The Proton Radius Puzzle- Why we all should care

Gerald A. Miller, University of Washington

Pohl et al Nature 466, 213 (8 July 2010)

 $dG_E(Q^2)$ $\begin{array}{c} \hline \end{array}$ 4 % Difference

*dQ*²

 $r_p^2 \equiv -6$

 $Q^2 = 0$ Pohl, Gilman, Miller, Pachucki (ARNPS63, 2013) C Carlson PPNP 82,59(2015)

muon H $r_p = 0.84184(67)$ fm electron H r_{p} =0.8768 (69)fm Large electron-p scattering $r_p = 0.875$ (10)fm PRad at JLab- lower Q² 1 Small Weizhi Xiong $3 \times 10^4 \le Q^2 \le 5 \times 10^{-2} \text{ GeV}^2$ $7.7 \times 10^{-3} \leq Q^2 \leq 0.13$ fm⁻² C. Weiss

4 % in radius: why care?

Can't be calculated to that accuracy? Sergey Syritsyn

Is the muon-proton interaction the same as the electron-proton interaction? violation of universality connections with muon g-2? connections with LHCb ? Something happening here ???

Outline - a) review history experiments b) List & explain possible resolutions

The proton radius puzzle In a picture

Randolf Pohl Mainz, $2nd$ June 2014 11

So far

- difference between muon and electron hydrogen spectroscopy re the proton radius
- similar effect in the deuteron, implies effect on neutron
- no effect in 3He (large error bars)
- no effect in 4He (small error bars)
- any explanation must account for the above, assumes there is a difference between muon and electron hydrogen spectroscopy

Beyer et al Science 358,79 electron H sidered in the analysis of only one of the previous $H_{\alpha\nu\alpha\beta\alpha}$ for the particular experimental scheme (19). The scheme (19). The scheme (19). The scheme (19). The scheme (19 \overline{a} and \overline{a} experience is the line of the line center \overline{AB} to belende oppoli \mathbf{t} and working at working at working \mathbf{t} d alactran H, the results of \sim cicl

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26 Jan., 2018 1801.08816 Paris 1S-3S hydrogen 4 Phys. Rev. Lett. 120, 183001 (2018)

 $\mathcal{O}(\mathcal{O}_2)$. As the signal height depends on \mathcal{O}_2 $\overline{\mathsf{S}}$ elements $\overline{\mathsf{S}}$ Is electron-hydrogen spectroscopy accurate enough?

Possible resolutions

- Electron H spectroscopy not so accurate
- Strong interaction effect in two photon exchange diagram
- Muon interacts differently than electron!- new scalar boson

The last two resolutions could be halved and not be in conflict with data

Shift from history to our efforts to explain

Suppose radii extracted in earlier H experiments is correct, some new muon effect is responsible

What energy difference corresponds to 4% in radius?

Measured = 206.2949(32) = 206.0573(45)-5.2262 r_p^2 +0.0347 r_p^3 meV computed

Explain puzzle with radius as in earlier H atom measures: increase 206.0573 meV by 0.31 meV-attractive effect on 2S state

Can go half way and not disagree with data

Two photon exchange

Measured = 206.2949(32) = 206.0573(45)-5.2262 r_p^2 +0.0347 r_p^3 meV computed

Miller PLB 2012 N F1(−q2) + F1(−q2)F(−q2)

will be the same as the same as the charge of a free demanded by current conservation as expressed through

the Ward-Takahashi identity [24, 25]. We assume

energy shift proportional to lepton mass⁴

 D_{0} where \mathbf{r} $\frac{1}{2}$ is a geven to about 1 GeV2 needed here. F($\frac{1}{2}$ DUL and LUL operator that projects on the on-mass-shell proton state. We use O^a unless otherwise stated. Γ Γ Γ \Box U. Carrier count f of virtual Compton sch Re part of virtual Compton scattering lepton; the solid line, the nucleon; the wavy lines photons; Im part is measured \mathbf{r} interference between one only part \mathbf{r} but unknown subtraction function is needed Can account for 0.31 meV, no conflict with e-H use dispersion relations Im *T*1(⌫*,Q*)

 e^+/e^- and μ^+/μ^- scattering on proton

So what? MUSE expt using RF time measurements. Magnet polarities can be reversed to allow the channel to transport

<http://www.physics.rutgers.edu/~rgilman/elasticmup/>

PSI proposal R-12-01.1 Paul Reimer

- constrains two photon effect, which still survives at significant level
- Fig. 3. A Georgia showing part of the MUSE experimental system. Here \mathbf{H} leptons see large radius • if large radius correct and no two photon all
- and a small beam spot (*x,y <* 1 cm) at the scattering target. The base line design for the MUSE beam detectors has a collimator and a scintillating fiber detector (SciFi) at the intermediate focus. **Some of the detectors in the target region are shown in Fig. 3. After the channel and immediate region are shown in Fig. 3. After the channel and immediate region and immediate region and immediate region and immediate re** • if small radius correct and no two photon all
- chambers. The second precision beam distributed beam and the target in the target of the target of the target The IFP collimator serves to cut the \mathcal{N} channel acceptance to reduce the beam flux to reduce th manageable levels. The IFP SciFi measures the RF time, for use in determining particle type, • will not see a new light particle, but all leptons see large radius

muon anomalous moment strong and electroweak interactions) by measuring *a*exp l and comparing the muon and and comparing the sense measurements with a structure with \mathbf{a}_0 calculated to much higher order in perturbation theory. Such comparisons test strong and electroweak interactions) by measuring *a*exp *^l* for the electron and muon ever muon anomalous colculation to much higher in period to much comparison theory.

H

Figure 1 The first-order QED correction to g-2 of the muon.

and its SU(3)*C*× SU(2)*L*× U(1)*^Y* standard-model (SM) extension (which includes

3.6 st. dev anomaly now fix add heavy photon interacts preferentially with muon Muon data is g-2 - BNL exp't, Hertzog- FermiLab now...

> Postulate new scalar boson!

 ϕ ϕ of 4 momentum conservation elastic ep scattering Dark Light interest

New scalar bosons

assumes puzzle exists

- give μ-p Lamb shift
- almost no hyperfine in μ proton
- small effect for D, almost no effect $4He$
- consistent with g-2 of ^μ and electron
- avoid many other constraints
- be found
PRL 117, 101801 (2016) PHYSICAL REVIEW LETTERS week ending

2 SEPTEMBER 2016

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Electrophobic Scalar Boson and Muonic Puzzles

Yu-Sheng Liu, David McKeen,[†] and Gerald A. Miller[‡] Department of Physics, University of Washington, Seattle, Washington 98195-1560, USA

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New scalar bosons

 \bigvee

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Electrophobic Scalar Boson and Muonic Puzzles

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mϕ (MeV)

 r om 200 KeV to 3 MeV Allowed mass range from 200 KeV to 3 MeV

Summary

- If the proton radius puzzle exists : new scalar boson of mass from 300 KeV to 3 MeV may exist- narrow target
- Direct detection is needed.

Does 4% matter?

Summary

- If the proton radius puzzle exists : new scalar boson of mass from 300 KeV to 3 MeV may exist- narrow target
- Direct detection is needed.

Does 4% matter?

Spares follow

 10^{-1} 1 1 10 10²

 m_{ϕ} (MeV)

 $Δa_μ$

19

Unshaded allowed by muon g-2 and muon-p Lamb shift Two anomalies have same scale

Unshaded allowed by muon g-2

Unshaded allowed by muon g-2 and muon-p Lamb shift Two anomalies have same scale

Unshaded allowed by muon g-2

Using new eH experiment $\epsilon_{\mu} \epsilon_{p}$ is reduced by a factor of 3: barely visible in loglog plot

Nuclear Physics constraints ϵ_n/ϵ_p

If $\epsilon_n = \epsilon_p$ scalar is ruled out by n-²⁰⁸Pb scattering If ϵ_n has opposite sign as ϵ_p parameter space widens Other constraints:

NN scattering, nuclear matter $&^3$ He $-{}^3$ H binding muonic D, ³He 2 $n - {}^{208}Pb$

electron Exclusion plot

More constrains Coupling to ϵ electron ϵ_e

 e^+e^- resonant (?) scattering not seen Hydrogen atom Lamb shift $\delta E^{H}_{L} < 45$ kHz

electron anomalous magnetic moment

BD is beam dump discussed next Electrophobic

Recast in EFT- parameters seem natural

So far size of this term cannot be determined from theory-experiment is needed lattice QCD may disprove above sentence

Science 12 Aug 2016: Vol. 353, Issue 6300, pp. 669-673

Deuteron is smaller too

Electron (D-H) isotope shift (2S-1S) 2 photon spectroscopy PRL 104, 233001 $r_d^2 - r_p^2 = 3.82007(65)$

 $\mu - D$ Lamb shift $r_d = 2.12562(78)$ fm Science 353 (2016) 669

CODATA $(2010)r_d = 2.1424(26)$ fm - mainly electron scattering

Use $r_p = 0.84087$ in $r_d^2 - r_p^2 = 3.82007(65)$ gives $r_d = 2.12769$ fm

 μ D and Electron (D-H) isotope shift are consistent \rightarrow redo *eD* scattering?

If NO proton radius puzzle, there still is a missing Lamb shift OR: CODATA deuteron radius is too large 2.1424 vs 2.12769 remeasure deuteron ?

Nuclear Physics constraints ϵ_n/ϵ_p

- Low energy scattering of neutrons on ²⁰⁸Pb using ϕ -nucleon coupling q_N . $\frac{g_N^2}{e^2} \to \frac{A-Z}{A} \epsilon_n^2 + \frac{Z}{A} \epsilon_p \epsilon_n$ cancellation evades previous limits
- NN charge-independence breaking scattering length $\Delta a = (a_{pp} + a_{nn})/2 - a_{np}$, measured: 5.64(60) fm, theory: 5.6(5) Scalar boson exchange: $\Delta a_{\phi} \propto \int_0^{\infty} \Delta V \bar{u} u_{np} dr \leq 1.6$ fm (2 S.D.)
- Change in binding energy/A infinite nuclear matter: less than 1 MeV
- binding energy $B(^3He) B(^3H) = 763.76$ keV due to Coulomb (693 keV) + strong force charge symmetry breaking (68 keV) ϕ exchange $<$ 30 keV

MUSE and scalar

$$
V_{\phi}(r) = -1.7 \times 10^{-6} \alpha \frac{e^{-m_{\phi}r}}{r}
$$

- No scattering experiment can detect a coupling this weak
- If this scalar exists (and other e H large radius experiments correct) MUSE will find electrons/positrons see the same large radius and
- muons and anti-muons will see the same large radius
- If no puzzle: all leptons see small radius
- In any case MUSE limits two photon (biggest uncertainty)

Randolf Pohl

Aldo Antognini

Fig. 1. (A) Formation of μ p in highly excited states and subsequent cascade with emission of "prompt" $K_{\alpha, \beta, \gamma}$. (B) Laser excitation of the 2S-2P transition with subsequent decay to the ground state with K_{α} emission. (C) 2S and 2P energy levels. The measured transitions v_s and v_t are indicated together with the Lamb shift, 2S-HFS, and 2P-fine and hyperfine splitting.

PHYSICAL REVIEW D 95, 036010 (2017)

Validity of the Weizsäcker-Williams approximation and the analysis of beam dump experiments: Production of a new scalar boson

Yu-Sheng Liu, $\check{\,}$ David McKeen, $\check{\,}$ and Gerald A. Miller $\check{\,}$

- previous cross sections obtained w. WW approximation cross sections obtained w.
- cross sections not accurate approximations [including the Weizsäcker-Williams (WW) approximation] or Monte-Carlo simulations.
- \bullet exclusion plots changed substantially exclusion plots changed substantially plots differ by substantial amounts when seen on a linear scale. In the event of a discovery, we generate
- if discovery, WW gives wrong parameters use of approximations to analyze the pseudodata for the future experiments is shown to lead to considerable errors in determining the parameters. Furthermore, a new region of parameter space can be explored
- not necessary to assume mass of new particle is much much greater than mass of electron not necessary t

Lepton-universality violating one boson exchange

- Tucker-Smith & Yavin PRD83, 101702 new particle scalar or vector coupling
- Brax & Burrage scalar particles PRD 83, 035020 &'14
- Batell, McKeen & Pospelov PRL 107, 011803 new gauge boson kinetically mixing with F^{μν} plus scalar for muon mag. mom. 1401.6154 W decays enhanced
- Carlson Rislow PRD 86, 035013 fine tune scalar pseudoscalar or polar and axial vector couplings
- Barger et al PRL106,153001 new particles ruled out but assumes universal coupling
- Kaon decays provide constraints 31

Looking for new scalars is not new
\nLow mass Higgs searches
\n
$$
p(1.88 \text{ MeV}) + {}^{19}F \rightarrow \alpha + {}^{16}O^*(6.05)
$$

\n ${}^{16}O^*(6.05) \rightarrow {}^{16}O(GS) + \phi$

Kohler et al PRL 33, 1628 (1974)

Freedman et al. PRL 52, 240 (1984)

 $p+3$ *H* \rightarrow ⁴ *He*(20*.1*) \rightarrow 4*He*(*GS*) + ϕ

No Scalars found, but assumed coupling constants were much larger than what we will use

Nuclear dependence of short-ranged mu-p effects The first example we consider is the model of Ref. \sim 14]. To make a prediction for the nucleus, one needs to know the nucleus

• Energy shift is proportional to square of muon wave function at the origin **the domain cone down a** resulting \mathbf{r} and \mathbf{r} is proportional to square of ρ result, if one considers the electron-deuteron atom, the e q

GAM • Suppose you have effect that gives energy shifts Ep (on proton) E_n (on neutron) $\mathcal{L}_{\text{Edd} \text{Odd} \text{Odd}}$ $\mathcal{L}_{\text{Edd} \text{Odd}}$ budgets the proton to respect that $\mathcal{L}_{\text{Edd} \text{Odd} \text{Odd} \text{Odd}}$ GAM 1501.01036

$$
E_A = \left(\frac{1 + \frac{m_{\mu}}{m_p}}{1 + \frac{m_{\mu}}{Am_p}}\right)^3 Z_{\mu}^3 (ZE_p + NE_n) \left(1 - \mathcal{O}\left(\frac{R_A^2}{a_{\mu}^2}\right)\right) \approx \left(\frac{1 + \frac{m_{\mu}}{m_p}}{1 + \frac{m_{\mu}}{Am_p}}\right)^3 Z^3 (ZE_p + NE_n),
$$

a¹

Nuclear shift interactions in the such contact interactions in the such at the such at the such at the such atomic number. The such at the such at the such at the such atomic number of the such at the such atomic number. T

In particular, the prediction of Ref. [4](with *^Eⁿ* = 2*/*3*Ep*) for ⁴He is a Lamb shift that is (1.27) 8 (2)5/3⇡ 10 meV, a huge number. The expression Eq. (23) applies to all models in which the contribution to the Lamb shift enters Square of wave fun Counting

 M_V model: ~ 0.3 meV (1+0.3)(9)(2) = 6.3 meV shout 6 st dev My model: ~0.3 meV (1+0.3)(8)(2) =-6.3 meV about 6 st. dev **E** any short range idea Almost unknown $\overline{T}_1(0,Q^2)$ Miller PLB 2012

$$
\Delta E^{\text{subt}} = \frac{\alpha^2}{m} \Psi_S^2(0) \int_0^\infty dQ^2 \frac{h(Q^2)}{Q^2} \overline{T}_1(0, Q^2) \text{ Soft proton}
$$

\n
$$
\lim_{Q^2 \to \infty} h(Q^2) \sim \frac{2m^2}{Q^2}, \text{ chiral PT}: \overline{T}_1(0, Q^2) = \frac{\beta_M}{\alpha} Q^2 + \cdots
$$

\n
$$
\rightarrow \text{Logarithmic divergence}
$$

\n
$$
\overline{T}_1(0, Q^2) \rightarrow \frac{\beta_M}{\alpha} Q^2 F_{\text{loop}}(Q^2) \text{ Cuts off integral}
$$

\nBirse & McGovern assume dipole: $\Delta E^{\text{subt}} = 0.004 \text{ meV}$ very small
\nMiller $F_{\text{loop}}(Q^2) = \left(\frac{Q^2}{M_0^2}\right)^n \frac{1}{(1 + aQ^2)^N}, n \ge 2, N \ge \text{ n+3}$
\nInfinite parameter set gets needed 0.31meV, NO constraint on neutron
\nChoose parameters so shift in proton mass $\le 0.5 \text{ MeV}$
\n(current uncertainty)

Recast in EFT- parameters seem natural

Two photon exchange ww.elsevier.com/locate/physletbane

by 0.31 meV-attractive effect on 2S state needed Measured = 206.2949(32) = 206.0573(45)-5.2262 r_p^2 +0.0347 r_p^3 meV computed **Explain puzzle with radius as in H atom increase 206.0573 meV** Proton polarizability contribution: Muonic hydrogen Lamb shift and elastic

> N F1(−q2) + F1(−q2)F1(−q2)F1(−q2)F1(−q2) a,b,c (2) M

energy shift proportional to lepton mass⁴

$$
T^{\mu\nu} = \underbrace{q^{\mu\nu} = \bigoplus_{\alpha} q^{\beta} \bigotimes_{\alpha} q^{\beta}}_{\alpha} \uparrow q^{\beta} \uparrow q^
$$

shell form factor, and *Dispersion rela* operator that projects on the on-mass-shell proton state. I_{α} $\text{Part} \& \text{crit}$ que de charge of the off-shell proton and $\frac{1}{2}$ ω as the charge of a demanded by current conservation by current conservation as expressed through \mathcal{R} Hill & Paz- big uncert on: $Im[T_1] \propto W_1$ mossured on. $I\mathfrak{m}[1] \propto W$ incasure α evaluation operator W_{ϵ} and integral σ _v inty (2π)⁴ $\dot{\ }$ in disp Hill & Paz- big uncertainty in dispersion approach ³⁵ $\sum_{i=1}^n$ by about $\sum_{i=1}^n$ Dispersion relation: $Im[T_1] \propto W_1$ measured \mathbf{T} standard deviations for the present \mathbf{T} drogen vil α is been to be very extensive and the very extensive and highly accurate muonic H σ uduracur account for the \mathbb{Z} meV. This is worthy of consideration be-Large virtual photon energy ν , $W_1 \sim \nu$ integral over energy diverges $\begin{array}{ccc} \n\bullet & \bullet & \bullet & \bullet & \bullet \end{array}$ atunction needed: *T* Subtraction function needed: $\bar{T}_1(0, Q^2)$ zero energy

Several new electron spectroscopy experiments

- Independent measurement of Rydberg constant. This would change only extracted $r_{\rm p}$ nothing else
- 2S-6S UK, 2S-4P Germany, IS-3S France
- 2S-2P classic, Canada
- Highly charged single electron ions NIST

2S-4P has reported preliminary results- small radius not yet published

Yes it really is G_E

- Non-relativistic reduction of one-photon exchange leads to the spin independent interaction being $G_E (Q^2)/Q^2$
- All recoil effects properly accounted for:Breit-Pauli Hamiltonian computed for non-zero lepton and proton momentum

U. D. Jentschura Phys.Rev.A88 (2013) 062514 Department of Physics, Missouri University of Science and Technology, Rolla, Missouri, MO65409-0640, USA and

If we assume an average of roughly 0.7×10-7 light sea positrons per valence quark, Thow that virtual electron-positron annihilation processes lead. then we can show that virtual electron-positron annihilation processes lead to an extra term in the electron-proton versus muon-proton interaction, which has the right sign perturbative processes could appear to lead to such a nonunversality. One of these is a coupling of and magnitude to explain the proton radius discrepancy. positron that is also part of the *ee*⁺ cloud of the proton. See Fig. 1. The term non-perturbative here refers to a term in the creen on procon versus muon procon incent

 $\frac{1}{2}$ n-perturbative lepton-pair exists in proton wave function. IIDI: energy shift ble for the stability of the universe (the hydrogen atom atom β $\propto 1/m_l^2$, from annihilation at rest. GAM: Shift $\propto 1/(\text{constituent quark mass})^2$ electron-positron-positron-positron-positron pair, a neutron \mathbf{r} have effect is small and same \mathbf{r} of the electromagnetic interaction and \mathbf{r} α priori, one would then α priori, one would then the think that the think that the think that the think the think that the think the think Non-perturbative lepton-pair exists in proton wave function. UDJ: energy shift \mathbf{f} to proton is non-universal, then it is conceivable in \mathbf{f} or electron and muon aton $\overline{1}$ a note $\overline{2}$ for the canceler of perturbative of p higher-order effects which could appear to \sim 3*m*² *e* ${\rm rom}$ annihilation at rest. ${\rm GAM}\colon {\rm Shift} \propto 1/2$ the virtual annihilation of a bound electron of a bound electron with a bound sea lepton" (positron) inside the proton. The up (*u*) and down (*d*) quarks, which carry non-integer charge numbers, interact $\frac{1}{\sqrt{2}}$ from one il il tipe of ref. $\frac{1}{\sqrt{4M}}$ Cl if $\frac{1}{\sqrt{4M}}$ and $\frac{1}{\sqrt{4M}}$ and $\frac{1}{\sqrt{4M}}$ $\propto 1/m_l^2$, from annihilation at rest. GAM: Shift $\propto 1/(\text{constituent quark mass})^2$ Any effect is small and same for electron and muon atoms light sea lepton that annihilates with the bound electron. An example of a non-perturbative quantum chromodynamic (QCD) Any effect is small and same for electron and muon atoms
arXiv:1501.01036

| Arbitrary functions We found the previous literature by including a form \mathcal{F} of *Q*2. Our aim here is to more fully explore the uncertainty in the subtraction term that arises from the logarithmic divergence. We shall use a form of *F*loop(*Q*2) that is consistent with the constraint on the *Q*⁴ term found Birse & McGovern [20]. This is done

$$
\overline{T}_1(0, Q^2) = \frac{\beta_M}{\alpha} Q^2 F_{\text{loop}}(Q^2).
$$
\n
$$
F_{\text{loop}}(Q^2) = \left(\frac{Q^2}{M_0^2}\right)^n \frac{1}{(1 + aQ^2)^N}, \ n \ge 2, \ N \ge n + 3,
$$
\n
$$
\overline{T}_1(0, Q^2) \sim \frac{1}{Q^4} \text{ or faster}, \ \beta_M \to \beta
$$

$$
\Delta E^{\rm subt} \approx 3 \alpha^2 m \Psi_S^2(0) \frac{\beta}{\alpha} \gamma^n B(N,n), \gamma \equiv \frac{1}{M_0^2 a}
$$

dipole form. However, we shall determine the substraction of $\frac{dI}{d\Omega}$ Choose parameters such that shift in proton mass \leq electromagnetic uncertainty of 0.5 MeV 3 parameters: n, N, a $(M_0 = M_\beta)$

Almost unknown $\overline{T}_1(0,Q^2)$ Miller PLB 2012 $\Delta E^{\rm subt} =$ α^2 *m* $\Psi_S^2(0)$ \int^∞ 0 $dQ^2 \frac{h(Q^2)}{Q^2}$ $\frac{(Q}{Q^2}$ $\overline{T}_1(0,Q^2)$ lim $Q^2 \rightarrow \infty$ $h(Q^2) \sim$ $\frac{2m^2}{Q^2}$, chiral PT : $\overline{T}_1(0,Q^2) = \frac{\beta_M}{\alpha}$ $Q^2 + \cdots$ \rightarrow Logarithmic divergence $\overline{T}_1(0,Q^2)\rightarrow \frac{\beta_M}{\alpha}$ α $Q^2F_{\text{loop}}(Q^2)$ Cuts off integral Soft proton **Typo bleow**

Birse & McGovern assume dipole : $\Delta E^{\text{subt}} = 0.004 \,\text{meV}$ very small

Miller
$$
F_{\text{loop}}(Q^2) = \left(\frac{Q^2}{M_0^2}\right)^n \frac{1}{(1 + aQ^2)^N}, n \ge 2, N \ge N + 3
$$

Infinite parameter set gets needed 0*.*31meV*,* NO constraint on neutron Recast in EFT- parameters seem natural Choose parameters so shift in proton mass <0.5 MeV (current uncertainty)

If recast into effective field theory strength seems natural

Pohl et al. Table of calculations

Lamb shift: vacuum polarization many, many terms

Resolution 1- QED calcs not OK

↵

Table 1: All known radius-**independent** contributions to the Lamb shift in μ p from different authors, and the one we selected. We follow the nomenclature of Eides $et al.^7$ Table 7.1. Item $#8$ in Refs.^{2,5} is the sum of items #6 and #7, without the recent correction from Ref.^{12} . The error of #10 has been increased to 100% to account for a remark in Ref.⁷. Values are in meV and the uncertainties have been added in quadrature.

Table 2: All relevant radius-**dependent** contributions as summarized in Eides et al.⁷, compared to Refs.^{2,5}. Values are in meV and radii in fm.

Pohl et al. Table of calculations

Lamb shift: vacuum polarization many, many terms

Mostly irrelevanttheory replaced by experiment

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 α

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QED calcs expand in α

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EFT of μp interaction cases is to use μp Caswell Lepage '86

symptom that an inefficient technique has been used in the first technique has been used α more efficient way to pro-

- Compute Feynman diagram, remove log divergence using dimensional regularization **In the follogarithmic mass to the fourth power of the fourth power of being dimension of being dimension of being**
	- \bullet include counter term in Lagrangian

larizability contributions, that enter in the two-photon enter in the two-photon exchange term, see Fig. 1, ca
1, can be a fig. 1, can b

 $\frac{1}{2}$ CHOOSE A to get U.51 HIE V SHIIT Choose λ to get 0.31 meV shift

$$
\Delta E^{\rm subt}(DR) = \alpha^2 m \frac{\beta_M}{\alpha} \Psi_S^2(0)(\lambda + 5/4)
$$

$$
\Delta E^{\rm subt}(DR) = 0.31 \text{ meV} \rightarrow \boxed{\lambda = 769}
$$

 α (α ^M is extracted number. However, α cancellation between α cancellation be parameters of an intermediate D and diamagnetic effects of the pion cloud ω Ivatural units $\frac{\rho_M}{\alpha} \approx 4\pi/(4\pi) \pi$ butter α bavage 32 β_M (magnetic polarizability) = $3.1 \times 10^{-4} \text{fm}^3$ very small Natural units $\beta_M/\alpha \sim 4\pi/(4\pi f_\pi)^3$ Butler & Savage '92

$$
\mathcal{M}_2^{DR} = i \, 3.95 \, \alpha^2 m \frac{4\pi}{\Lambda_\chi^3} \overline{u}_f u_i \overline{U}_f U_i.
$$

The coefficient 3.95 \equiv natural size. Thus standard EFT techniques result in an effective result in an eff

Muonic hydrogen experiment and $r_{\rm p}$ | Thuoni

Electronic Hydrogen -Pohl

shift-have ~ 20 available $E(nS) \approxeq$ Need two levels to get Rydberg and Lamb

 $\frac{R_{\infty}}{1}$

 L_{1S}