over m_{jj} performed in a sliding window to estimate background

Three functions are used to fit to the data

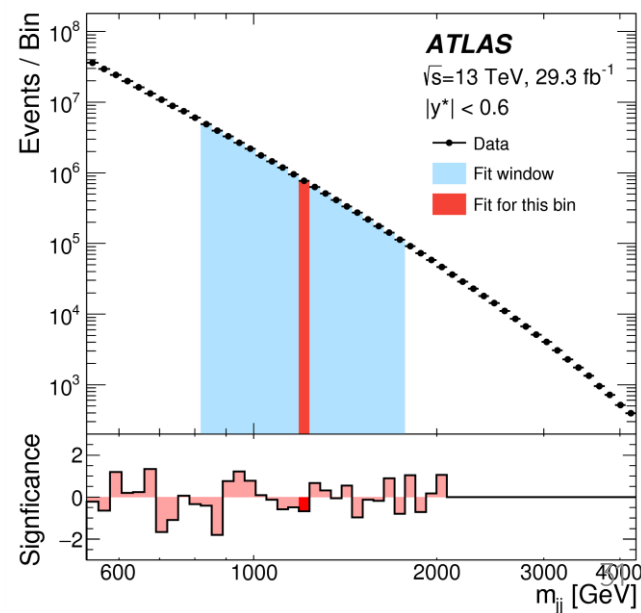
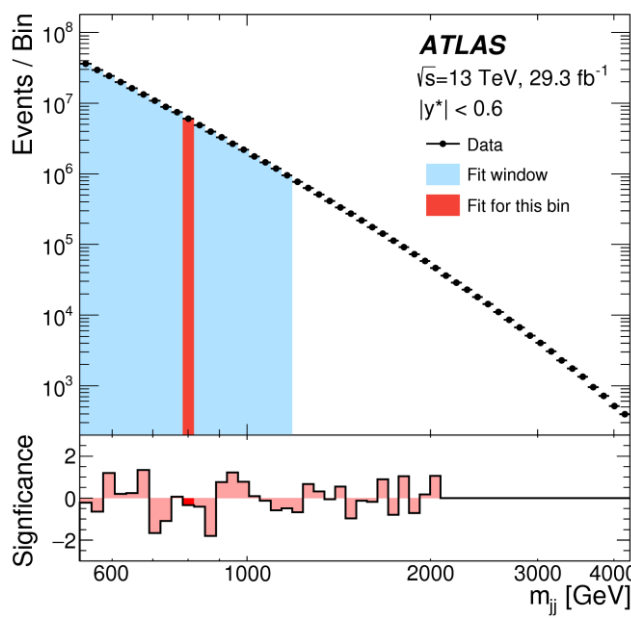
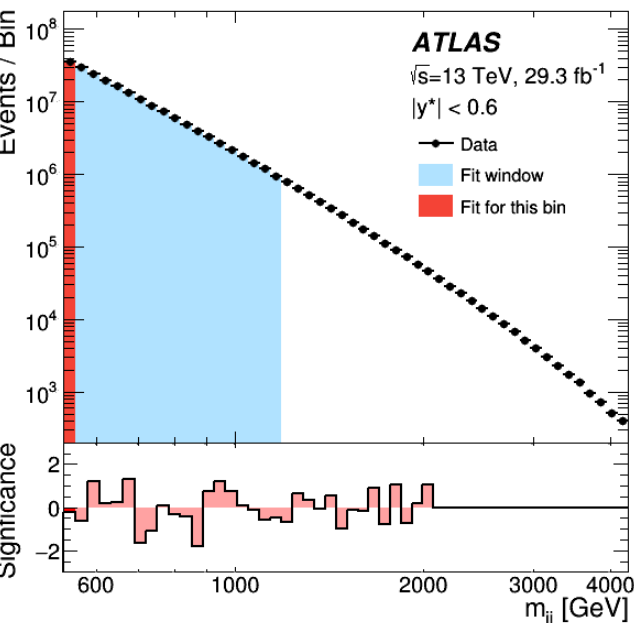
- Final function used has best chi squared over the full range
- Systematic uncertainty uses alternate function

$|y^*| < 0.6$ (0.3) uses first (second) function

$$f(x) = p_1(1-x)^{p_2} x^{p_3+p_4 \ln x + p_5 \ln x^2}$$

$$f(x) = p_1(1-x)^{p_2} x^{p_3+p_4 \ln x} \quad (\text{ISR})$$

$$f(x) = \frac{p_1}{x^{p_2}} e^{-p_3 x - p_4 x^2}$$

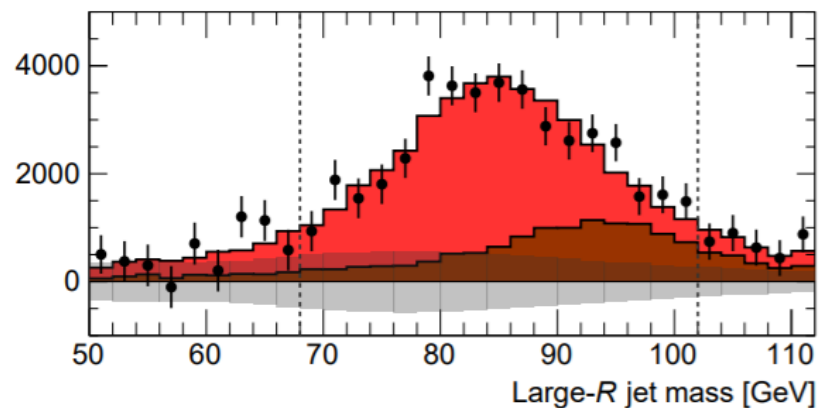
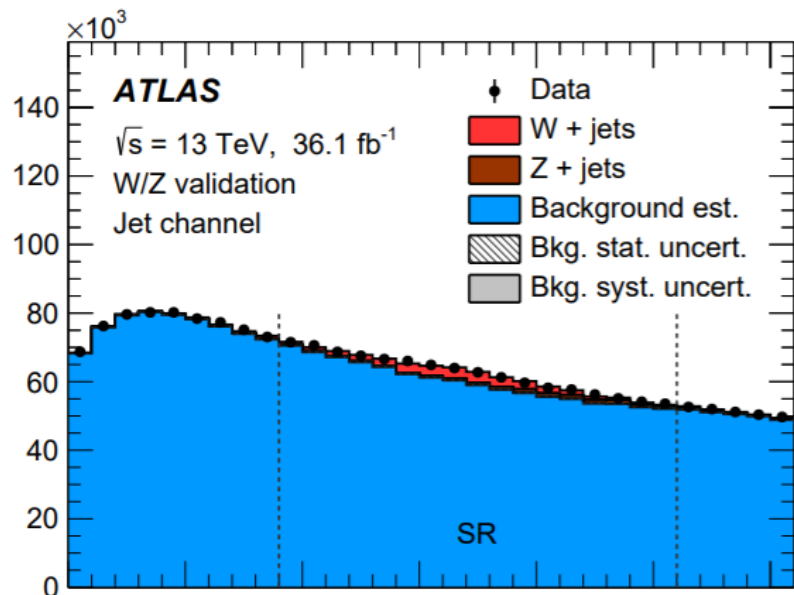


Control region inverts τ_{21}^{DDT} selection

- Background estimated in SR using transfer factor
- Method Validated on W/Z peak

Largest (systematic) uncertainty is from the transfer factor

- 90% for $M_{Z'}$ of 160 GeV



BG = TF * (Number in CR - Expected number from production with an associated vector boson)

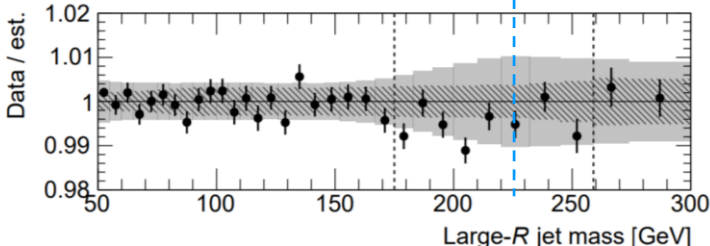
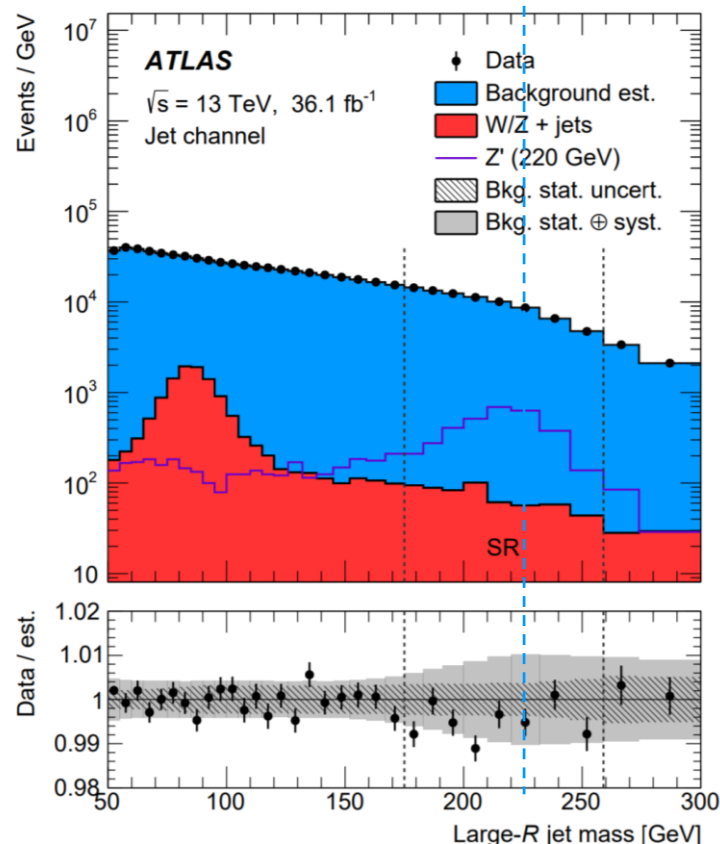
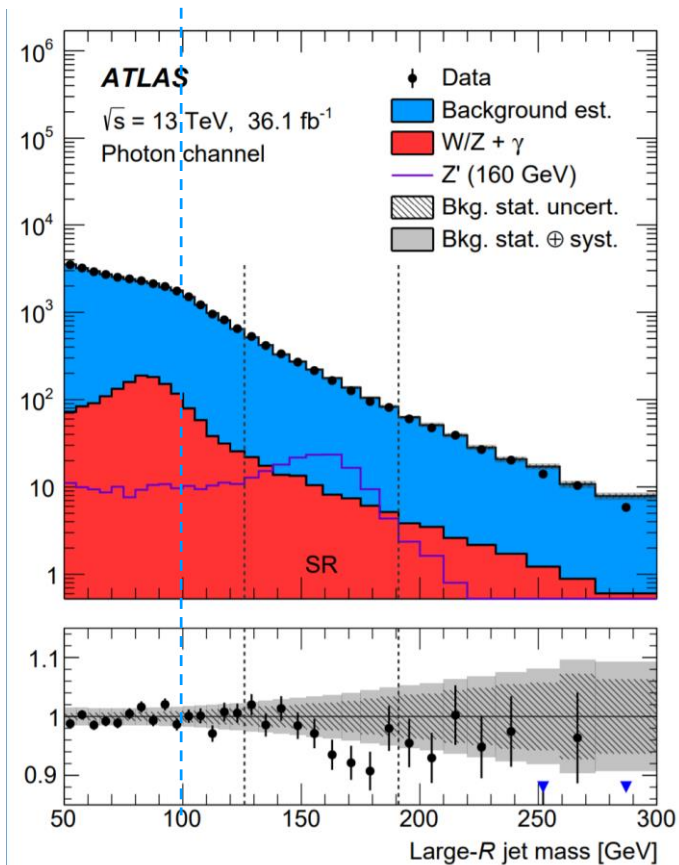
Transfer factor = expected ratio of events which pass τ_{21}^{DDT} to those which fail

- Data measured away from Z' mass under consideration (20% buffer)
- Parameterised by $\log(\text{Large radius jet } p_T / \mu)$ and ρ^{DDT}

Uncertainty source	$\Delta\mu/\mu$ [%]	
	$m_{Z'} = 160$ GeV	$m_{Z'} = 220$ GeV
Transfer factor	90	88
Large- R jet	25	17
Total systematic uncertainty	93	91
Statistical uncertainty	10	11

Boosted requirement of large jet momentum $> 2 \times$ the jet mass results in altered slope

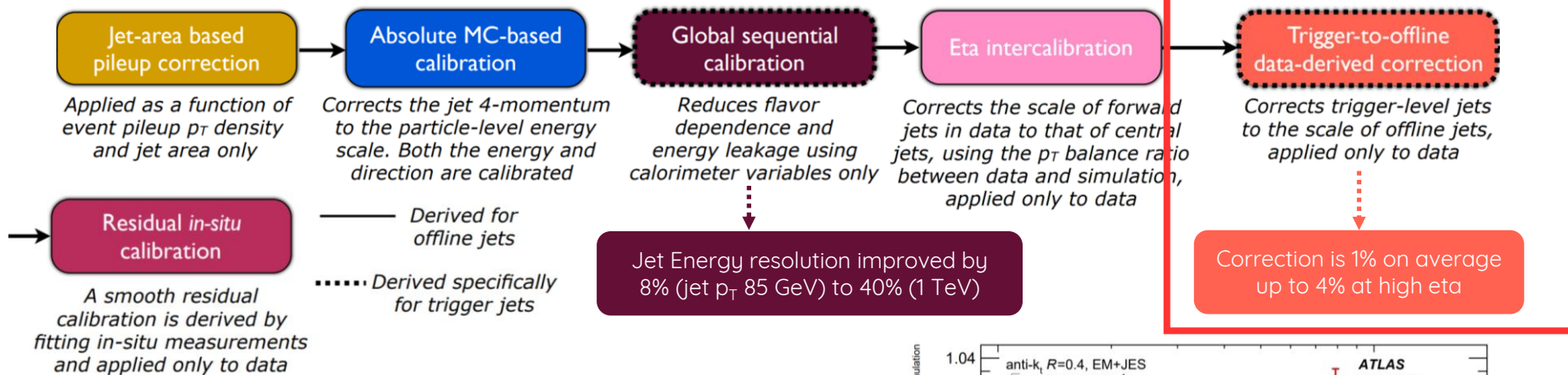
Large Radius Jet p_T	$> 200 \text{ GeV}$	$> 450 \text{ GeV}$
------------------------	---------------------	---------------------



$$\mathcal{L}_{\text{vector}} = -g_{\text{DM}} Z'_{\mu} \bar{\chi} \gamma^{\mu} \chi - g_q \sum_{q=u,d,s,c,b,t} Z'_{\mu} \bar{q} \gamma^{\mu} q,$$

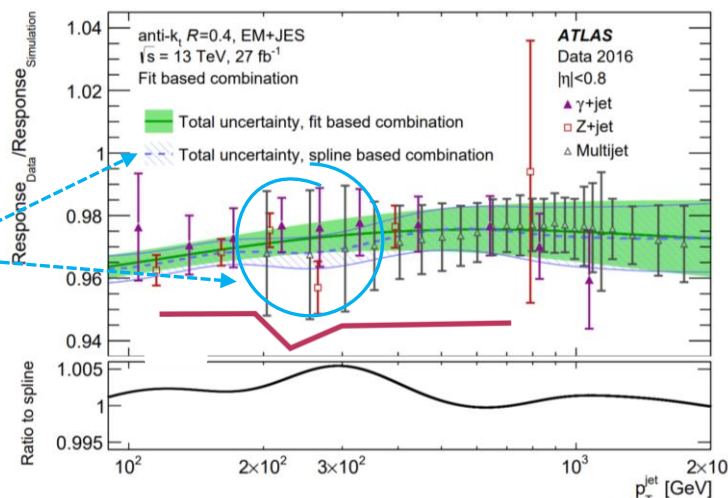
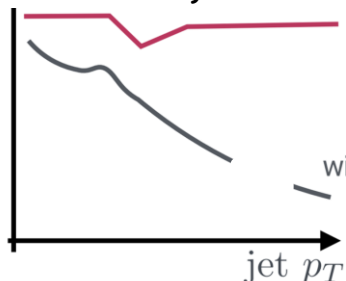
$$\mathcal{L}_{\text{axial-vector}} = -g_{\text{DM}} Z'_{\mu} \bar{\chi} \gamma^{\mu} \gamma_5 \chi - g_q \sum_{q=u,d,s,c,b,t} Z'_{\mu} \bar{q} \gamma^{\mu} \gamma_5 q.$$

Jet energies are calibrated similarly to offline jets



Calibrate using momentum balance of offline jets against well calibrated objects

- Needs to be smooth as can introduce bumps
- Use polynomial fit in $\log(p_T)$ rather than spline



After calibration energy of trigger and offline jets agree within 0.05%

Jet energy scale uncertainty is 1-2%, similar to that of offline jets