

Single-top production in the Standard Model and beyond

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- Higher-order soft-gluon corrections
- t -channel and s -channel production
- tW production
- tZ production via anomalous couplings



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Higher-order corrections

QCD corrections are very significant for single-top production

Soft-gluon corrections are important
and they approximate exact results very well

We calculate/resum these soft corrections at NNLL accuracy
for the double-differential cross section

Finite-order expansions

Approximate NNLO (aN³NLO) and N³LO (aN³LO) predictions
for cross sections and differential distributions

Soft-gluon corrections

partonic processes

$$f_1(p_1) + f_2(p_2) \rightarrow t(p_t) + X$$

define

$$s = (p_1 + p_2)^2, \quad t = (p_1 - p_t)^2, \quad u = (p_2 - p_t)^2$$

and $s_4 = s + t + u - \sum m^2$

At partonic threshold $s_4 \rightarrow 0$

Soft corrections $\left[\frac{\ln^k(s_4/m_t^2)}{s_4} \right]_+$

For the order α_s^n corrections $k \leq 2n - 1$

Resum these soft corrections for the double-differential cross section

At NNLL accuracy we need two-loop soft anomalous dimensions

Soft-gluon Resummation

moments of the partonic cross section with moment variable N :

$$\hat{\sigma}(N) = \int (ds_4/s) e^{-Ns_4/s} \hat{\sigma}(s_4)$$

factorized expression for the cross section in $4 - \epsilon$ dimensions

$$\begin{aligned} \sigma^{f_1 f_2 \rightarrow tX}(N, \epsilon) &= H_{IL}^{f_1 f_2 \rightarrow tX}(\alpha_s(\mu_R)) S_{LI}^{f_1 f_2 \rightarrow tX}\left(\frac{m_t}{N\mu_F}, \alpha_s(\mu_R)\right) \\ &\quad \times \prod J_{\text{in}}(N, \mu_F, \epsilon) \prod J_{\text{out}}(N, \mu_F, \epsilon) \end{aligned}$$

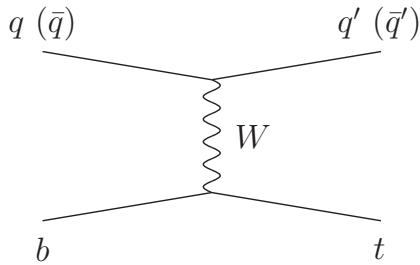
$H_{IL}^{f_1 f_2 \rightarrow tX}$ is hard function and $S_{LI}^{f_1 f_2 \rightarrow tX}$ is soft function

S_{LI} satisfies the renormalization group equation

$$\left(\mu \frac{\partial}{\partial \mu} + \beta(g_s) \frac{\partial}{\partial g_s} \right) S_{LI} = -(\Gamma_S^\dagger)_{LK} S_{KI} - S_{LK} (\Gamma_S)_{KI}$$

Soft anomalous dimension Γ_S controls the evolution of the soft function which gives the exponentiation of logarithms of N

t-channel production



At one loop

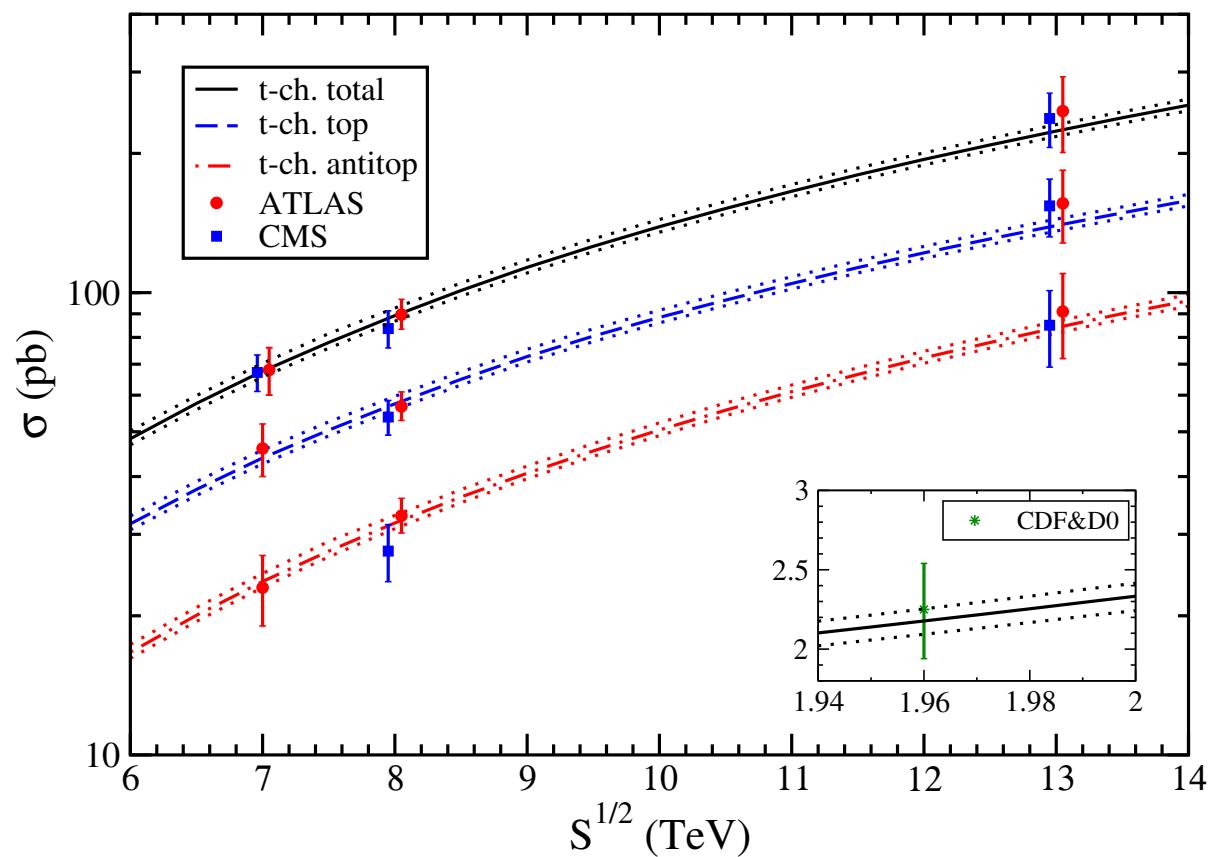
$$\begin{aligned}\Gamma_{S\ 11}^{t\ (1)} &= C_F \left[\ln \left(\frac{t(t - m_t^2)}{m_t s^{3/2}} \right) - \frac{1}{2} \right] & \Gamma_{S\ 12}^{t\ (1)} &= \frac{C_F}{2N} \ln \left(\frac{u(u - m_t^2)}{s(s - m_t^2)} \right) & \Gamma_{S\ 21}^{t\ (1)} &= \ln \left(\frac{u(u - m_t^2)}{s(s - m_t^2)} \right) \\ \Gamma_{S\ 22}^{t\ (1)} &= C_F \ln \left(\frac{s - m_t^2}{m_t \sqrt{s}} \right) - \frac{1}{2N} \ln \left(\frac{t(t - m_t^2)}{s(s - m_t^2)} \right) + \frac{(N^2 - 2)}{2N} \ln \left(\frac{u(u - m_t^2)}{s(s - m_t^2)} \right) - \frac{C_F}{2}\end{aligned}$$

At two loops

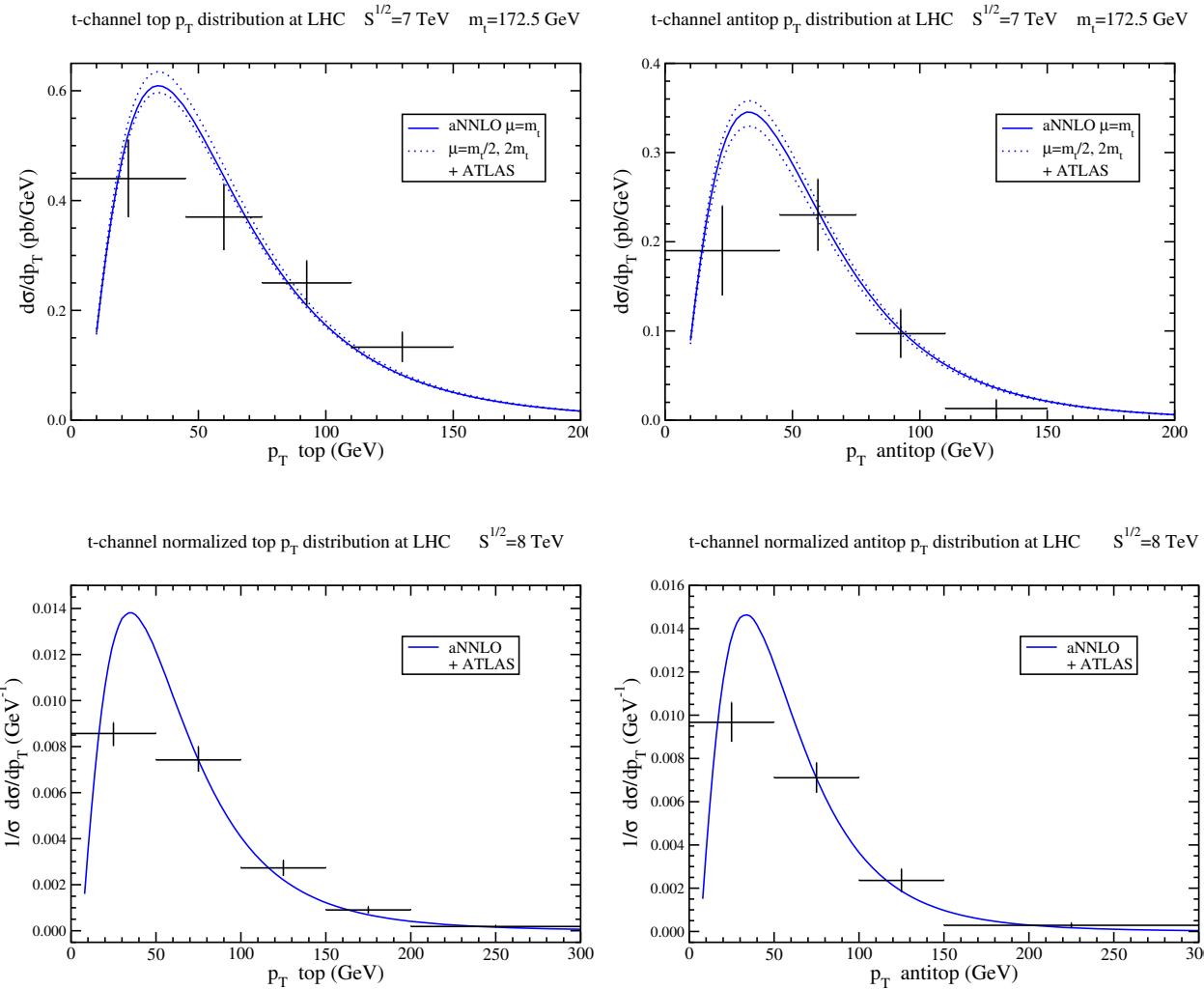
$$\Gamma_{S\ 11}^{t\ (2)} = \left[C_A \left(\frac{67}{36} - \frac{\zeta_2}{2} \right) - \frac{5}{18} n_f \right] \Gamma_{S\ 11}^{t\ (1)} + C_F C_A \frac{(1 - \zeta_3)}{4}$$

t-channel production at aNNLO

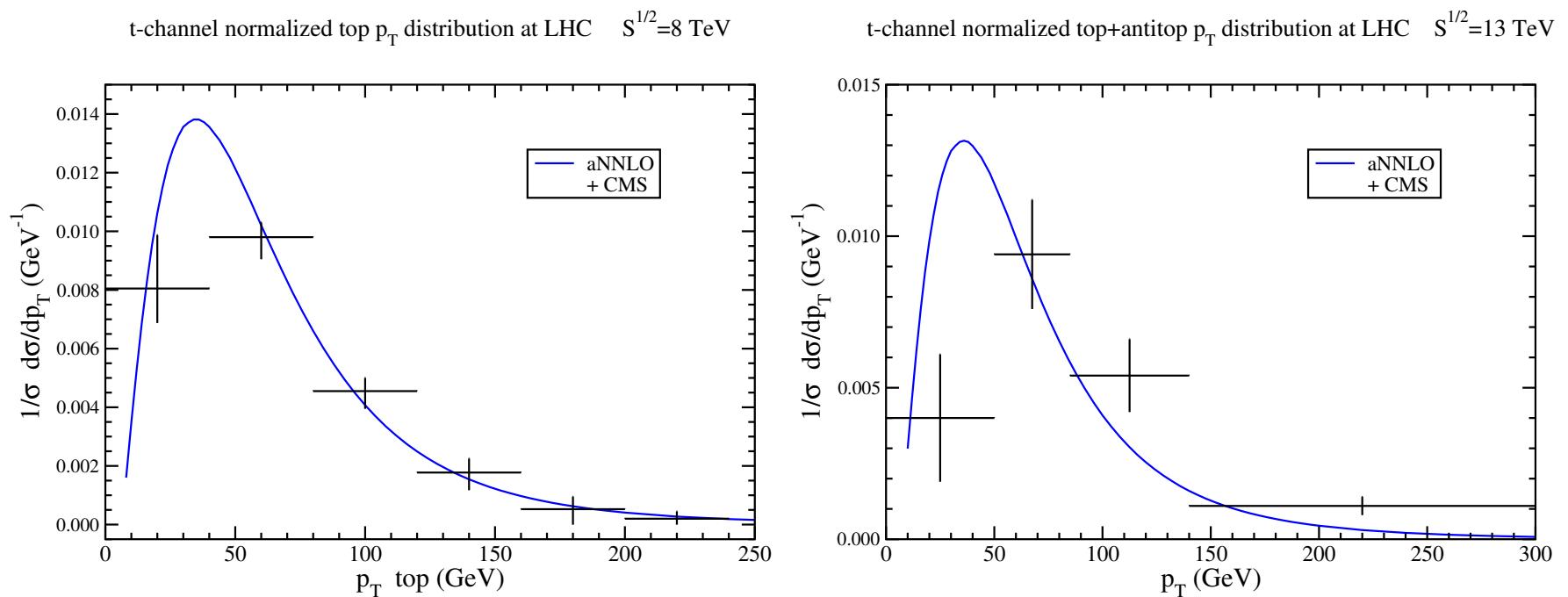
Single-top *t*-channel aNNLO cross sections $m_t = 172.5 \text{ GeV}$



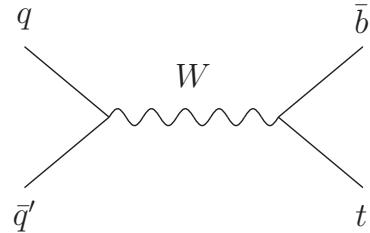
Top p_T distributions in t -channel production at the LHC



Top p_T distributions in t -channel production at the LHC



s-channel production



At one loop

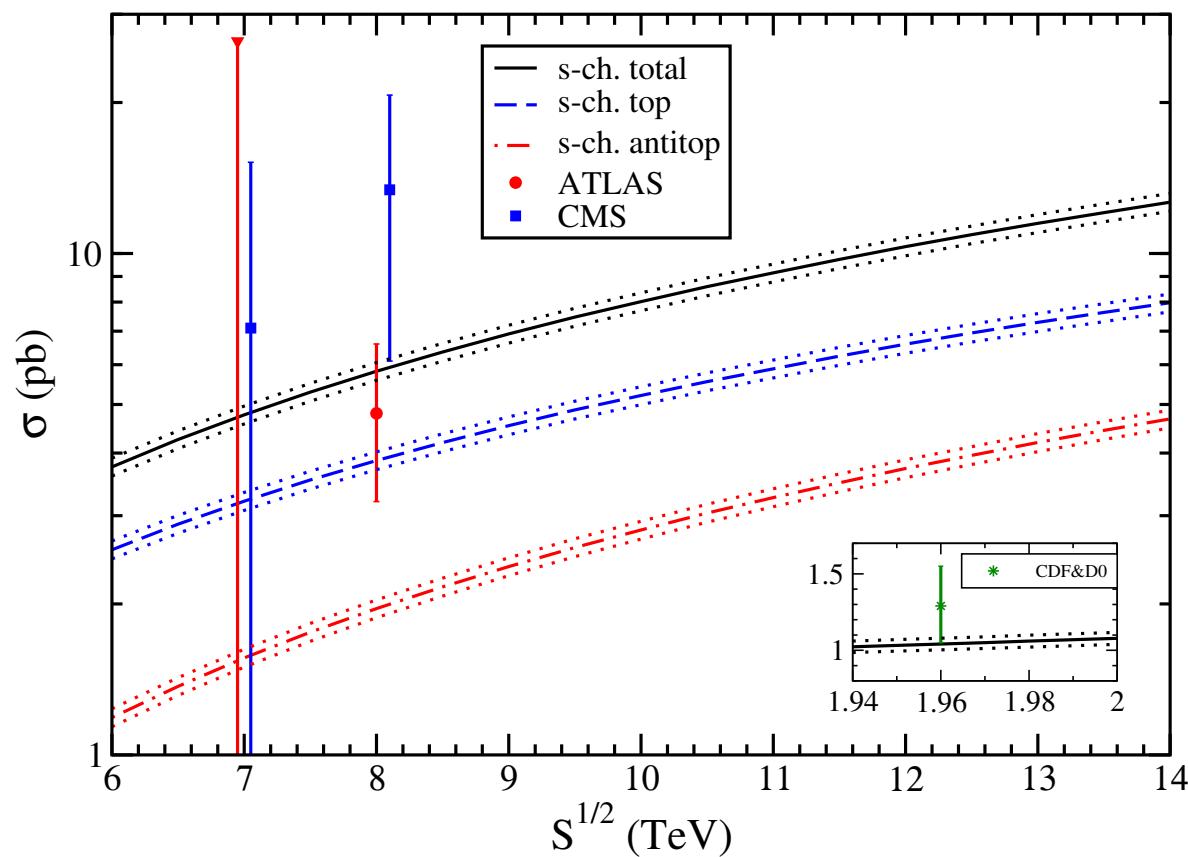
$$\begin{aligned}\Gamma_{S\ 11}^{s\ (1)} &= C_F \left[\ln \left(\frac{s - m_t^2}{m_t \sqrt{s}} \right) - \frac{1}{2} \right] & \Gamma_{S\ 12}^{s\ (1)} &= \frac{C_F}{2N} \ln \left(\frac{u(u - m_t^2)}{t(t - m_t^2)} \right) & \Gamma_{S\ 21}^{s\ (1)} &= \ln \left(\frac{u(u - m_t^2)}{t(t - m_t^2)} \right) \\ \Gamma_{S\ 22}^{s\ (1)} &= C_F \ln \left(\frac{s - m_t^2}{m_t \sqrt{s}} \right) - \frac{1}{N} \ln \left(\frac{u(u - m_t^2)}{t(t - m_t^2)} \right) + \frac{N}{2} \ln \left(\frac{u(u - m_t^2)}{s(s - m_t^2)} \right) - \frac{C_F}{2}\end{aligned}$$

At two loops

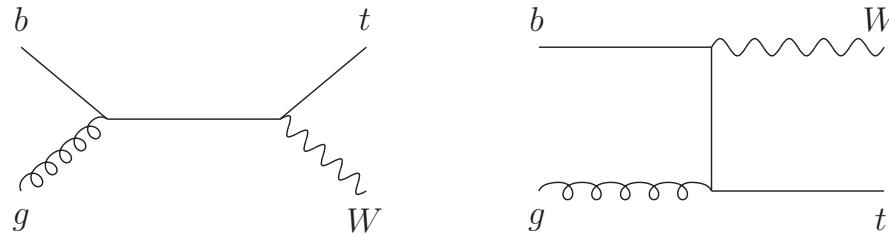
$$\Gamma_{S\ 11}^{s\ (2)} = \left[C_A \left(\frac{67}{36} - \frac{\zeta_2}{2} \right) - \frac{5}{18} n_f \right] \Gamma_{S\ 11}^{s\ (1)} + C_F C_A \frac{(1 - \zeta_3)}{4}$$

s-channel production at aNNLO

Single-top s-channel aNNLO cross sections $m_t = 172.5 \text{ GeV}$



Associated tW production



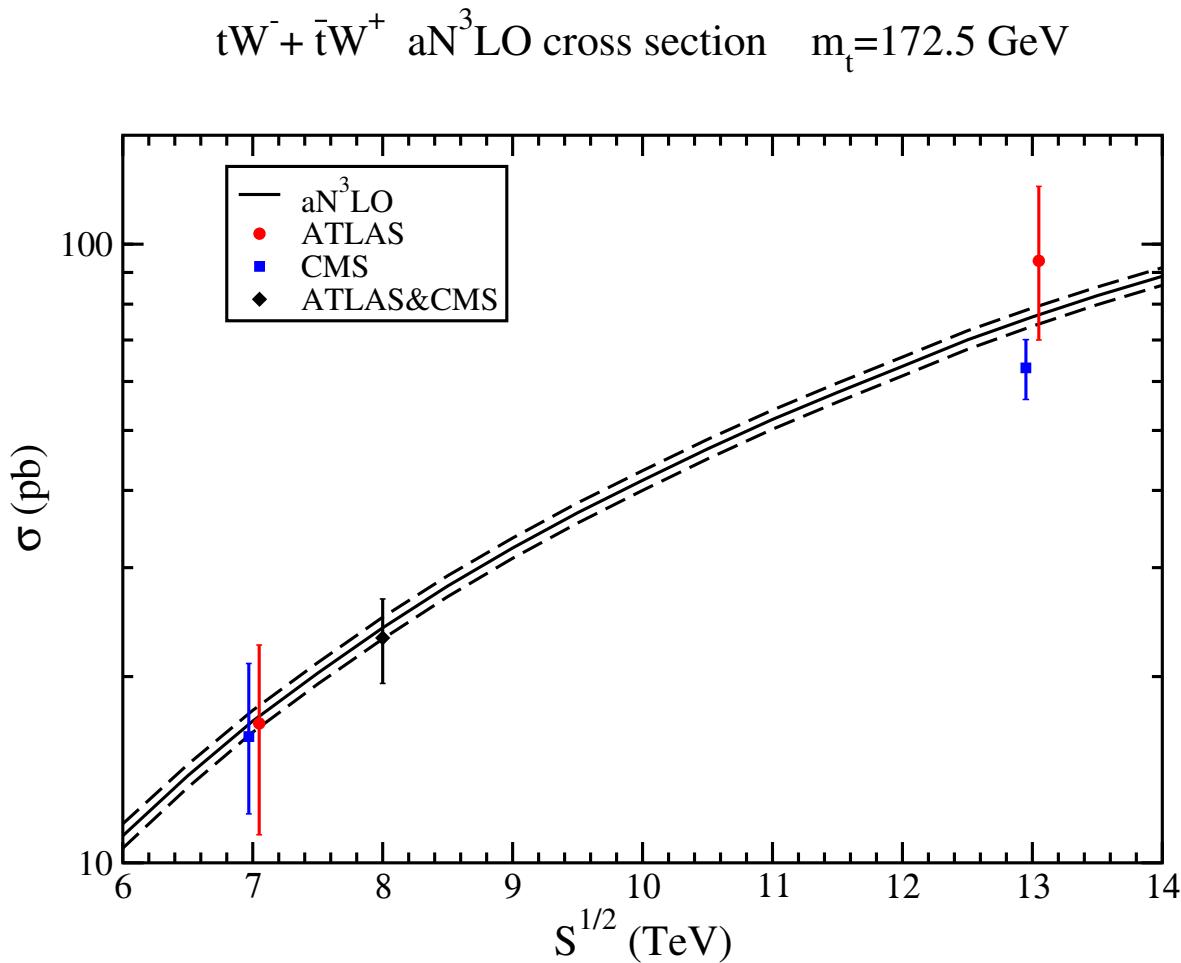
At one loop

$$\Gamma_S^{tW^- \text{ (1)}} = C_F \left[\ln \left(\frac{m_t^2 - t}{m_t \sqrt{s}} \right) - \frac{1}{2} \right] + \frac{C_A}{2} \ln \left(\frac{m_t^2 - u}{m_t^2 - t} \right)$$

At two loops

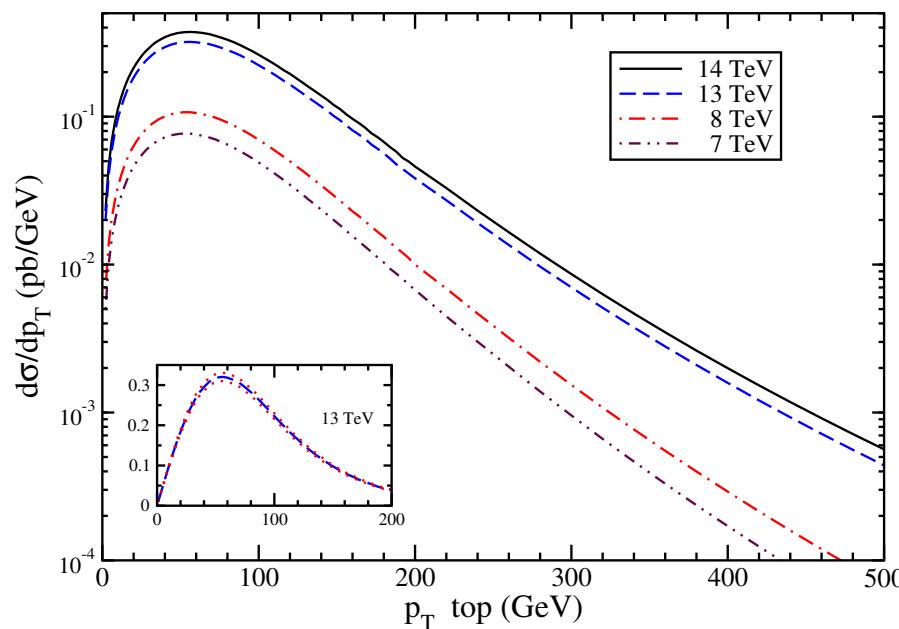
$$\Gamma_S^{tW^- \text{ (2)}} = \left[C_A \left(\frac{67}{36} - \frac{\zeta_2}{2} \right) - \frac{5}{18} n_f \right] \Gamma_S^{tW^- \text{ (1)}} + C_F C_A \frac{(1 - \zeta_3)}{4}$$

tW production at aN³LO

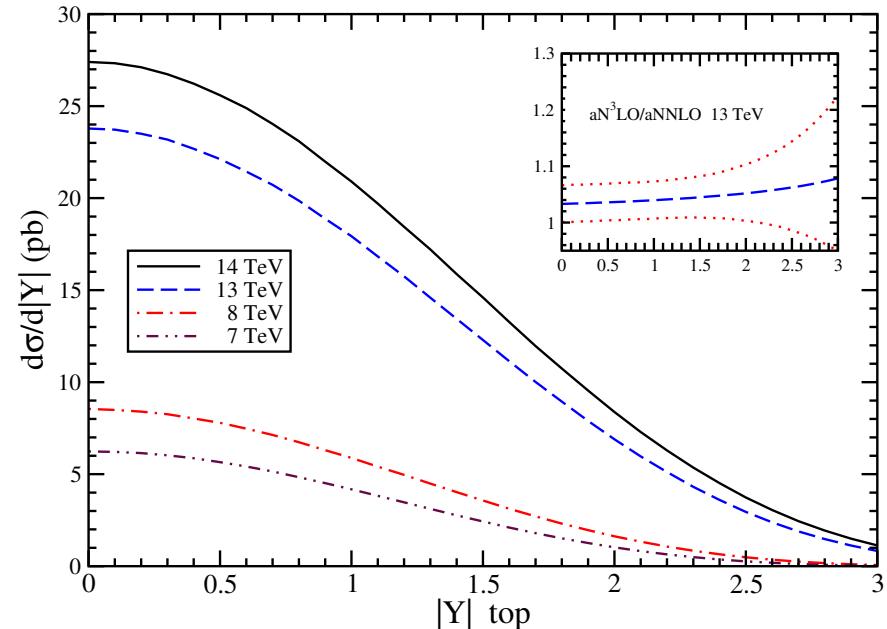


Top quark p_T and rapidity distributions in tW^- production

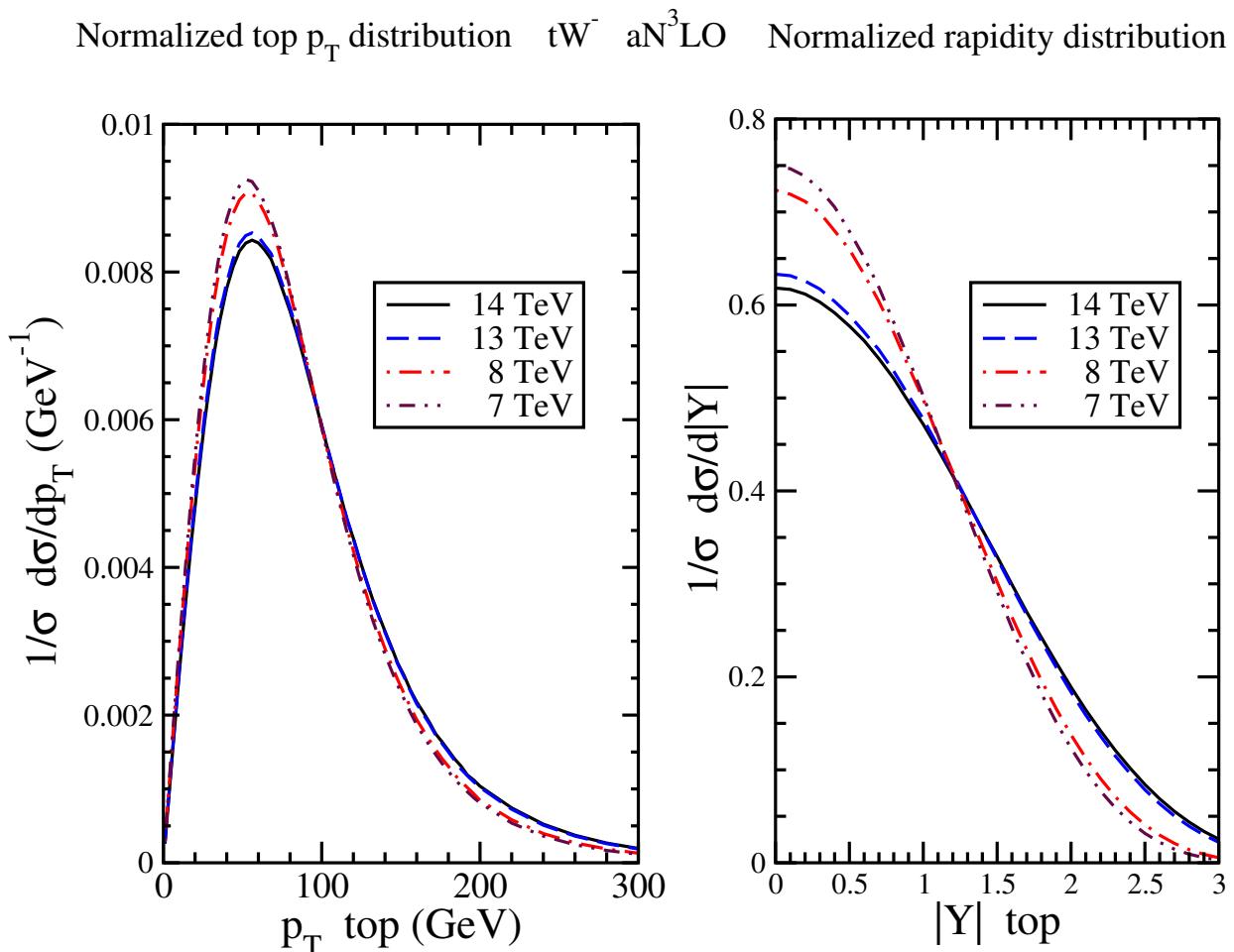
Top p_T distribution in tW^- production at LHC aN³LO $m_t=173.3$ GeV



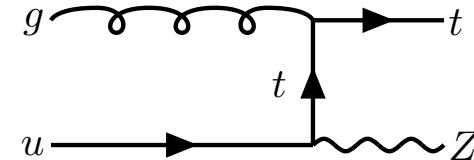
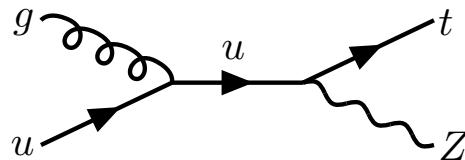
Top rapidity distribution in tW^- production at LHC aN³LO $m_t=173.3$ GeV



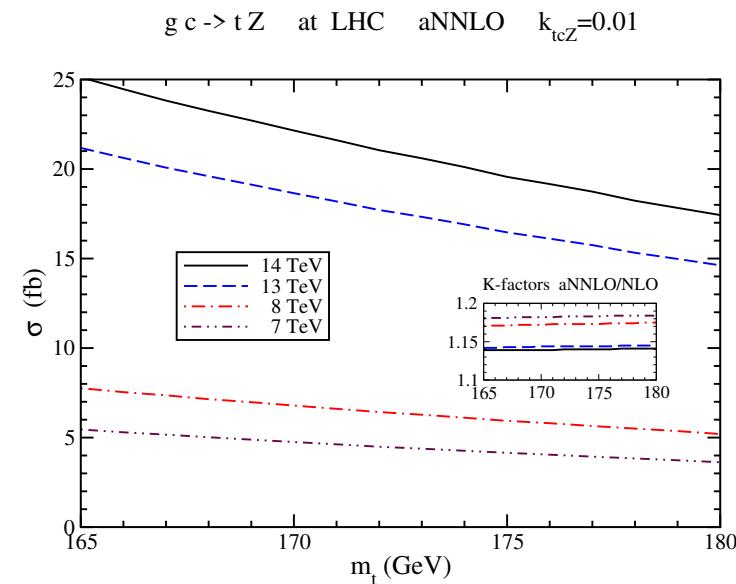
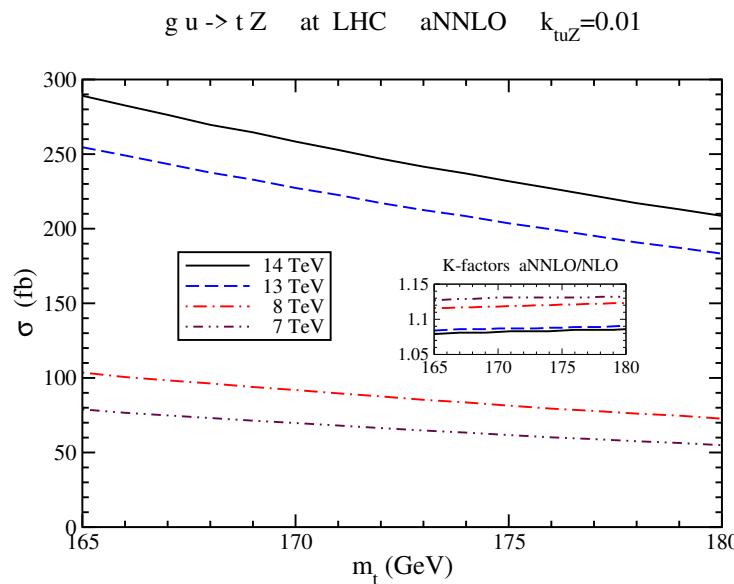
Normalized top quark distributions in tW^- production



tZ production via anomalous couplings

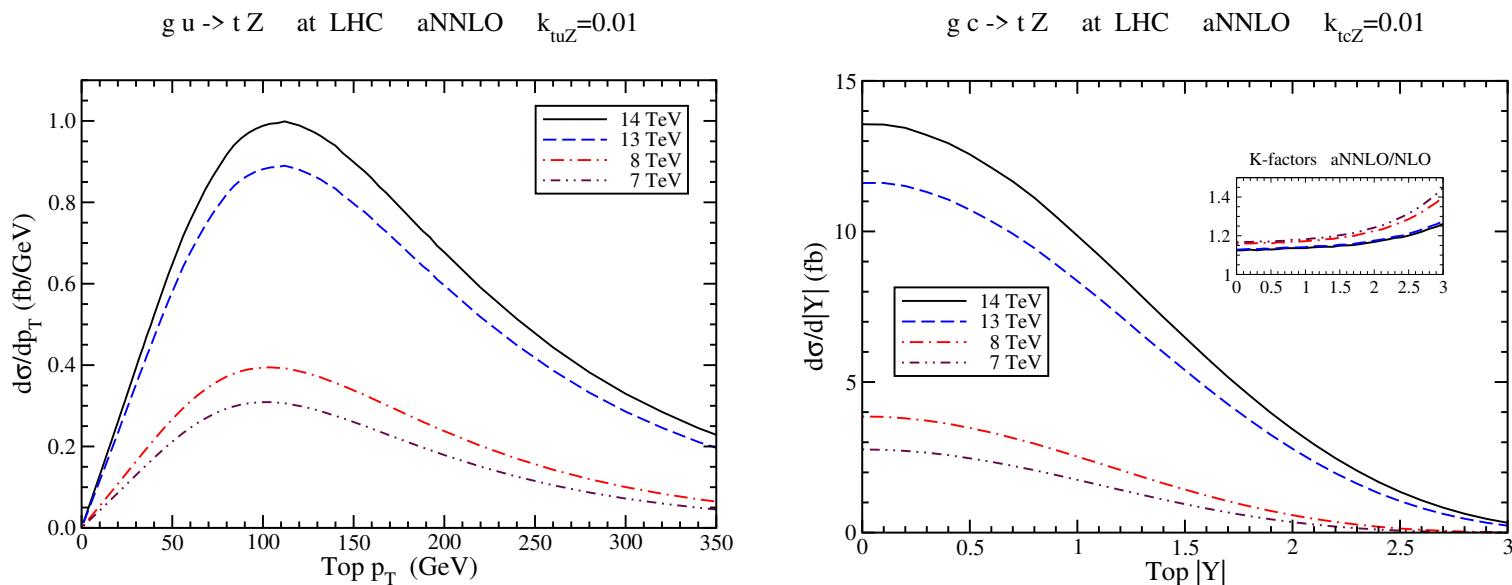


$$\Delta\mathcal{L}^{eff} = \frac{1}{\Lambda} \kappa_{tqZ} e \bar{t}(i/2)(\gamma_\mu \gamma_\nu - \gamma_\nu \gamma_\mu) q F_Z^{\mu\nu} + h.c.$$



tZ production via anomalous couplings

Top-quark p_T and rapidity distributions



Summary

- cross sections and distributions for single-top production
- soft-gluon corrections
- t -channel at aNNLO
- s -channel at aNNLO
- tW production at aN³LO
- excellent agreement with LHC and Tevatron data
- tZ production via anomalous couplings
- high-order corrections are very significant