

What the Milky Way's Dwarfs have to tell us about the Galactic Center extended excess

Ryan Keeley

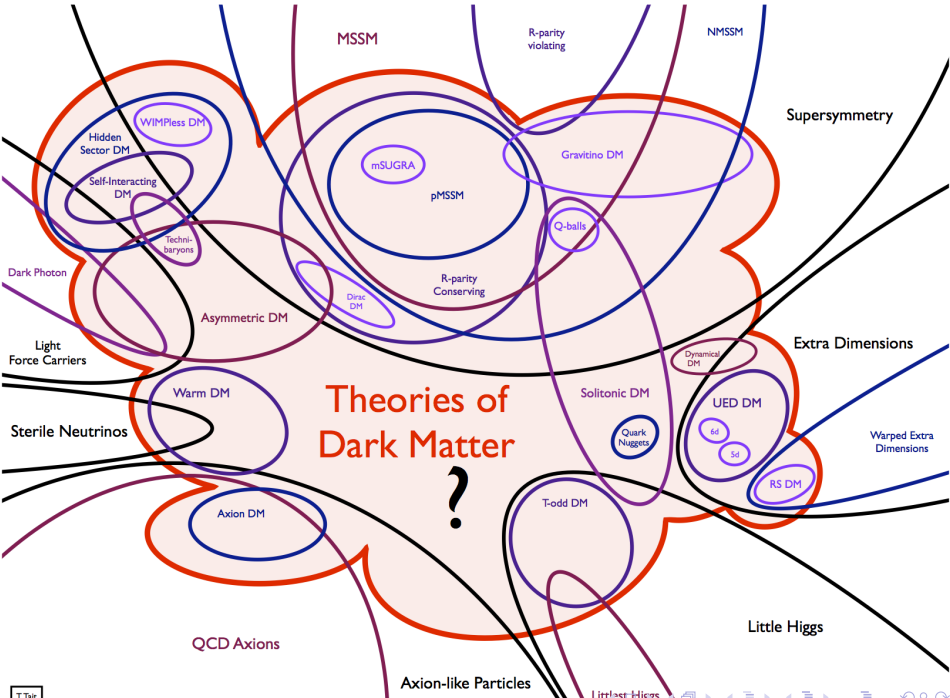
CIPANP

June 1, 2018

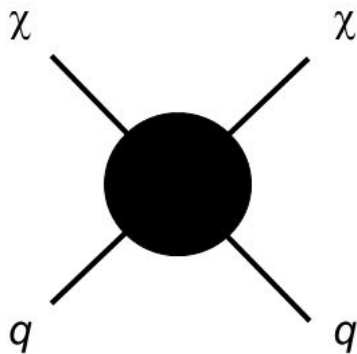
ArXiv: 1710.03215

Collaborators: Kevork Abazajian, Anna Kwa, Nick Rodd, Ben Safdi

Theories of Dark Matter



Efficient annihilation now
(Indirect detection)



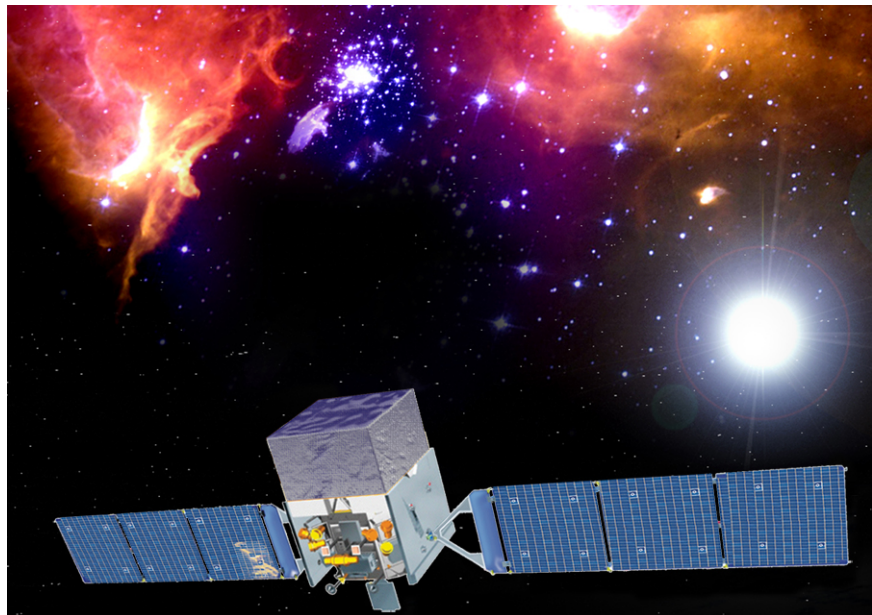
Efficient scattering now
(Direct detection)



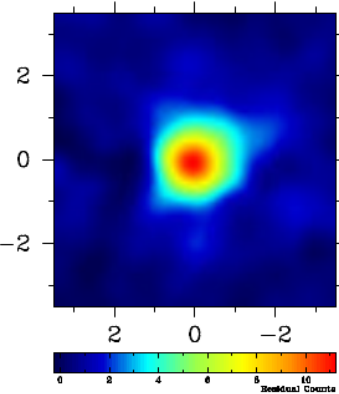
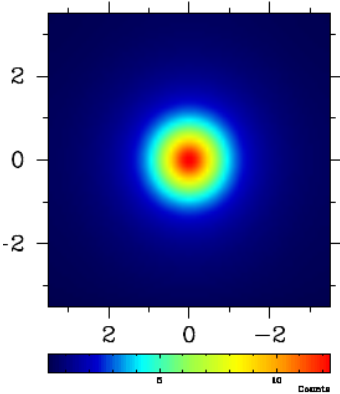
Efficient production now
(Particle colliders)



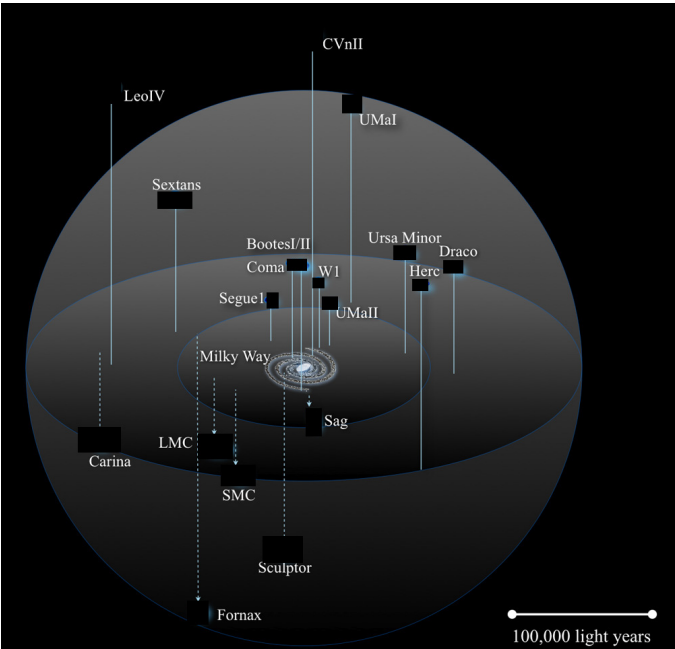
Fermi Gamma Ray Space Telescope



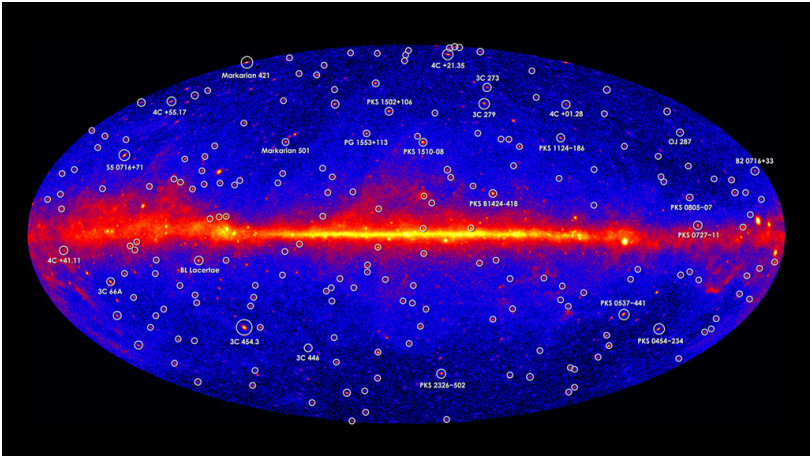
Dark Matter?



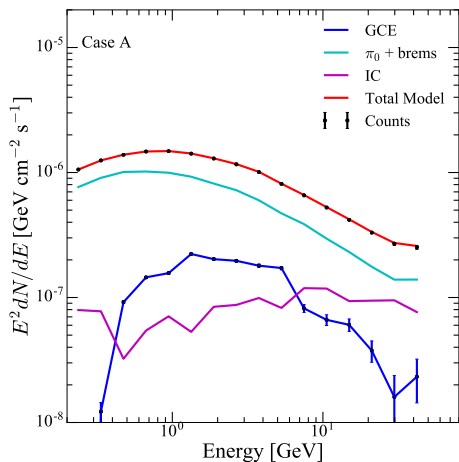
Dim Dwarf Galaxies



Complicated Gamma Ray Sky



An Accounting of the Fermi Gamma Rays

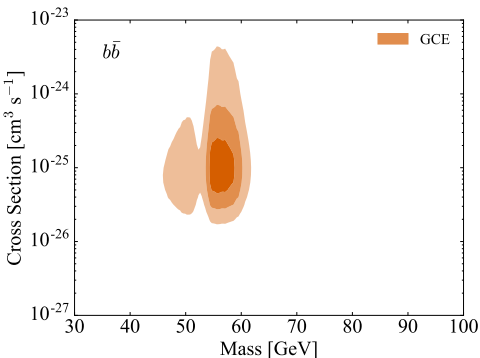


Templates for:

- ▶ Bremsstrahlung
- ▶ π_0
- ▶ inverse Compton

In each energy bin the flux of each template is varied to minimize Poisson likelihood

Flux from Dark Matter Annihilation

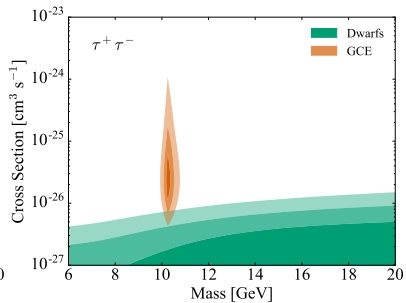
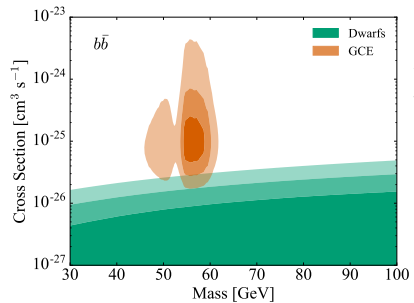


$$\frac{d\Phi}{dE} = \frac{1}{4\pi} \frac{J}{m_\chi^2} \frac{\langle\sigma v\rangle}{2} \frac{dN}{dE} \quad (1)$$

$$J = \int d\Omega \int dz \rho^2(z, \Omega) \quad (2)$$

$$\rho(r) = \frac{\rho_\odot}{\left(\frac{r}{R_\odot}\right)^\gamma \left(\frac{1+r/R_s}{1+R_\odot/R_s}\right)^{3-\gamma}} \quad (3)$$

Tension in Dark Matter Interpretation



Evidence Ratios

$$ER = \frac{p(D_1)p(D_2)}{p(D_1, D_2)} = \frac{\int d\theta_1 p(D_1|\theta_1)p(\theta_1) \int d\theta_2 p(D_2|\theta_2)p(\theta_2)}{\int p(D_1, D_2|\theta)p(\theta)} \quad (4)$$

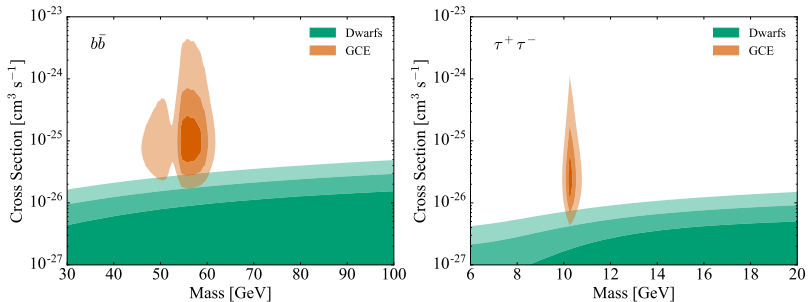
Bayes factor where the model for the numerator is the same as the one for the denominator but has an additional copy of the parameter space

$ER > 1 \rightarrow$ The data prefers to be described by two sets of parameters

$ER < 1 \rightarrow$ The data prefers to be described by one set of parameters

$ER = 1 \rightarrow$ No mutual information between the data sets

DM Tension



Model	
DM: $b\bar{b}$	3600
DM: $\tau^+\tau^-$	2300

Beyond DM Explanations

Log-parabola

$$\frac{dN}{dE} = N_0 \left(\frac{E}{E_s} \right)^{-\alpha - \beta \log(E/E_s)} \quad (5)$$

Exponential cutoff

$$\frac{dN}{dE} = N_0 \left(\frac{E}{E_s} \right)^\gamma e^{-E/E_c}, \quad (6)$$

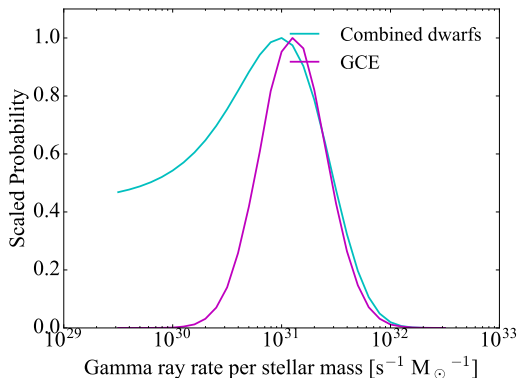
Flux Scalings

- ▶ Break any degeneracies between the normalization parameter and the spectral parameters
- ▶ Answer ‘How well can these models fit the GCE?’ and ‘How well do these models address the difference in flux?’ independently
- ▶ N_0 is set such that the integral of dN/dE over the energy is one.
- ▶ Allow the flux to scale as the stellar mass of the system.

$$\frac{d\Phi}{dE} = \frac{\dot{N}}{4\pi R^2} \frac{M_*}{M_0} \frac{dN}{dE}, \quad (7)$$

- ▶ \dot{N}/M_0 , analogous to cross section
- ▶ M_*/R^2 , analogous to the J-factor

Evidence Ratios: Astrophysics



Model	
Log-Parabola	0.69
Exponential Cutoff	0.73

- ▶ Both the GCE and the slight excesses in the Dwarfs pick out a scale of $\dot{N}/M_0^2 = 10^{31 \pm 1} \text{ sec}^{-1} M_{\odot}^{-1}$
- ▶ This explains why the ER is less than unity

Representative SIDM model

$$\frac{dn_\gamma}{dEdt} = \sigma_T c n_e n_{\text{ISRF}} \frac{dN_\gamma}{dE}, \quad (8)$$

GCE is produced by an IC process where high energy electrons from DM annihilation scatter off the interstellar radiation field

Representative SIDM model

$$\frac{dn_\gamma}{dEdt} = \sigma_T c n_e n_{\text{ISRF}} \frac{dN_\gamma}{dE}, \quad (9)$$

$$\frac{dN_\gamma}{dE} = \int dE_e dE_{\text{ISRF}} p(E_\gamma | E_e, E_{\text{ISRF}}) p(E_e) p(E_{\text{ISRF}}). \quad (10)$$

The energy distribution of the gamma rays is the convolution of the electron energy distribution and the energy distribution of the ISRF, with the IC process.

Representative SIDM model

$$\frac{dn_\gamma}{dEdt} = \sigma_T c n_e n_{\text{ISRF}} \frac{dN_\gamma}{dE}, \quad (11)$$

$$\frac{dN_\gamma}{dE} = \int dE_e dE_{\text{ISRF}} p(E_\gamma | E_e, E_{\text{ISRF}}) p(E_e) p(E_{\text{ISRF}}). \quad (12)$$

$$\frac{d\Phi_\gamma}{dE} = \int dV' \frac{1}{4\pi(\vec{R} - \vec{R}')^2} \frac{dn_\gamma}{dEdt}(\vec{R}'). \quad (13)$$

Deriving how the flux scales...

Representative SIDM model

$$\frac{dn_\gamma}{dEdt} = \sigma_T c n_e n_{\text{ISRF}} \frac{dN_\gamma}{dE}, \quad (14)$$

$$\frac{dN_\gamma}{dE} = \int dE_e dE_{\text{ISRF}} p(E_\gamma | E_e, E_{\text{ISRF}}) p(E_e) p(E_{\text{ISRF}}). \quad (15)$$

$$\frac{d\Phi_\gamma}{dE} = \int dV' \frac{1}{4\pi(\vec{R} - \vec{R}')^2} \frac{dn_\gamma}{dEdt}(\vec{R}'). \quad (16)$$

$$\frac{d\Phi_\gamma}{dE} \propto \int d\Omega dz \frac{1}{4\pi} n_{\text{ISRF}} n_\chi^2 \frac{dN}{dE}. \quad (17)$$

...shows that the flux scales as a three-body process

Representative SIDM model

$$\frac{dn_\gamma}{dEdt} = \sigma_T c n_e n_{\text{ISRF}} \frac{dN_\gamma}{dE}, \quad (18)$$

$$\frac{dN_\gamma}{dE} = \int dE_e dE_{\text{ISRF}} p(E_\gamma | E_e, E_{\text{ISRF}}) p(E_e) p(E_{\text{ISRF}}). \quad (19)$$

$$\frac{d\Phi_\gamma}{dE} = \int dV' \frac{1}{4\pi(\vec{R} - \vec{R}')^2} \frac{dn_\gamma}{dEdt}(\vec{R}'). \quad (20)$$

$$\frac{d\Phi_\gamma}{dE} \propto \int d\Omega dz \frac{1}{4\pi} n_{\text{ISRF}} n_\chi^2 \frac{dN}{dE}. \quad (21)$$

$$\frac{d\Phi_\gamma}{dE} \propto \frac{J}{m_\chi^2} \frac{M_*}{M_{*,\text{GC}}}. \quad (22)$$

Putting this into a useful form, take the flux to scale as the J-factor times the stellar mass of the system.

Evidence Ratios

Model	
DM: $b\bar{b}$	3600
DM: $\tau^+\tau^-$	2300
Log-Parabola	0.69
Exponential Cutoff	0.73
SIDM	1.1

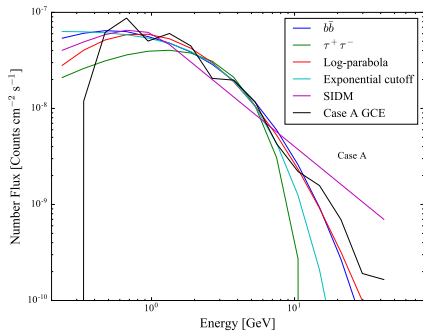
Bayes Factors

$$K_{12} = \frac{p(M_1|D)p(M_2)}{p(M_2|D)p(M_1)} = \frac{p(D|M_1)}{p(D|M_2)}. \quad (23)$$

Bayes factors for the considered models, relative to the $b\bar{b}$ model, for each of the different background cases. Values larger than one indicate the data prefer that model over $b\bar{b}$

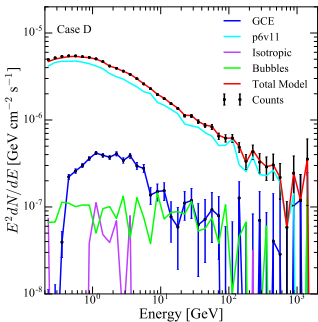
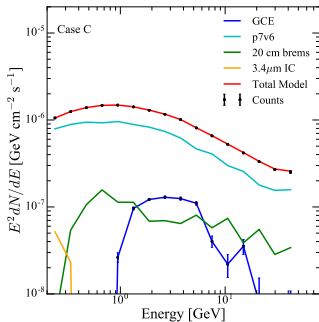
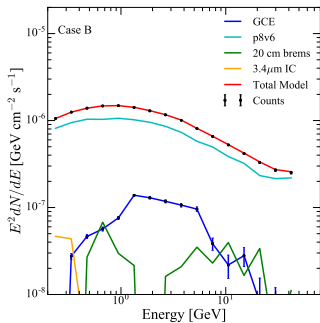
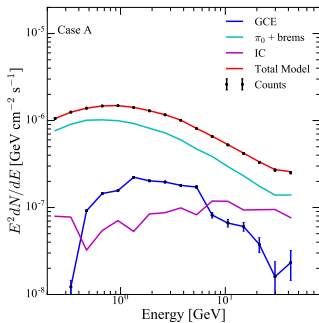
Model	
DM: $\tau^+\tau^-$	4×10^{-24}
Log-Parabola	3×10^{12}
Exponential Cutoff	8×10^7
SIDM	5×10^{-20}

Best Fits

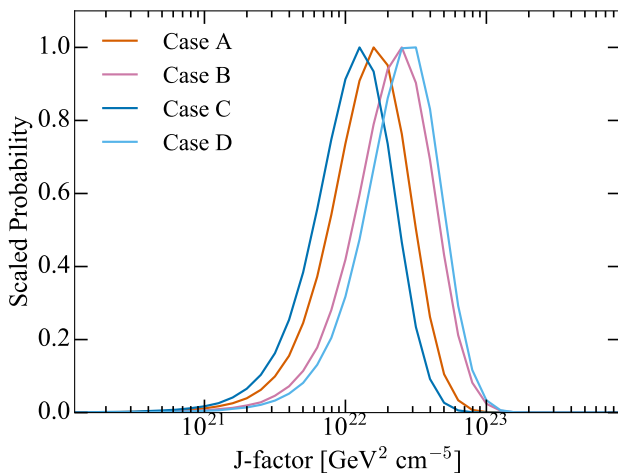


τ leptons cannot simultaneously explain the peak of the spectra and the high energy tail
Log-parabola can better explain the low energy data

Data Cases



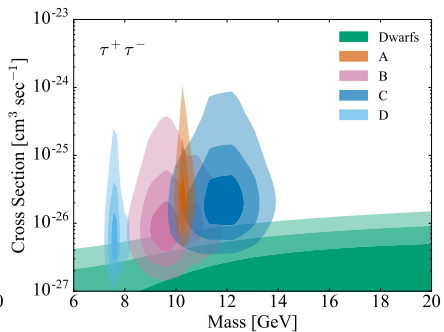
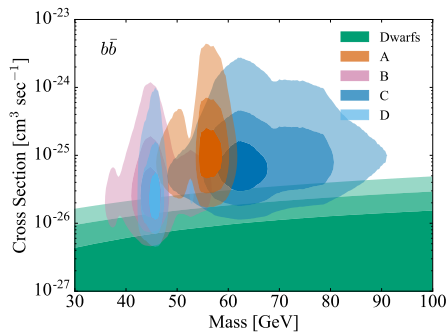
J-factors



$$\frac{d\Phi}{dE} = \frac{1}{4\pi} \frac{J}{m_{\chi}^2} \frac{\langle\sigma v\rangle}{2} \frac{dN}{dE} \quad (24)$$

$$J = \int d\Omega \int dz \rho^2(z, \Omega) \quad (25)$$

Tension



Evidence Ratios

Evidence ratios for our five models using the diffuse templates for our various background cases.

Model	Case A	Case B	Case C	Case D
DM: $b\bar{b}$	3600	21	220	15
DM: $\tau^+\tau^-$	2300	25	230	29
Log-Parabola	0.69	0.58	0.71	0.54
Exponential Cutoff	0.73	0.59	0.78	0.54
SIDM	1.1	1.2	1.2	1.1

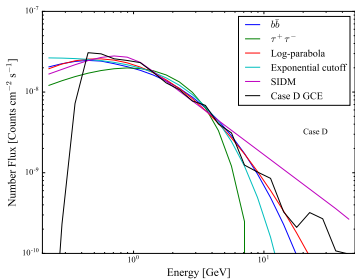
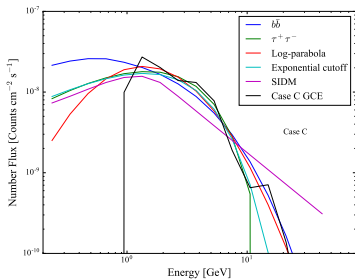
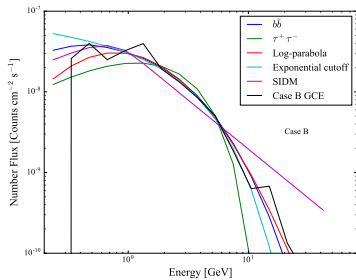
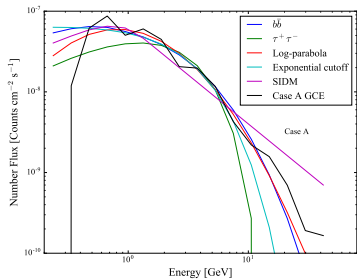
Bayes Factors

$$K_{12} = \frac{p(M_1|D)p(M_2)}{p(M_2|D)p(M_1)} = \frac{p(D|M_1)}{p(D|M_2)}. \quad (26)$$

Bayes factors for the considered models, relative to the $b\bar{b}$ model, for each of the different background cases. Values larger than one indicate the data prefer that model over $b\bar{b}$

Model	Case A	Case B	Case C	Case D
DM: $\tau^+\tau^-$	4×10^{-24}	1×10^{-5}	7×10^4	1×10^{-22}
Log-Parabola	3×10^{12}	9×10^6	2×10^{12}	5×10^9
Exponential Cutoff	8×10^7	2×10^4	4×10^{10}	0.1
SIDM	5×10^{-20}	8×10^{-19}	6×10^{-2}	0.1

Best Fits



Concluding Remarks

- ▶ DM interpretations display a strong to definitive tension between the bright GCE and the dim dwarfs
- ▶ Astrophysical interpretations pick out a consistent scale between the GCE and the low significance excesses.
- ▶ The GCE is precise enough to prefer (in all data cases) the log-parabola spectrum over any DM spectrum