

aCORN

A Precision Measurement of the Electron-Antineutrino
Correlation "a" in Neutron Beta Decay



F. E. Wietfeldt



Physics Department
Tulane University
New Orleans, LA



aCORN Collaboration

G. Darius, C. DeAngelis, T. Hassan, F. E. Wietfeldt
Tulane University

M. S. Dewey, M. P. Mendenhall, J. S. Nico, H. Park (visitor)
National Institute of Standards and Technology

G. Noid, E. Stephenson
Indiana University

B. Collett, G. L. Jones
Hamilton College

A. Komives
DePauw University

graduate student

Neutron Decay Parameters

Phenomenological ($J = 1/2 \rightarrow J = 1/2$) beta decay formula [Jackson, Treiman, Wyld, 1957] :

$$dW \propto \frac{1}{\tau} F(E_e) \left[1 + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + b \frac{m_e}{E_e} + A \frac{\vec{\sigma}_n \cdot \vec{p}_e}{E_e} + B \frac{\vec{\sigma}_n \cdot \vec{p}_\nu}{E_\nu} + D \frac{\vec{\sigma}_n \cdot (\vec{p}_e \times \vec{p}_\nu)}{E_e E_\nu} \right]$$

For allowed beta decay, neglecting recoil order terms, the standard electroweak model (Weinberg, Glashow, Salam, et al.) predicts:

$$a = \frac{1 - \lambda^2}{1 + 3\lambda^2} \quad b = 0 \quad A = -2 \frac{\lambda^2 + \text{Re}(\lambda)}{1 + 3\lambda^2} \quad B = 2 \frac{\lambda^2 - \text{Re}(\lambda)}{1 + 3\lambda^2}$$

$$D = 2 \frac{\text{Im}(\lambda)}{1 + 3\lambda^2} \approx 0 \quad \tau \propto \frac{1}{g_v^2 + 3g_A^2} \quad \text{where} \quad \lambda \equiv \frac{g_A}{g_v}$$

Why do we measure neutron decay parameters?

Within Standard Model: Get G_A , G_V

Beyond Standard Model:

Mostovoy Parameters, model-independent consistency test of SM:

predicted	actual	
$F_1 = 1 + A - B - a = 0$	$F_1 = 0.0025 \pm 0.0064$	uncertainties dominated by "a"
$F_2 = aB - A - A^2 = 0$	$F_2 = 0.0034 \pm 0.0050$	

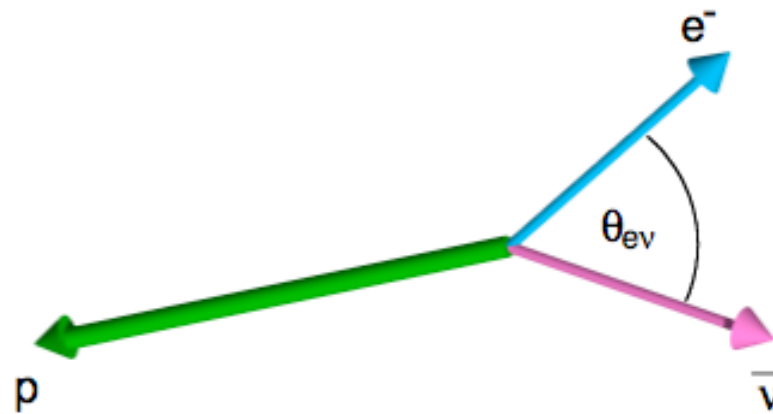
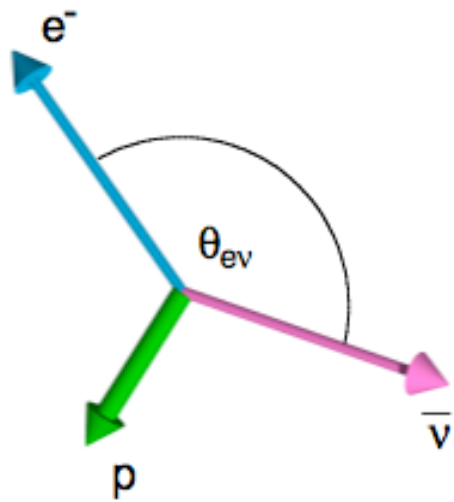
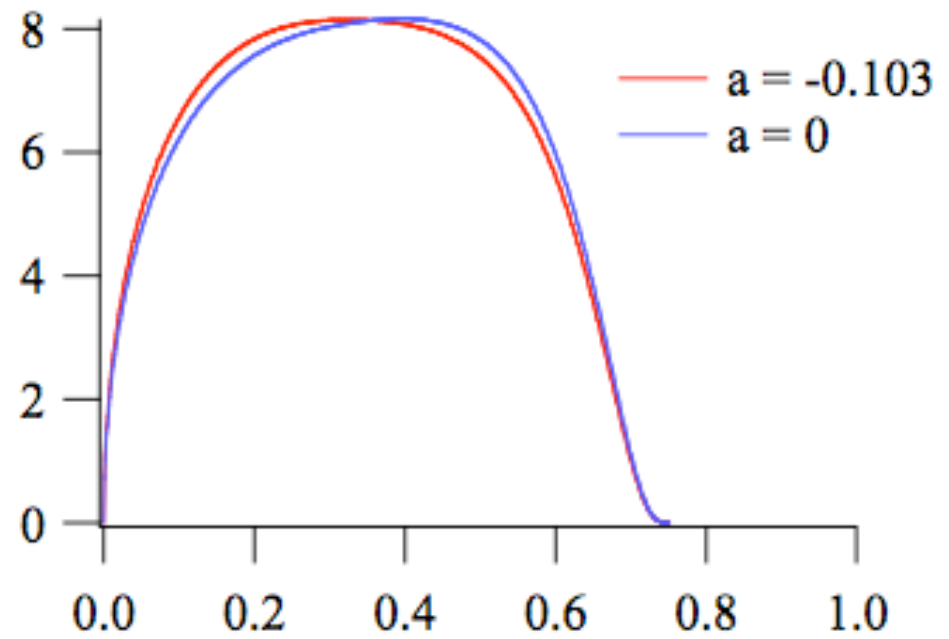
Precise comparisons of a , b , A , B , D are sensitive to:

- scalar and tensor weak currents
- right handed weak currents
- new CP violation
- CVC violation and second-class currents (Gardner and Zhang, 2000)
- SUSY (Profumo, Ramsey-Musolf, and Tulin, 2007)

Standard method for measuring the e- $\bar{\nu}$ correlation:

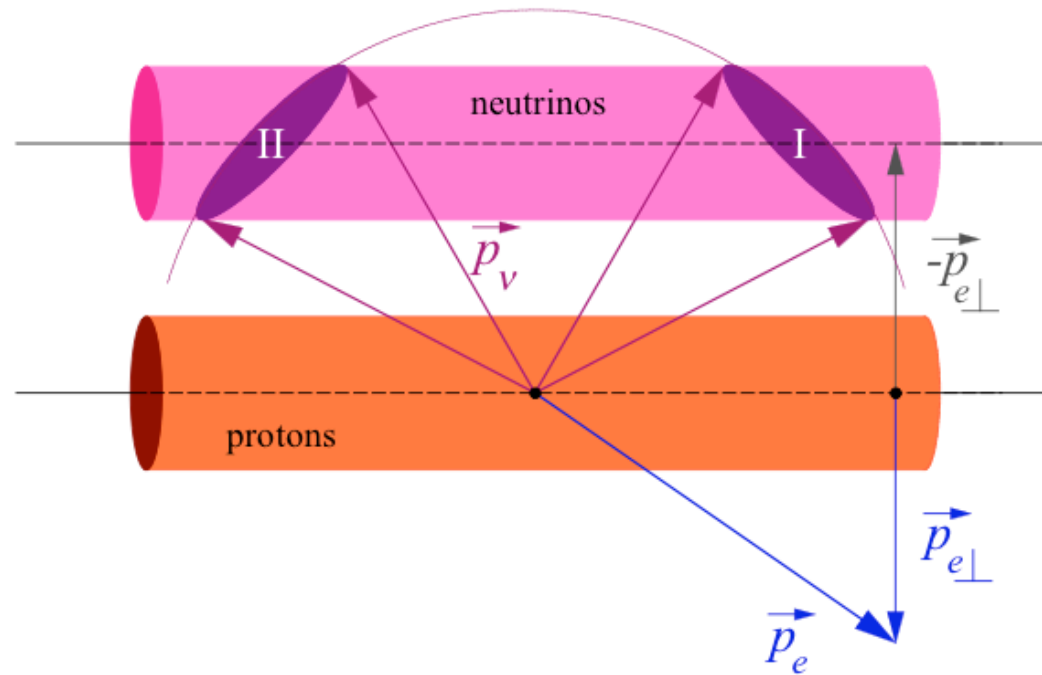
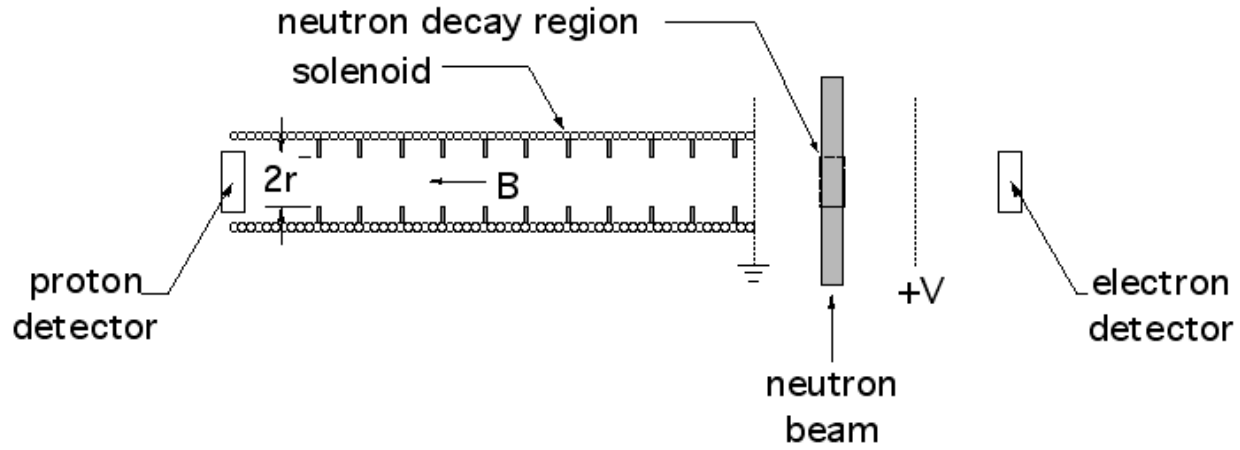
recoil energy spectrum

statistically most advantageous

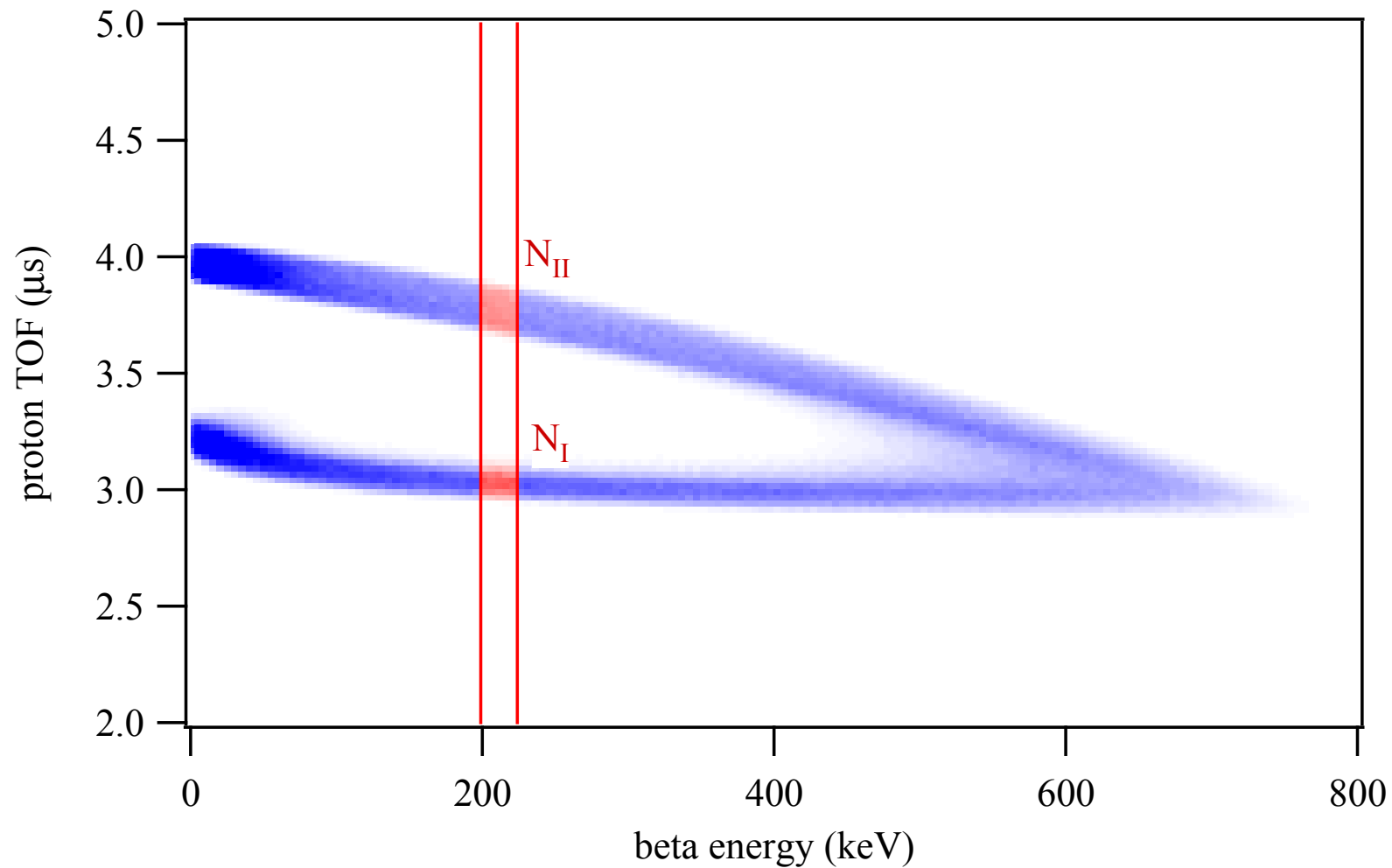


A Novel Method to Measure α

(Yerozolimsky and Mostovoy, 1996)



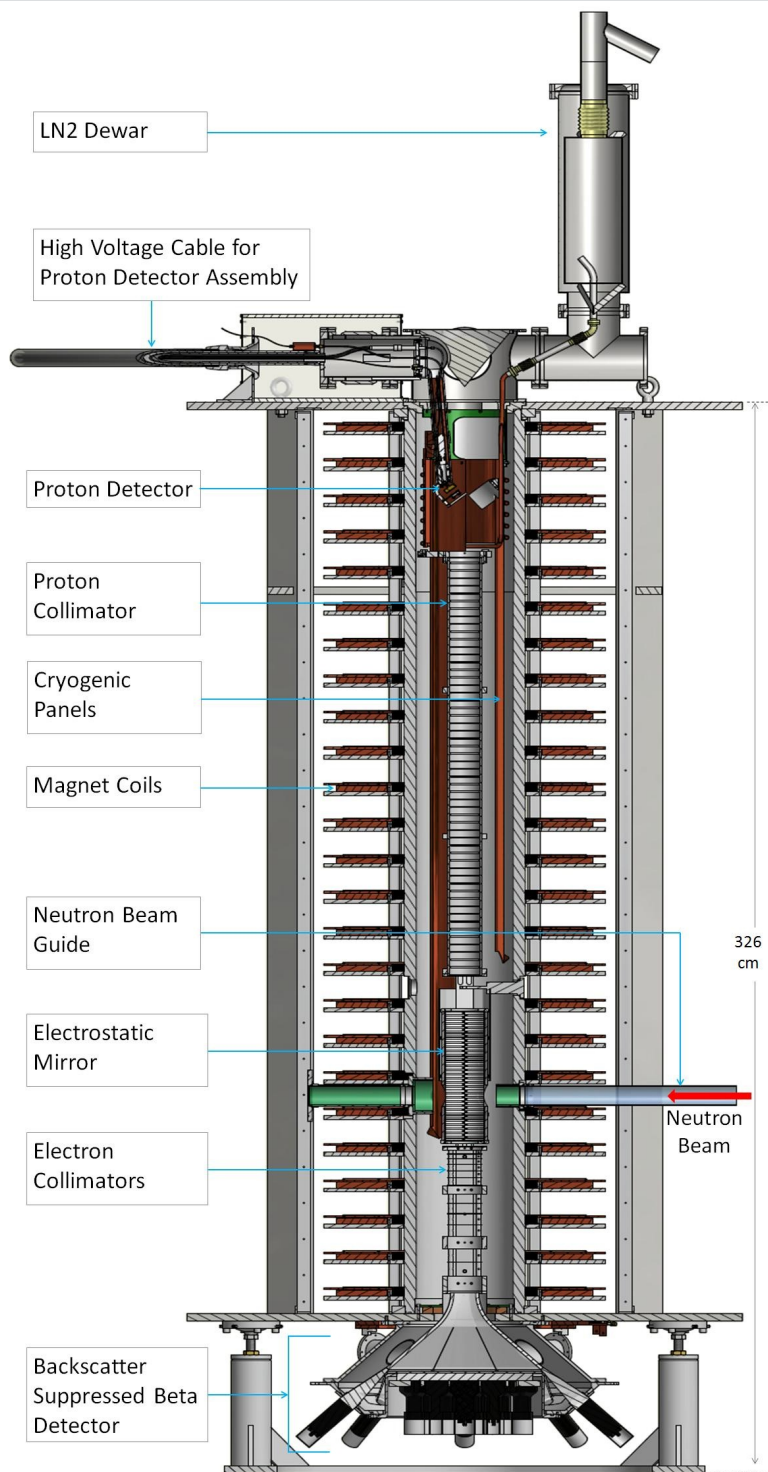
We separate groups I and II by beta energy and proton time-of-flight (TOF)



$$X(E) = \frac{N_I - N_{II}}{N_I + N_{II}} = a f_a(E)$$

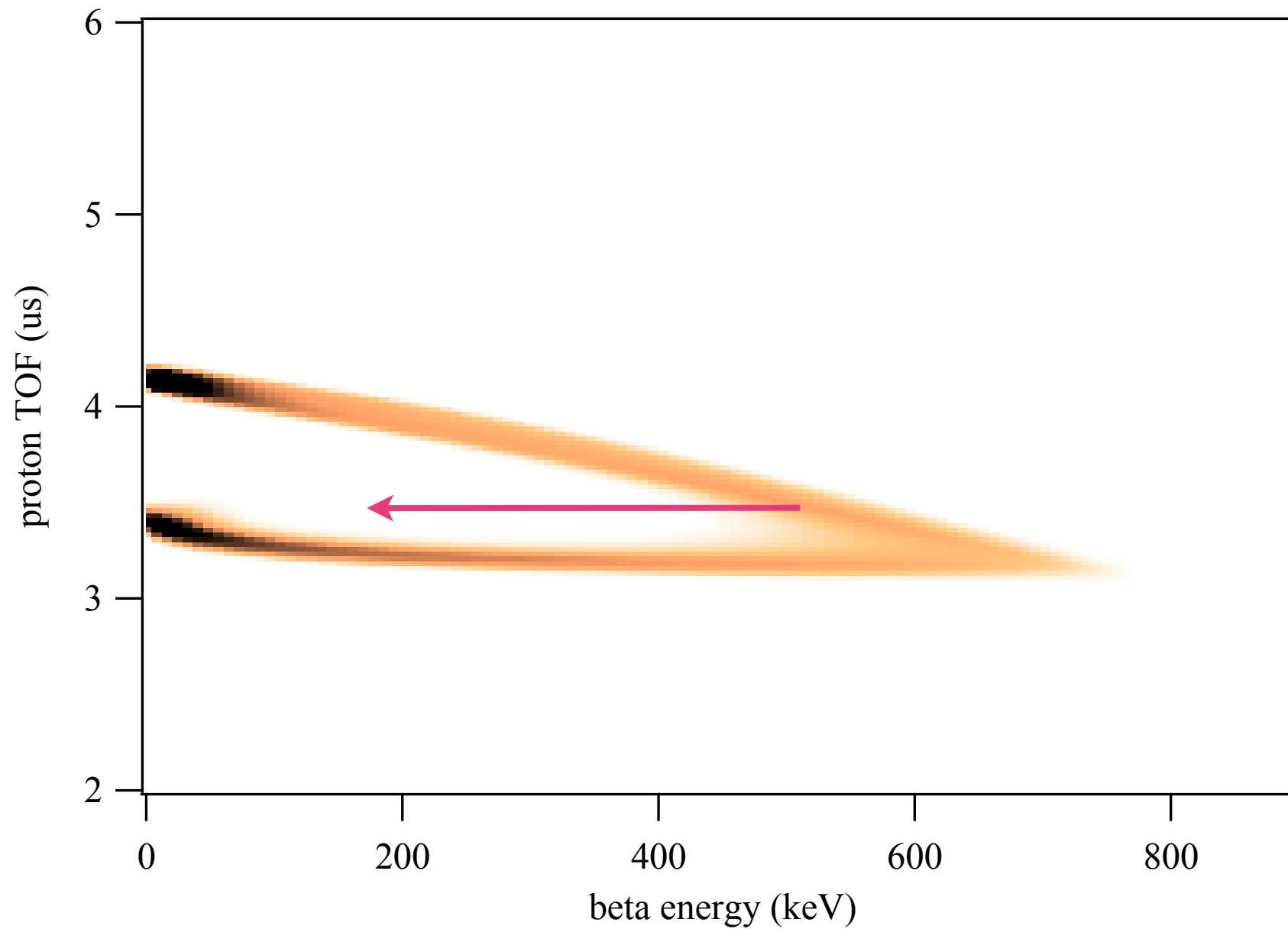
wishbone asymmetry

geometric function



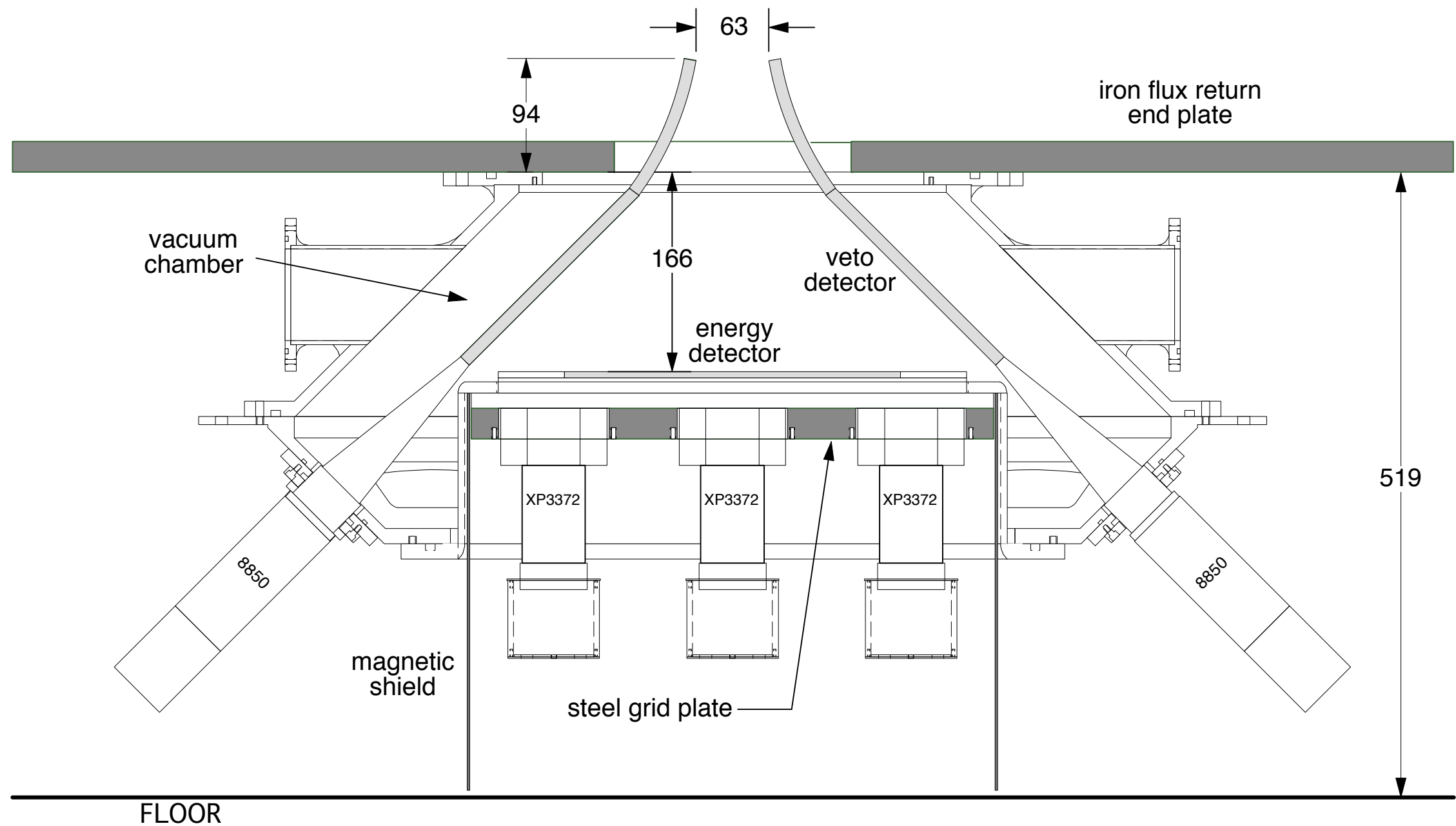
aCORN

Electron backscatter

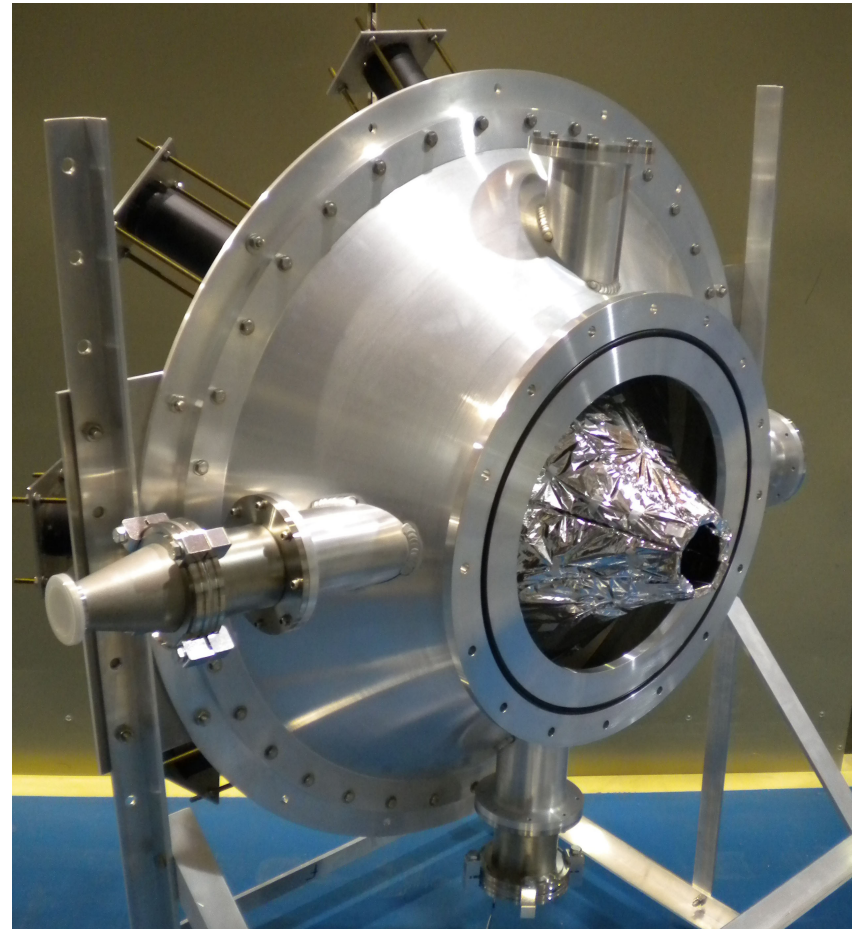
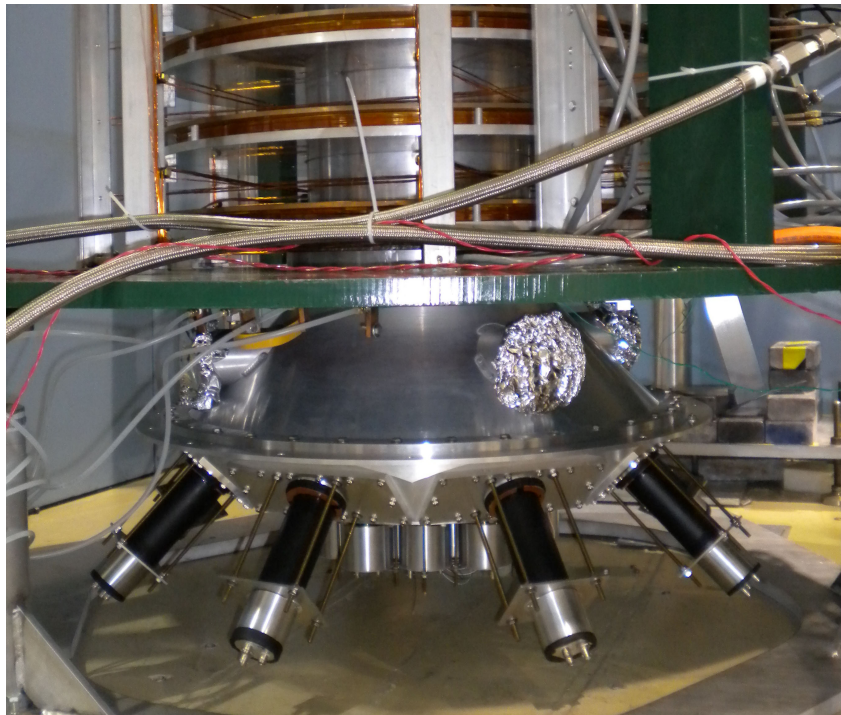
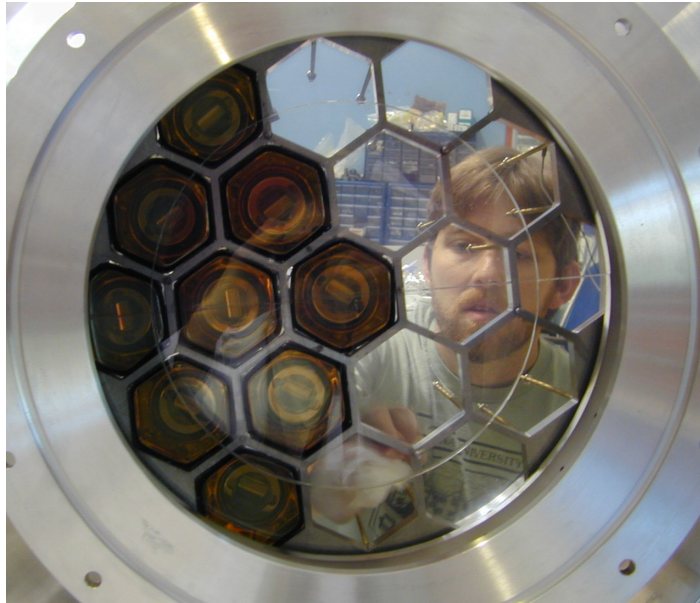


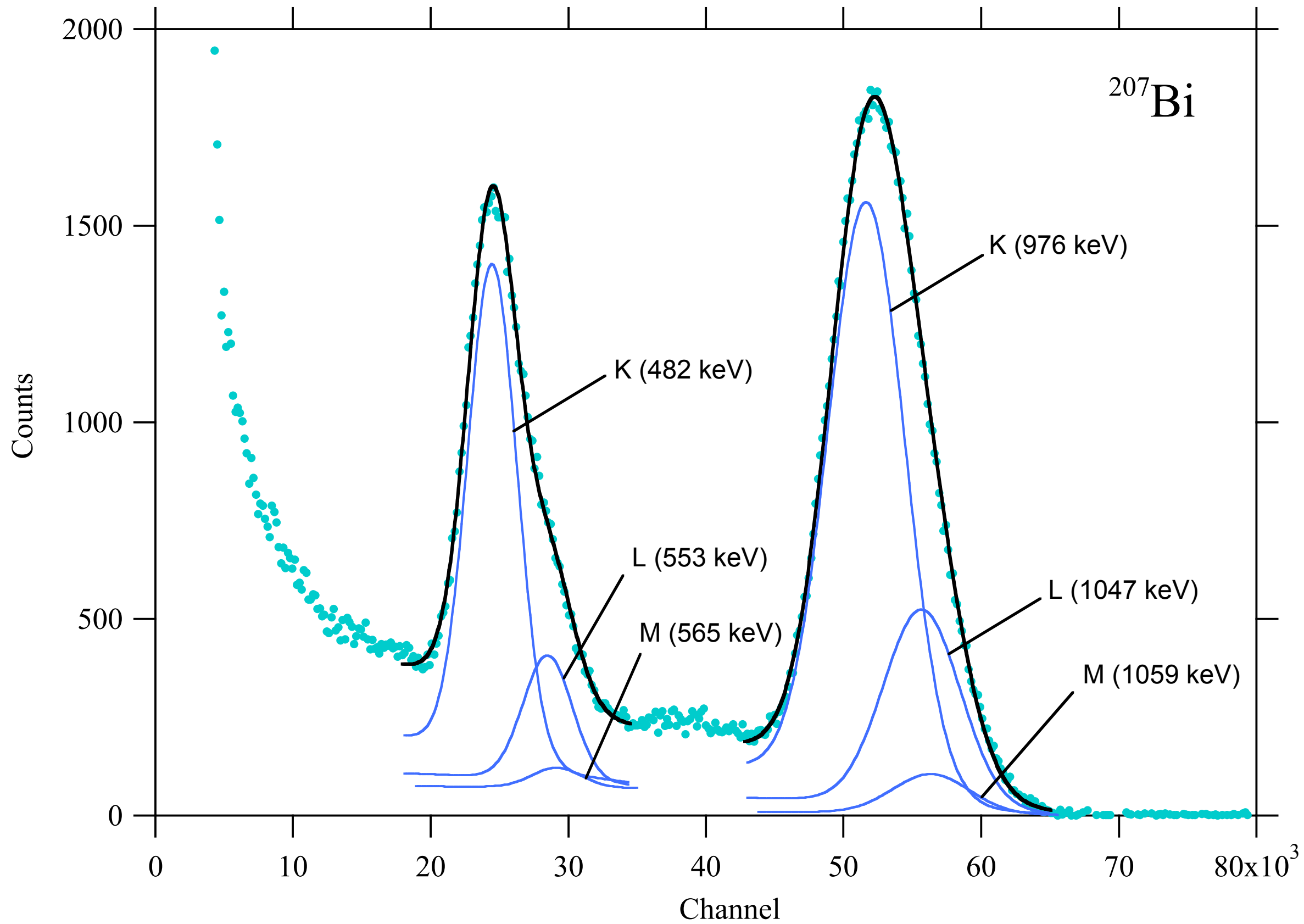
Electron backscatter will cause electrons to appear at a lower, incorrect energy, filling in the gap between the branches.

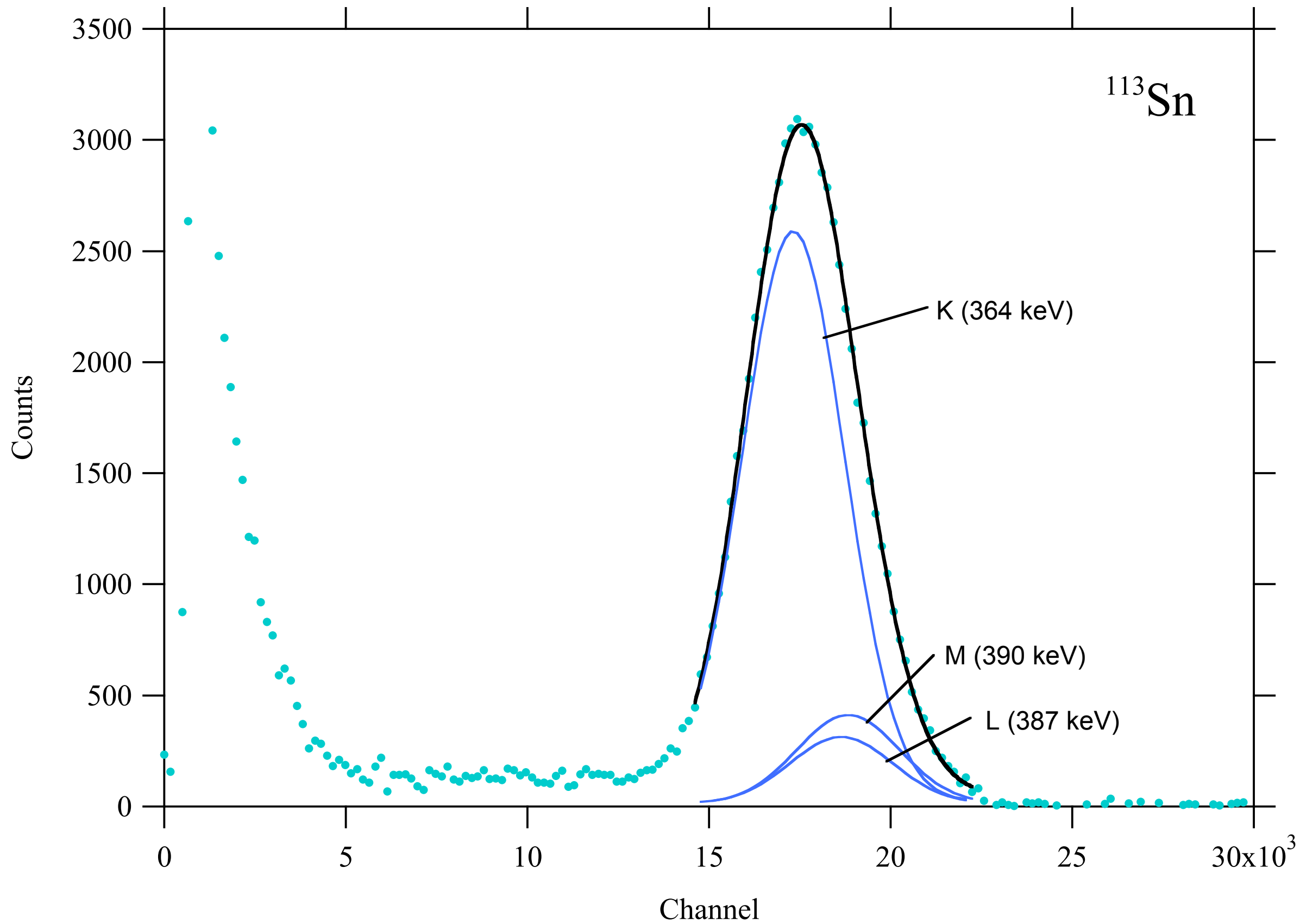
aCORN backscatter suppressed beta spectrometer



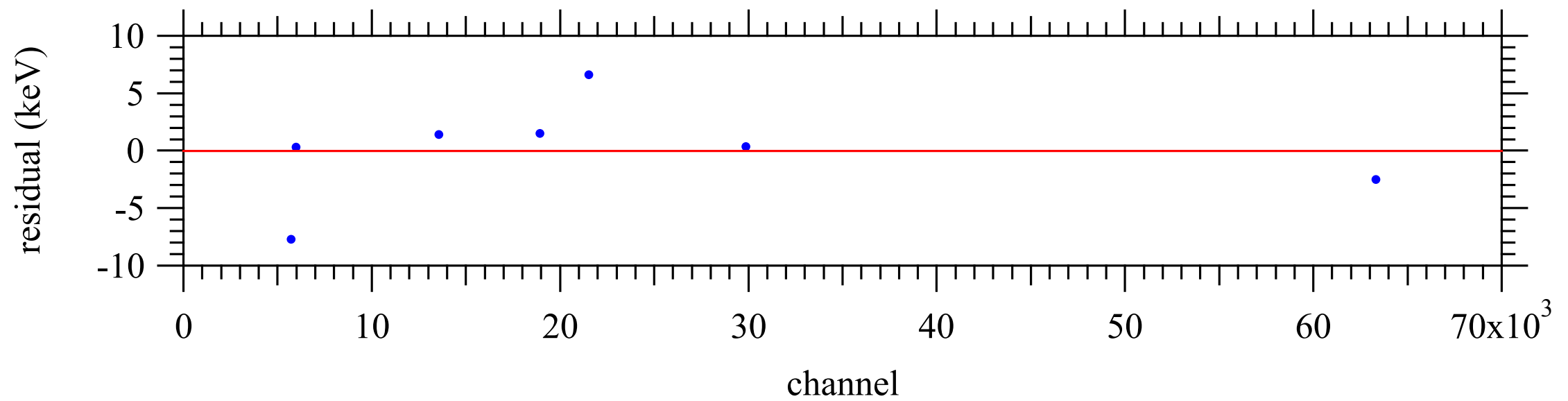
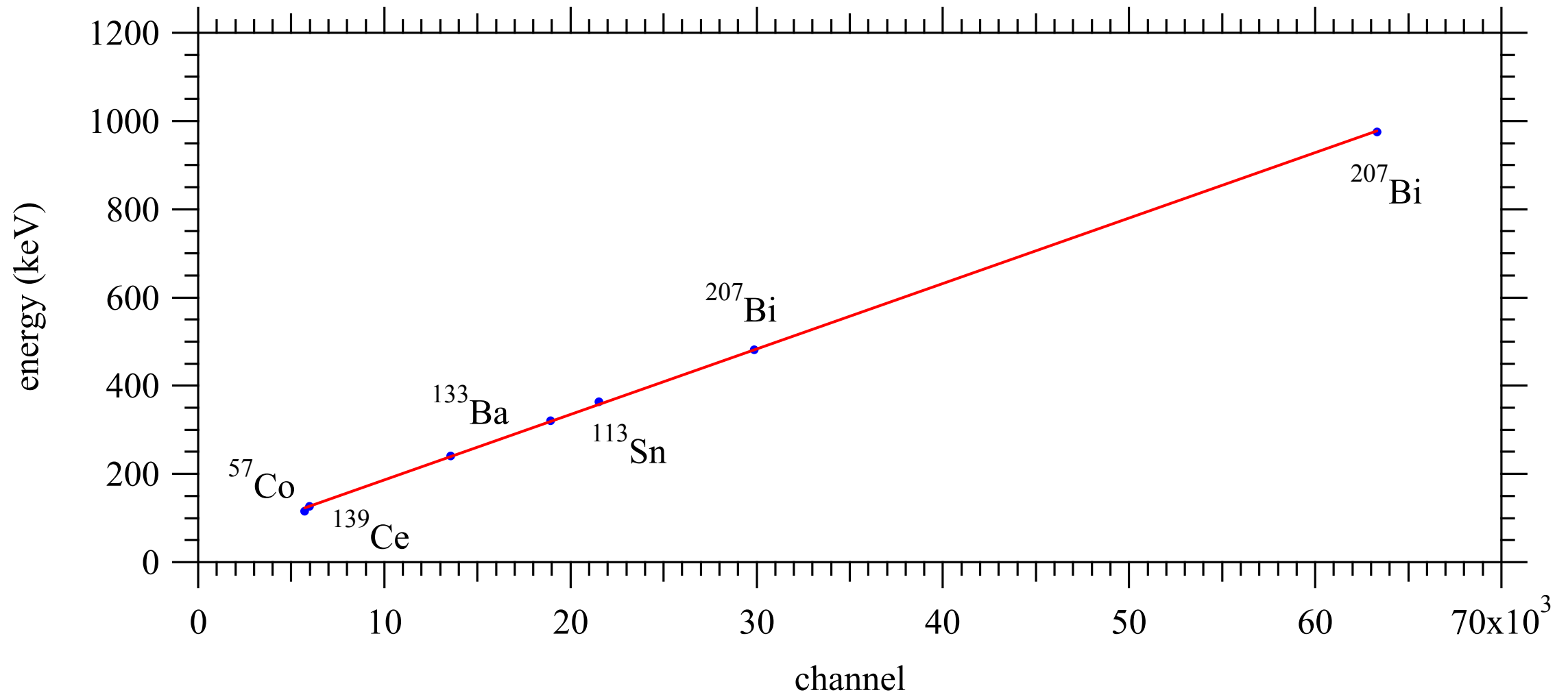
aCORN Beta Spectrometer



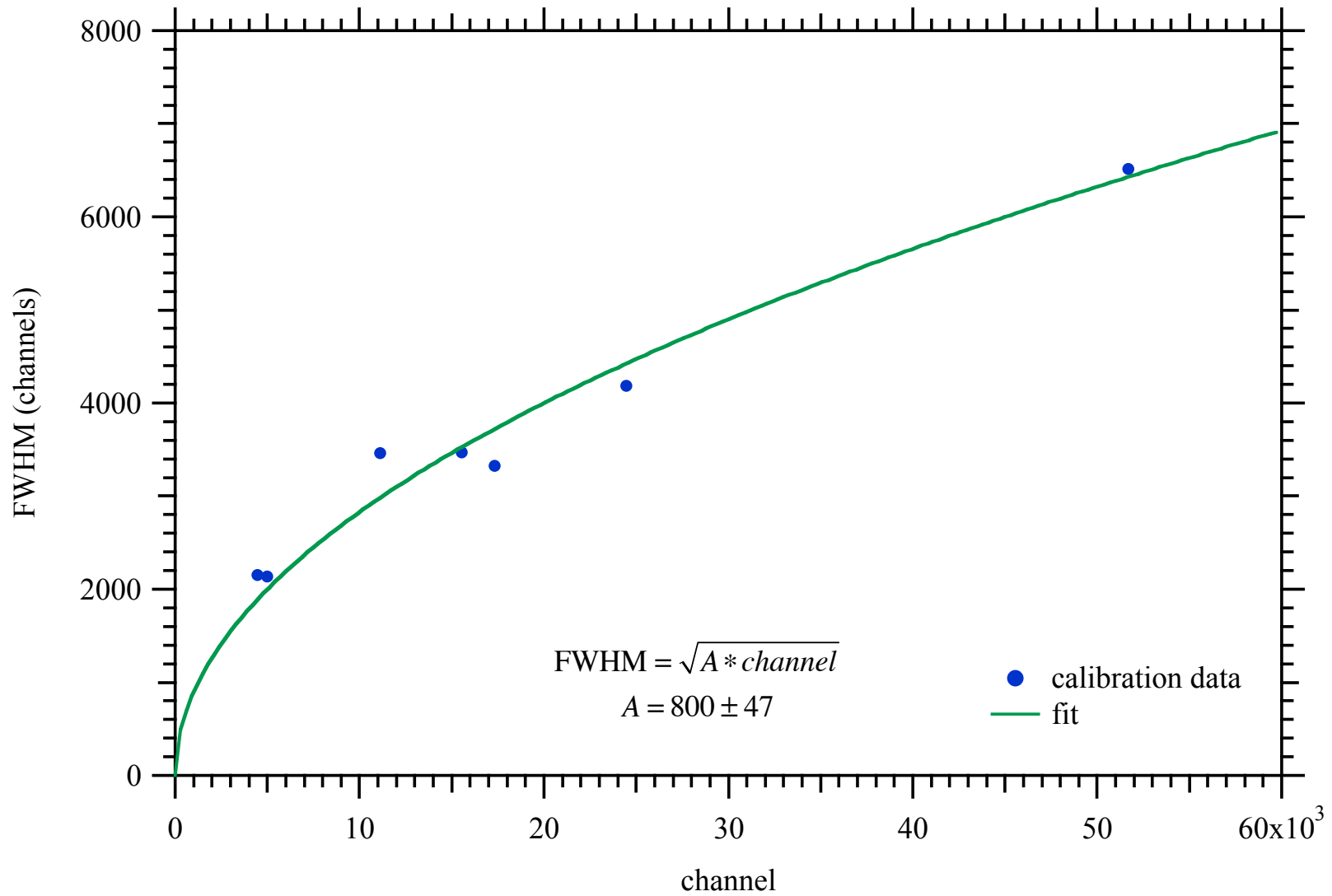




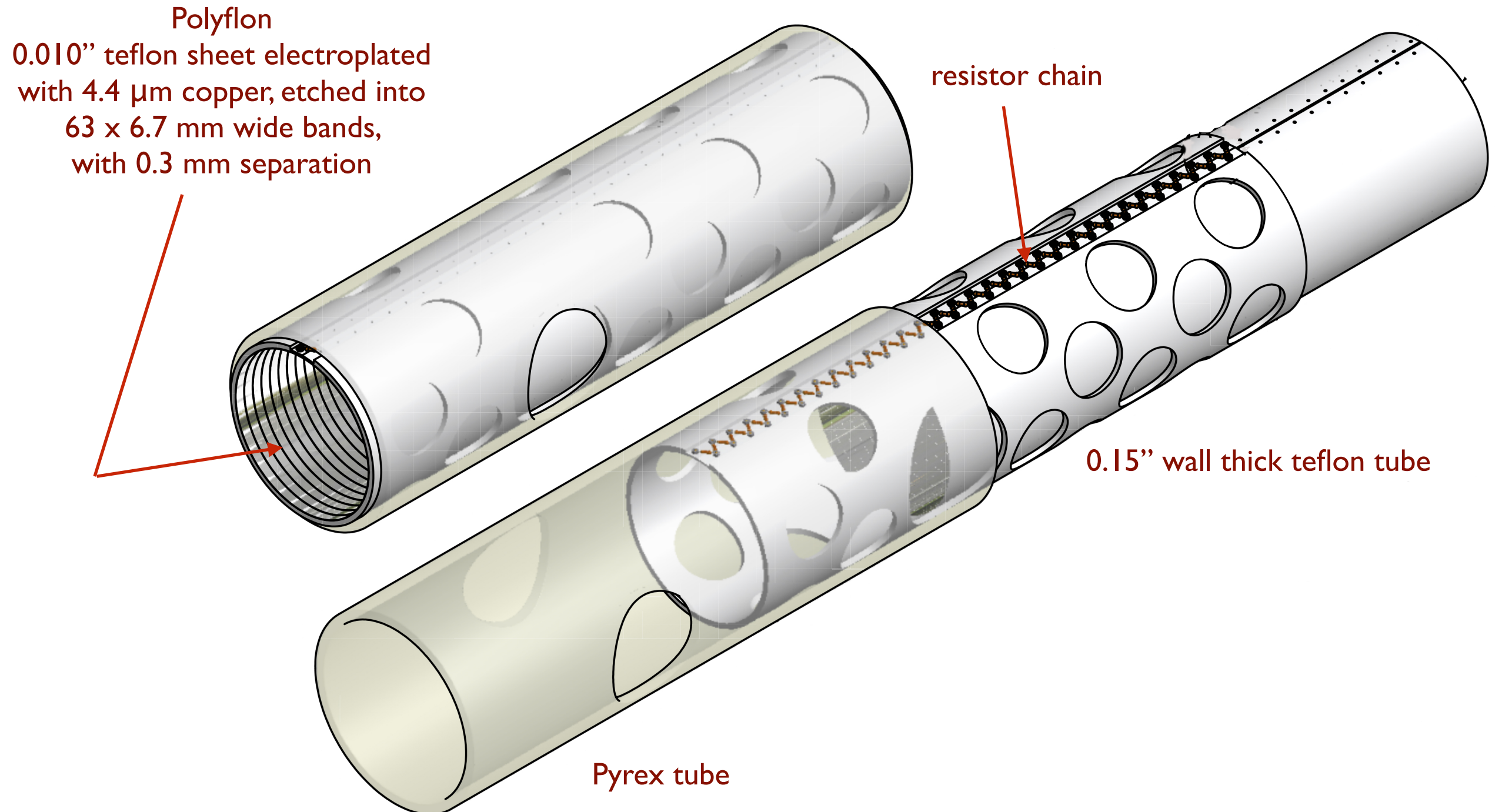
Beta Spectrometer Energy Response



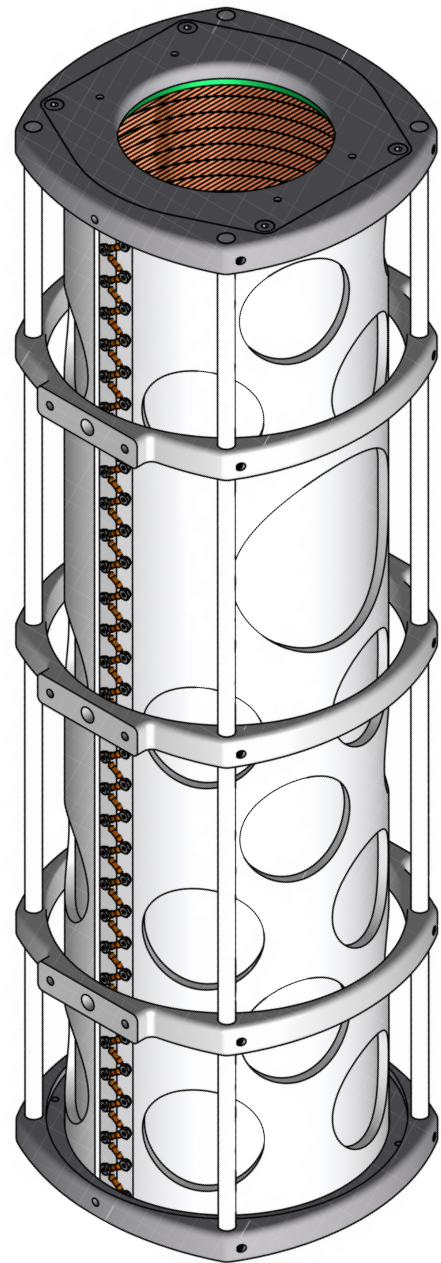
Beta Spectrometer Energy Resolution (FWHM)



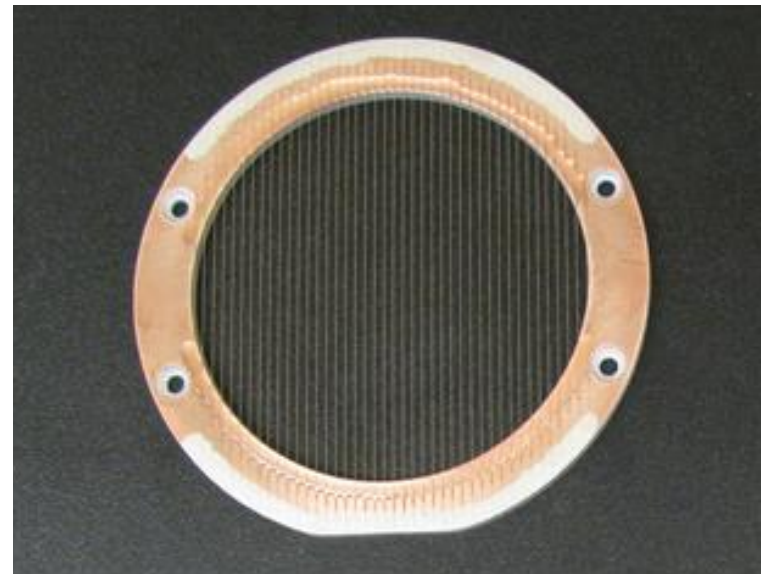
Electrostatic mirror



Electrostatic mirror



proton collimator →



ground grid

100 μm gold-plated BeCu wire grid,
2-mm spacing

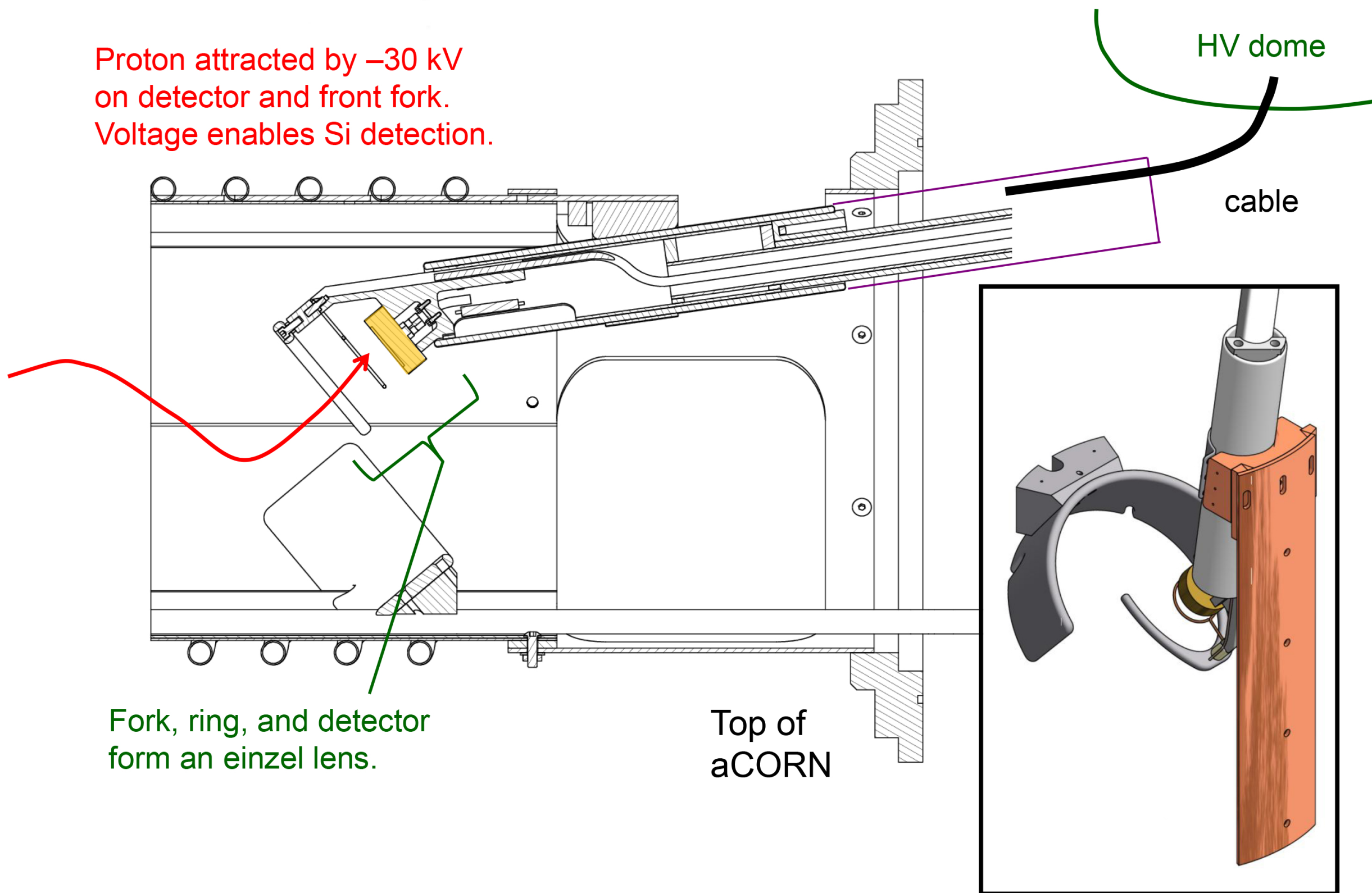
+3 kV grid

beta collimator →



Proton detector

Proton attracted by -30 kV
on detector and front fork.
Voltage enables Si detection.



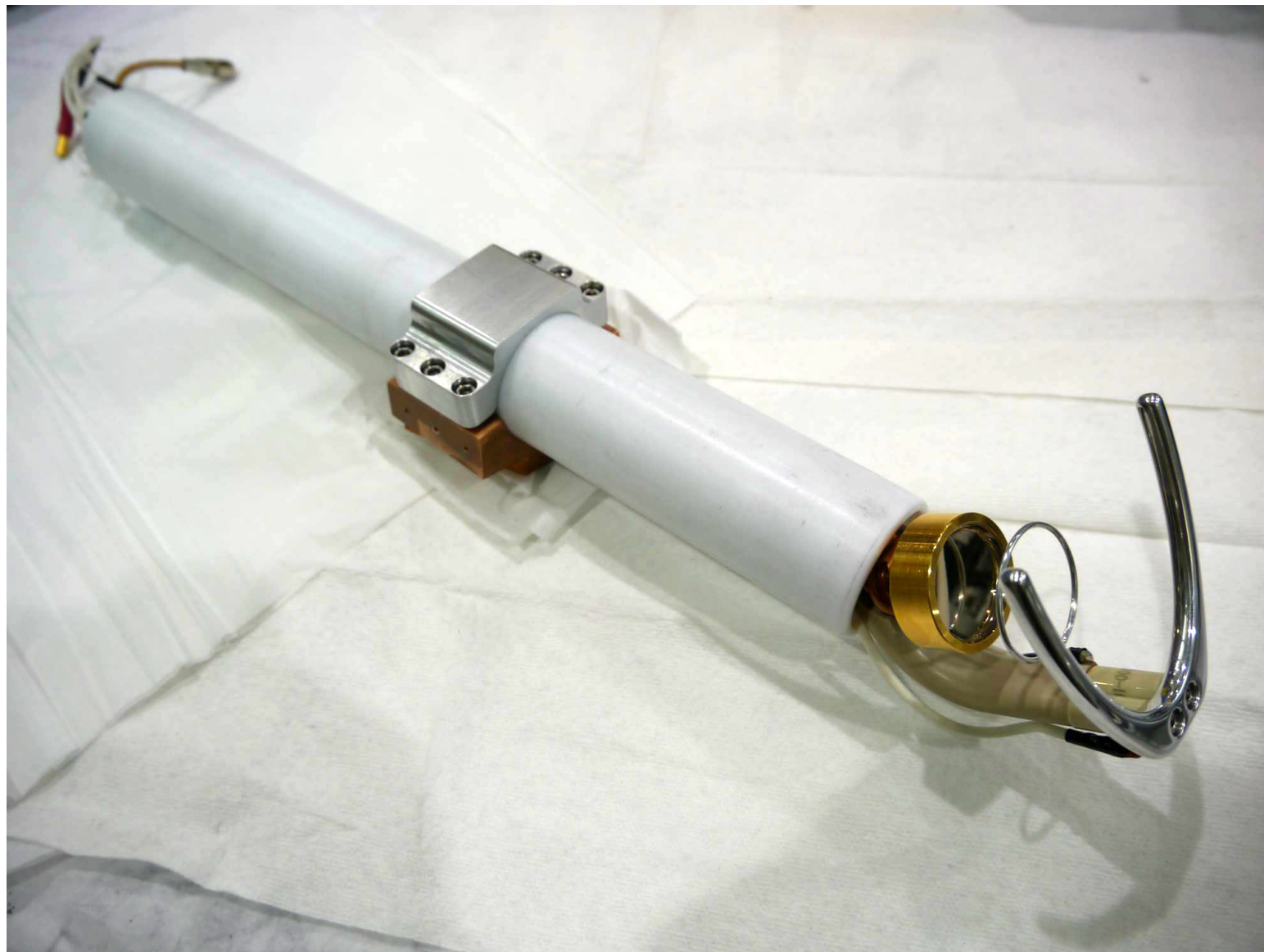
Fork, ring, and detector
form an einzel lens.

Top of
aCORN

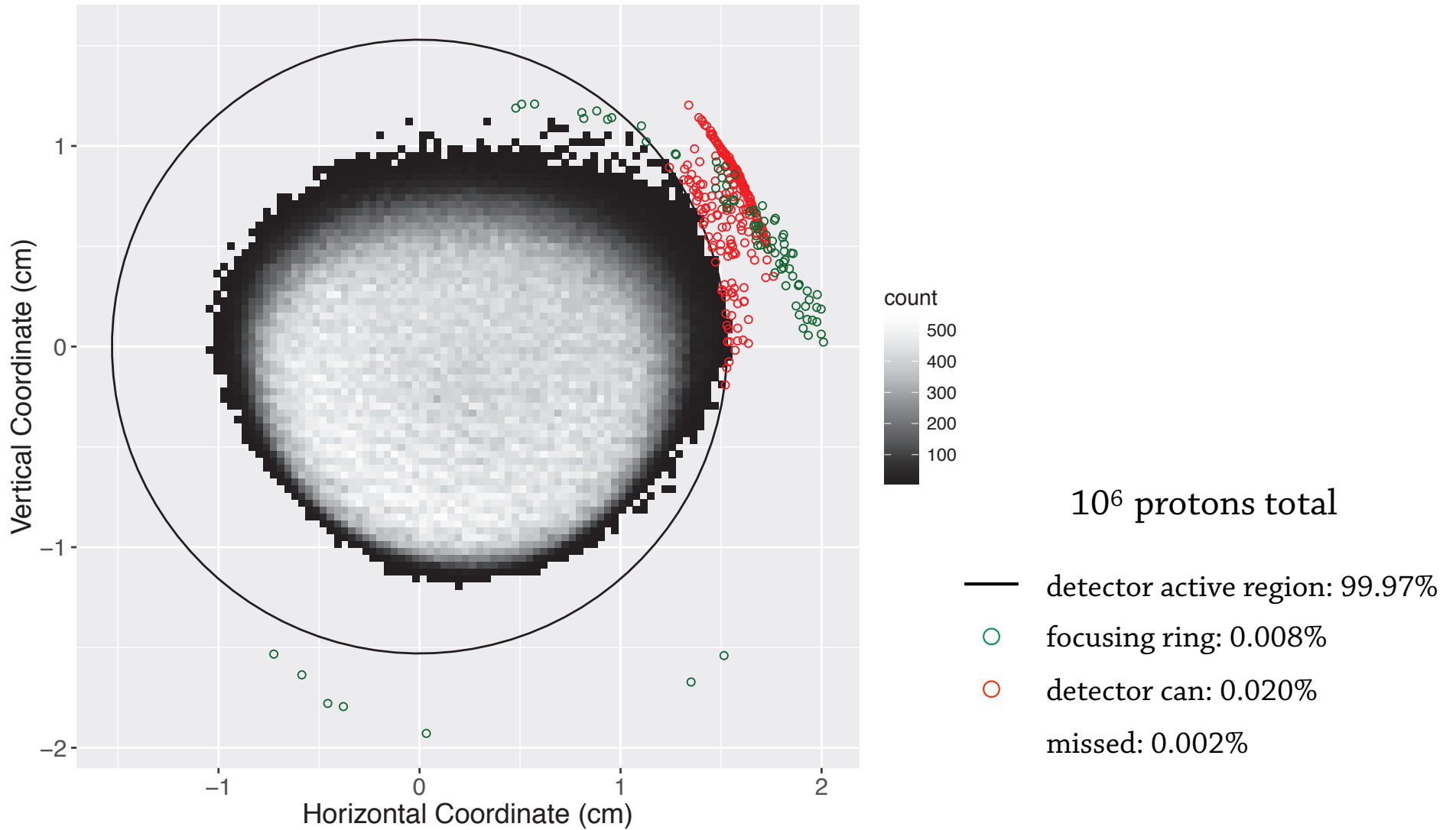
HV dome

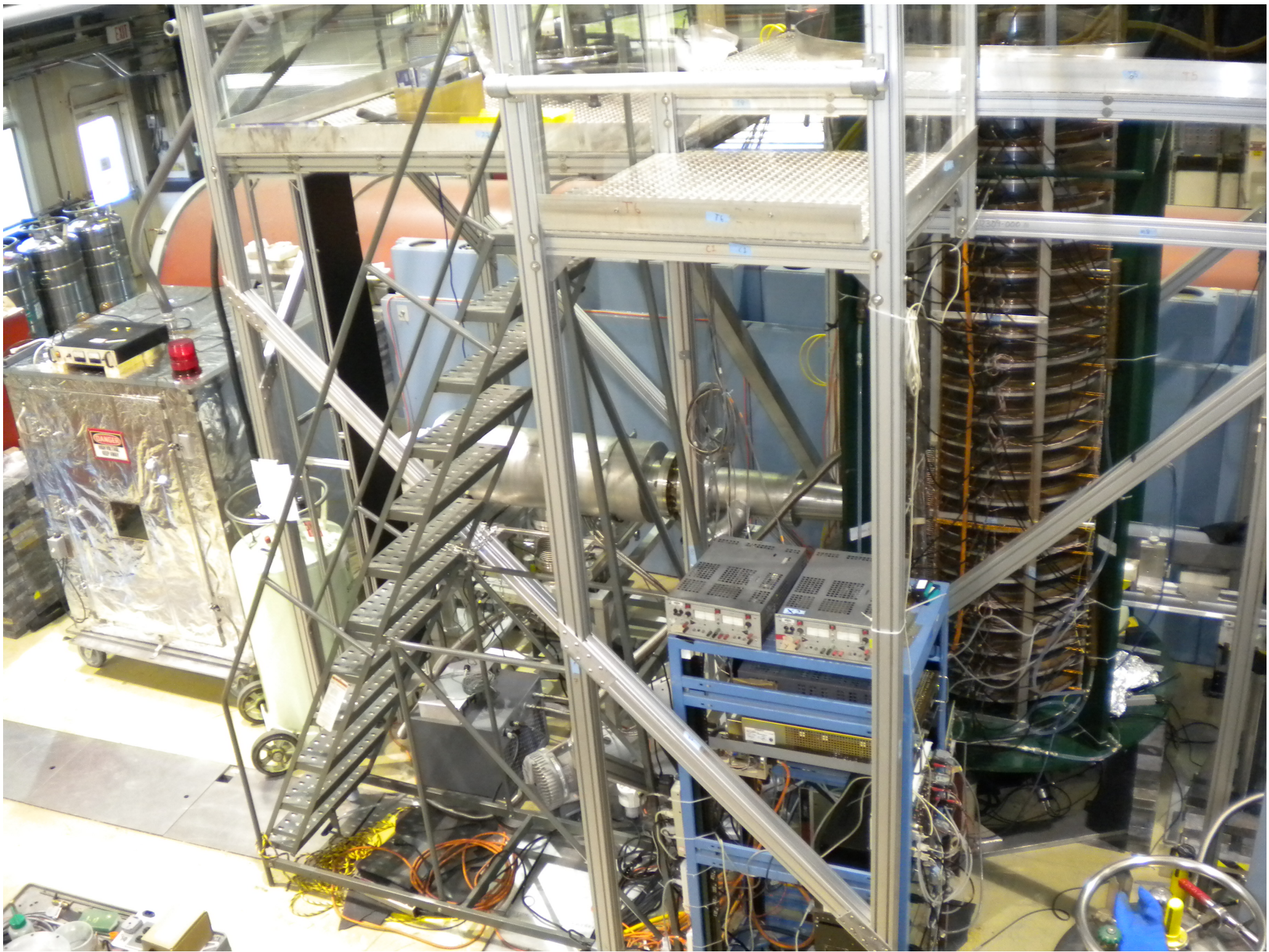
cable

aCORN proton detector

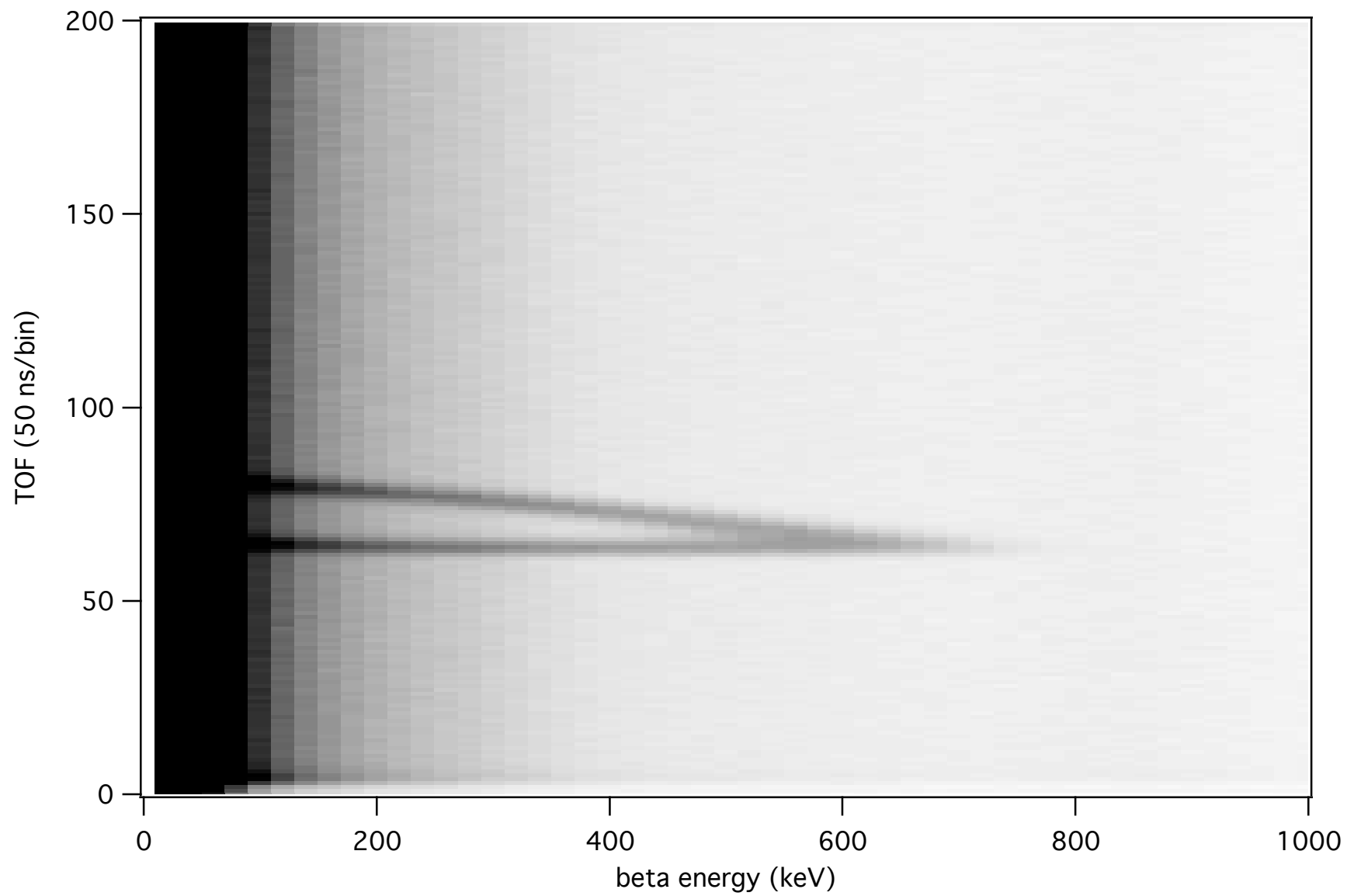


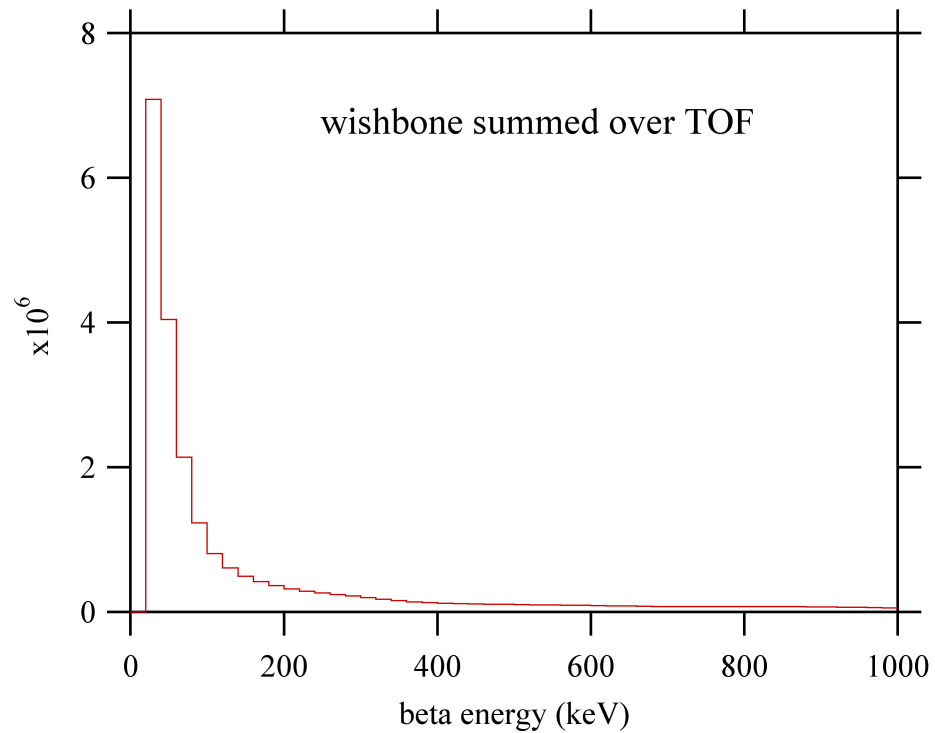
Proton Focusing Simulation



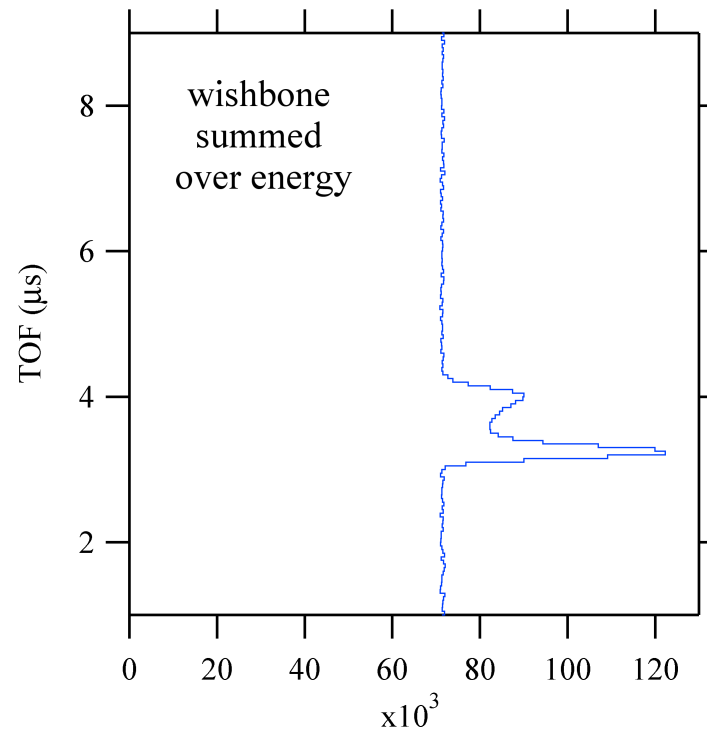
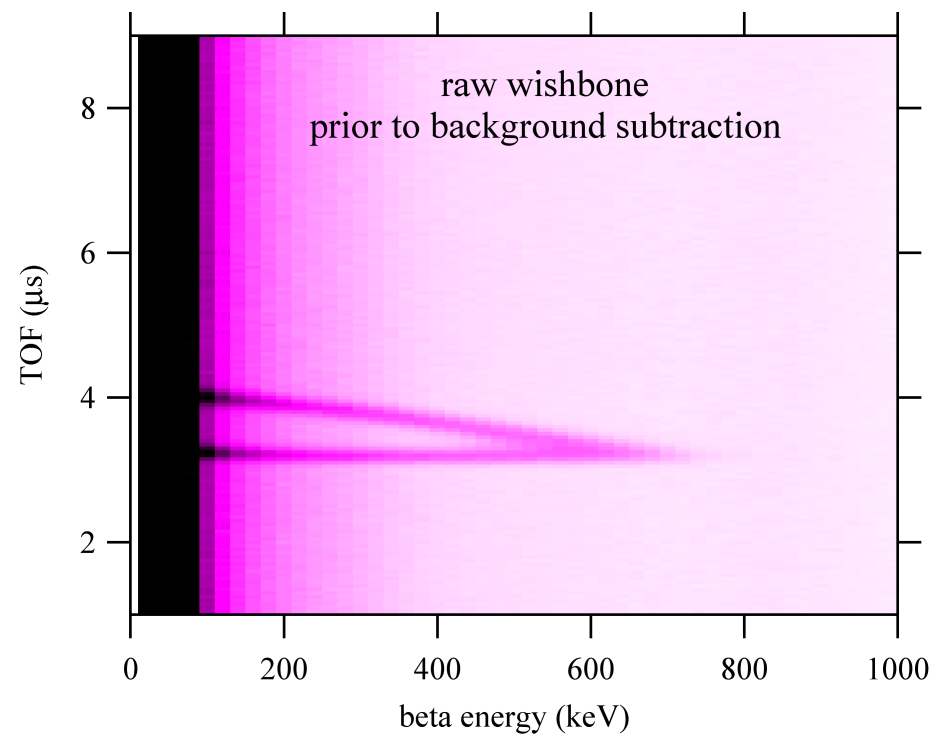


Typical Wishbone

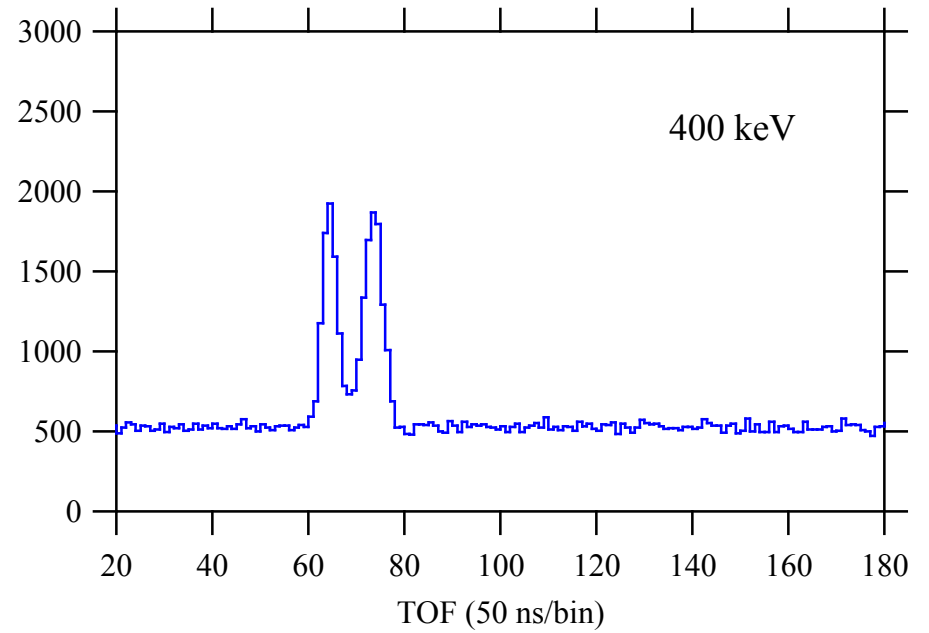
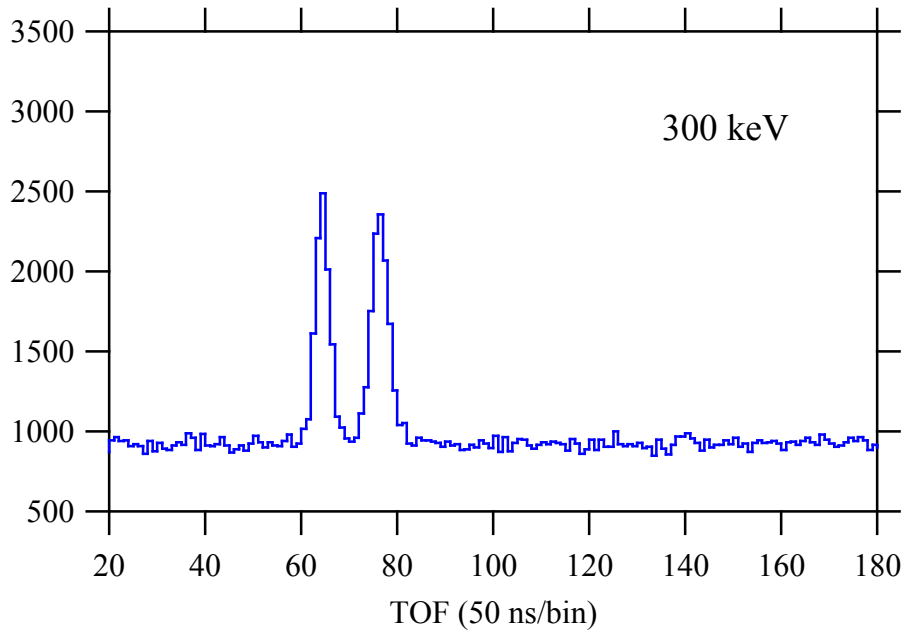
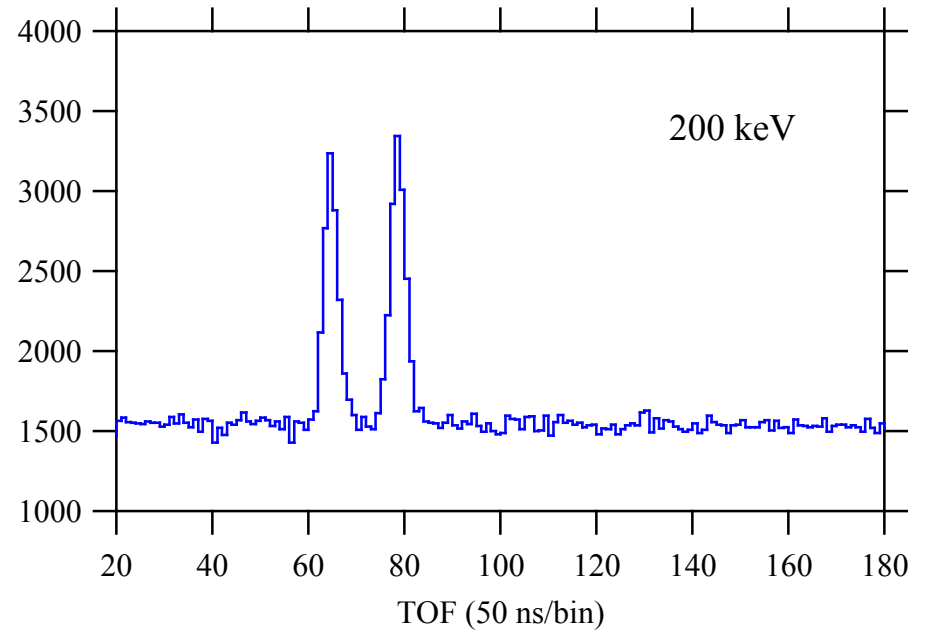
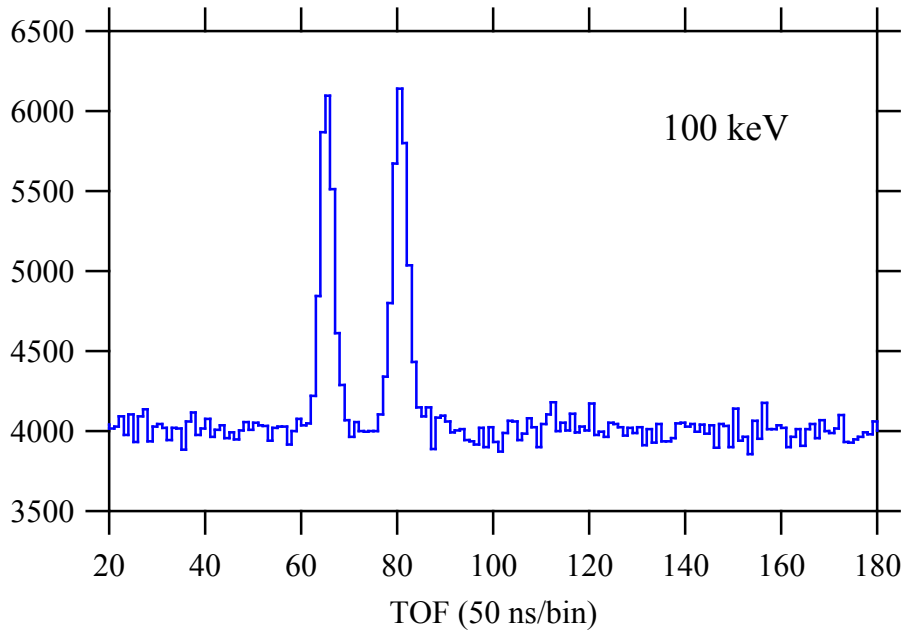


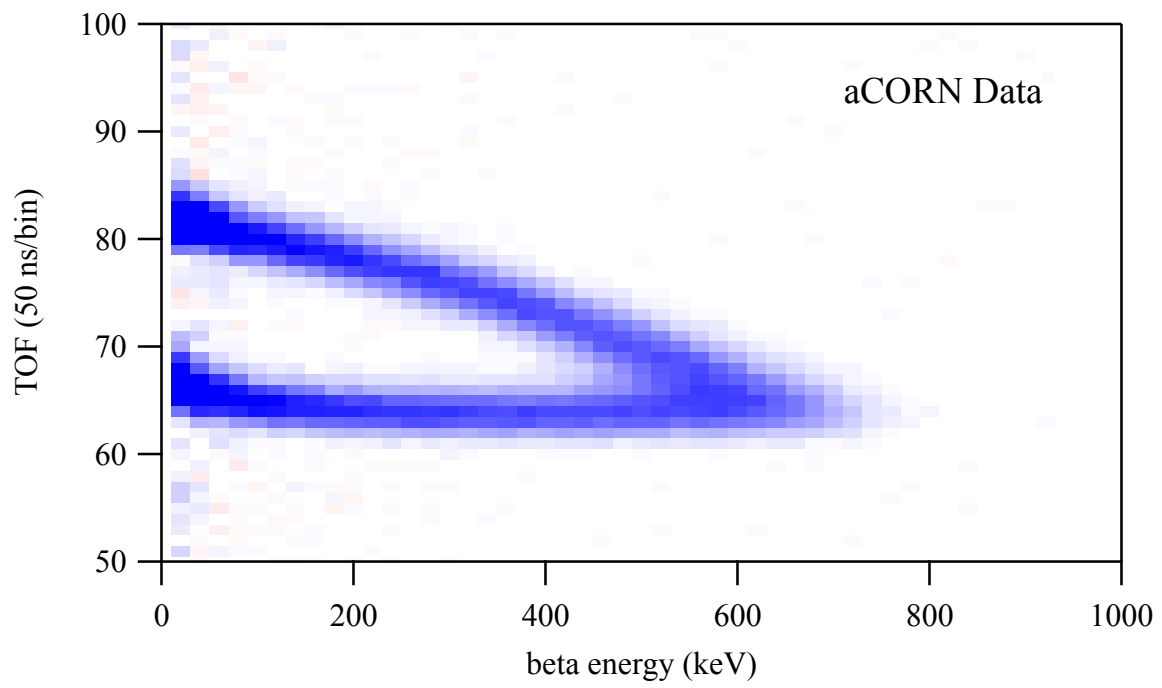


aCORN
background
(NG-6)

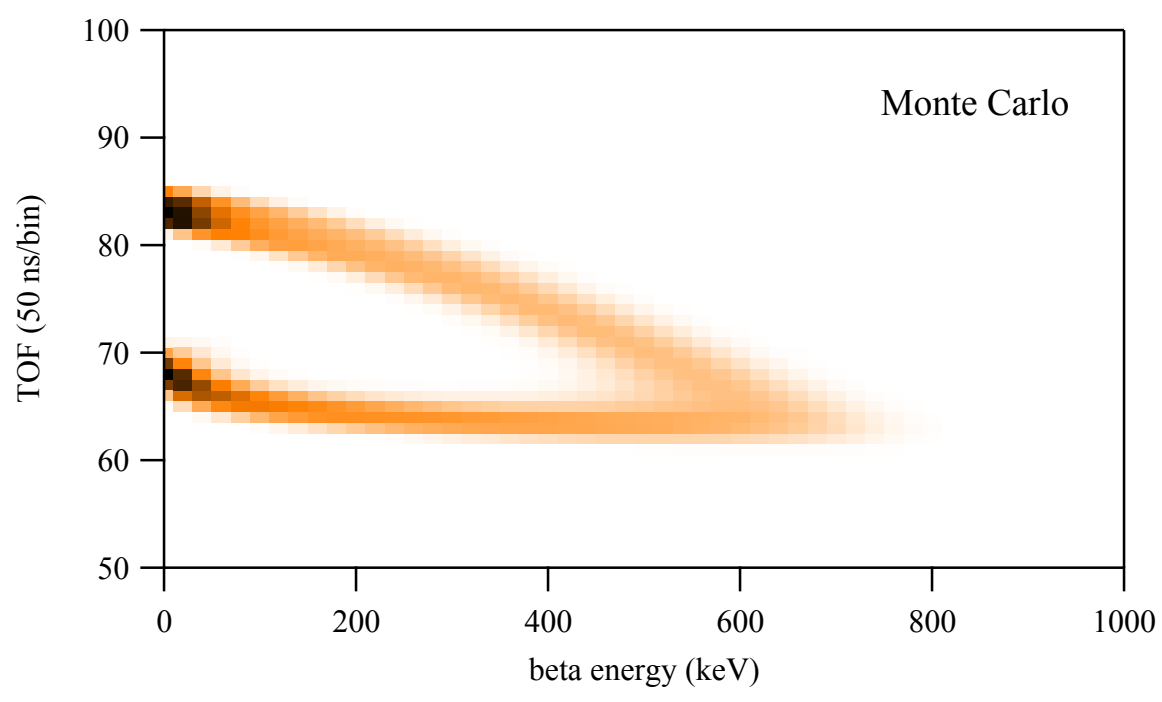


Wishbone Slices

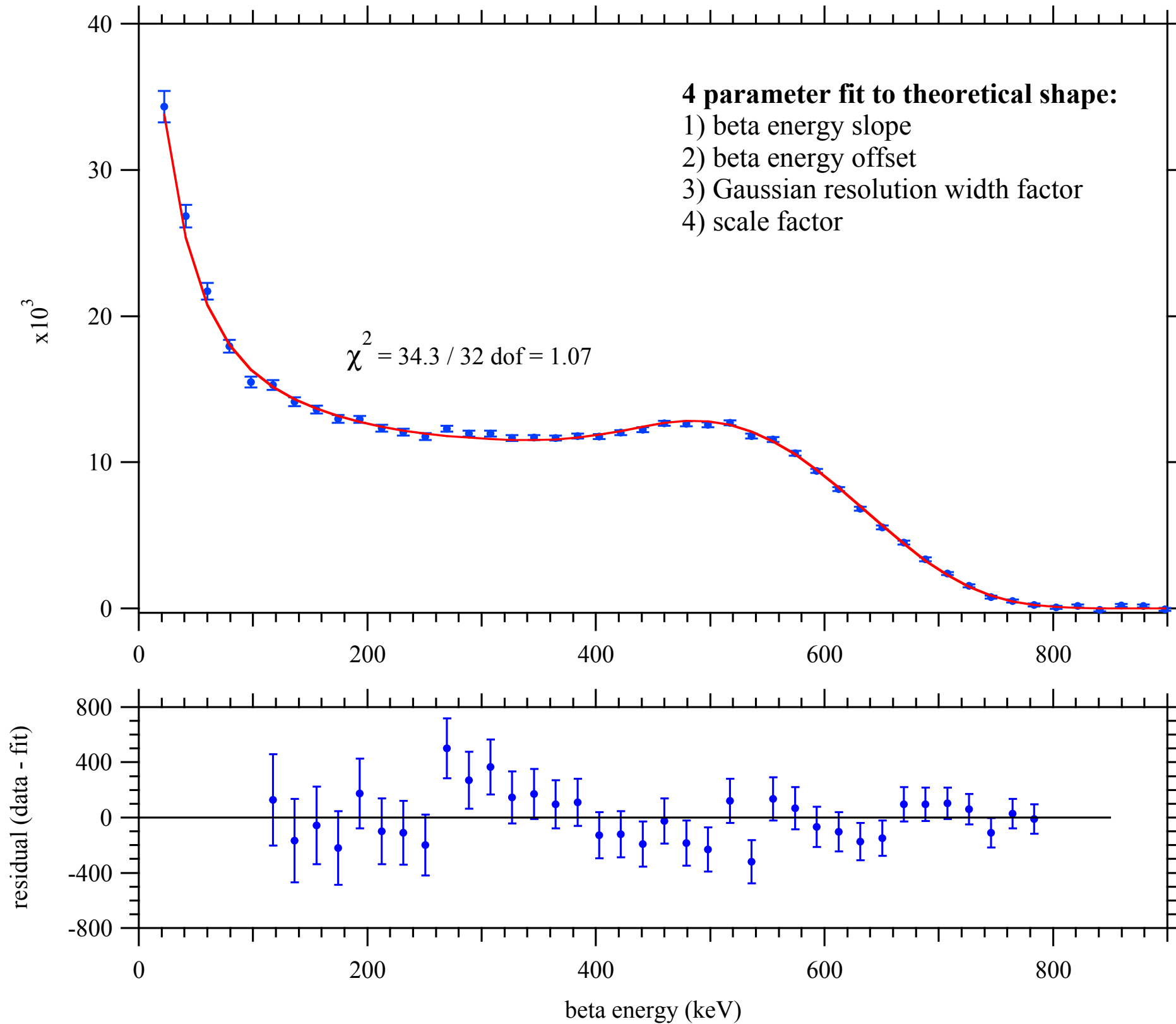




Background
Subtracted
Wishbone



Energy Calibration Fit



The geometric function $f_a(E)$

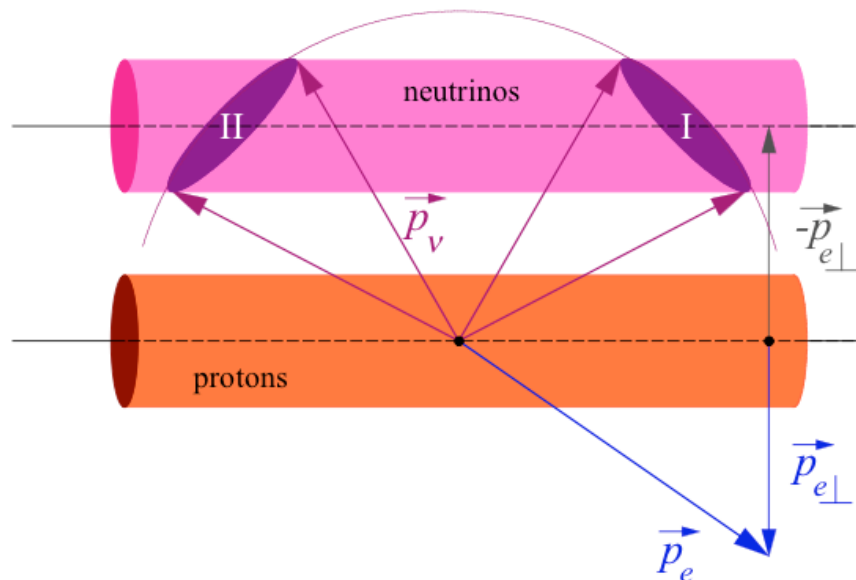
wishbone asymmetry: $X(E) = a f_a(E)$

$$f_a(E) = \frac{1}{2} v (\phi^I(E) - \phi^{II}(E))$$

$\phi^I(E), \phi^{II}(E)$ are the average angle between electron and antineutrino momentum vectors for all momenta within the aCORN acceptance, independent of the beta decay distributions.

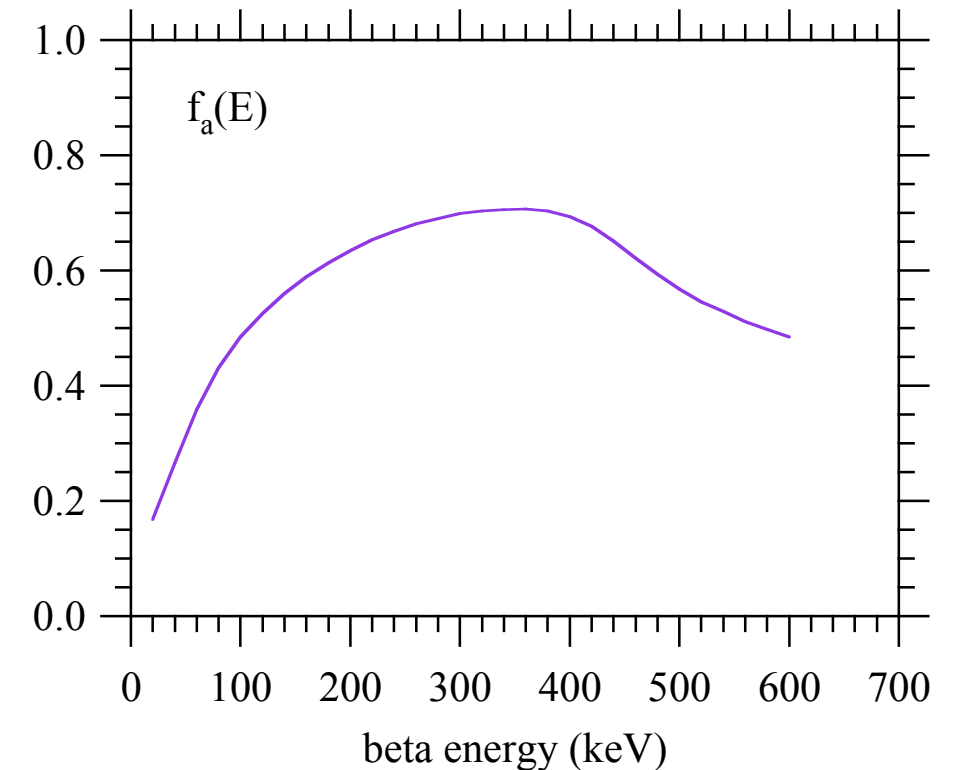
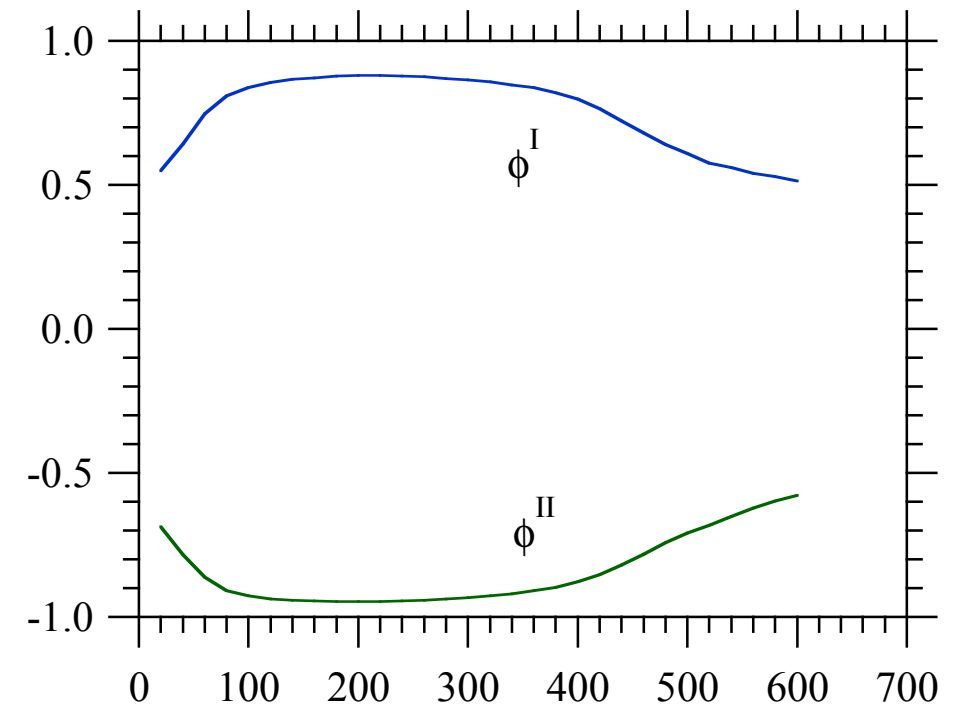
$f_a(E)$ depends ONLY on:

- magnetic field strength
- collimator geometry
- neutron beam density distribution (weakly)

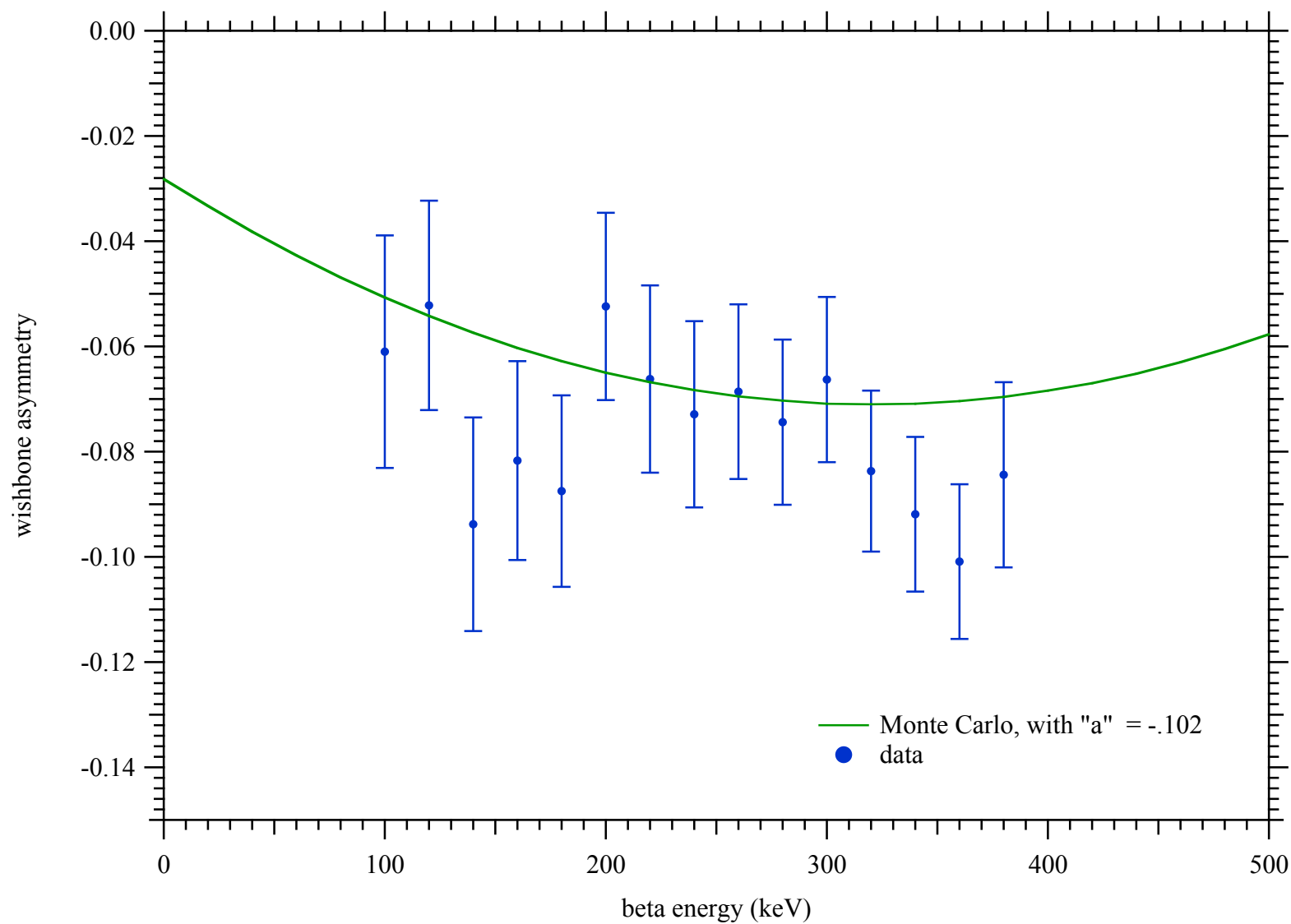


$$\phi^I(E) = \frac{\int d\Omega_e \int_I d\Omega_v \cos \theta_{ev}}{\Omega_e \Omega_v^I}$$

$$\phi^{II}(E) = \frac{\int d\Omega_e \int_{II} d\Omega_v \cos \theta_{ev}}{\Omega_e \Omega_v^{II}}$$



Uncorrected wishbone asymmetry



Energy-dependent corrections

Energy-dependent corrections

Electrostatic mirror

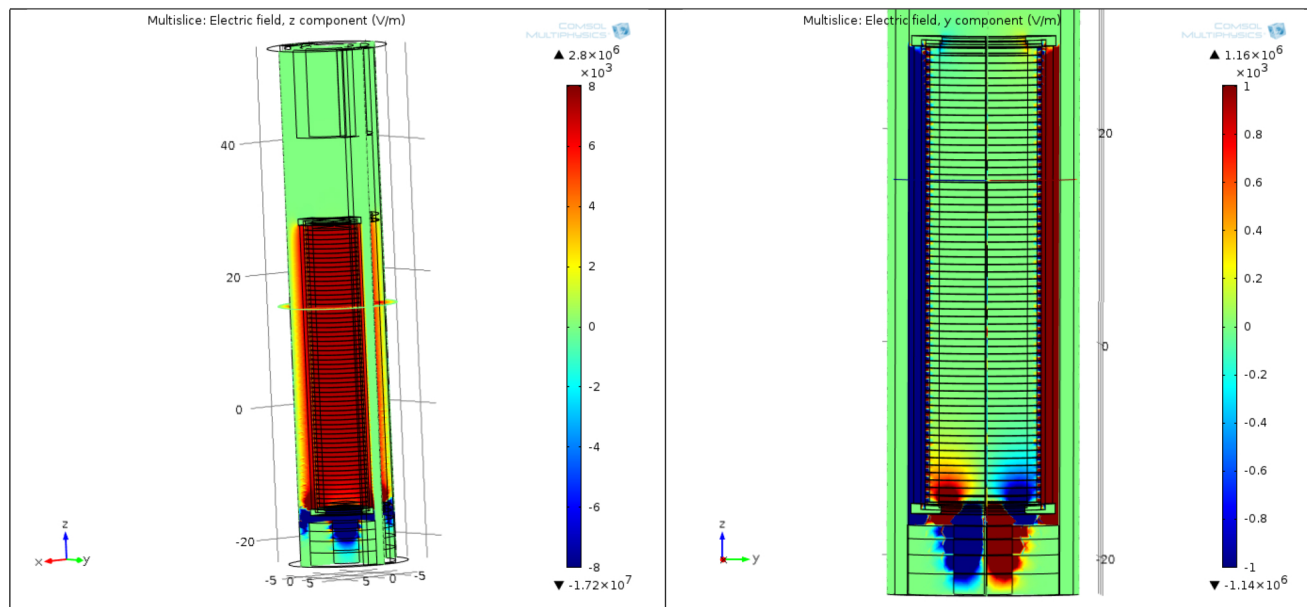
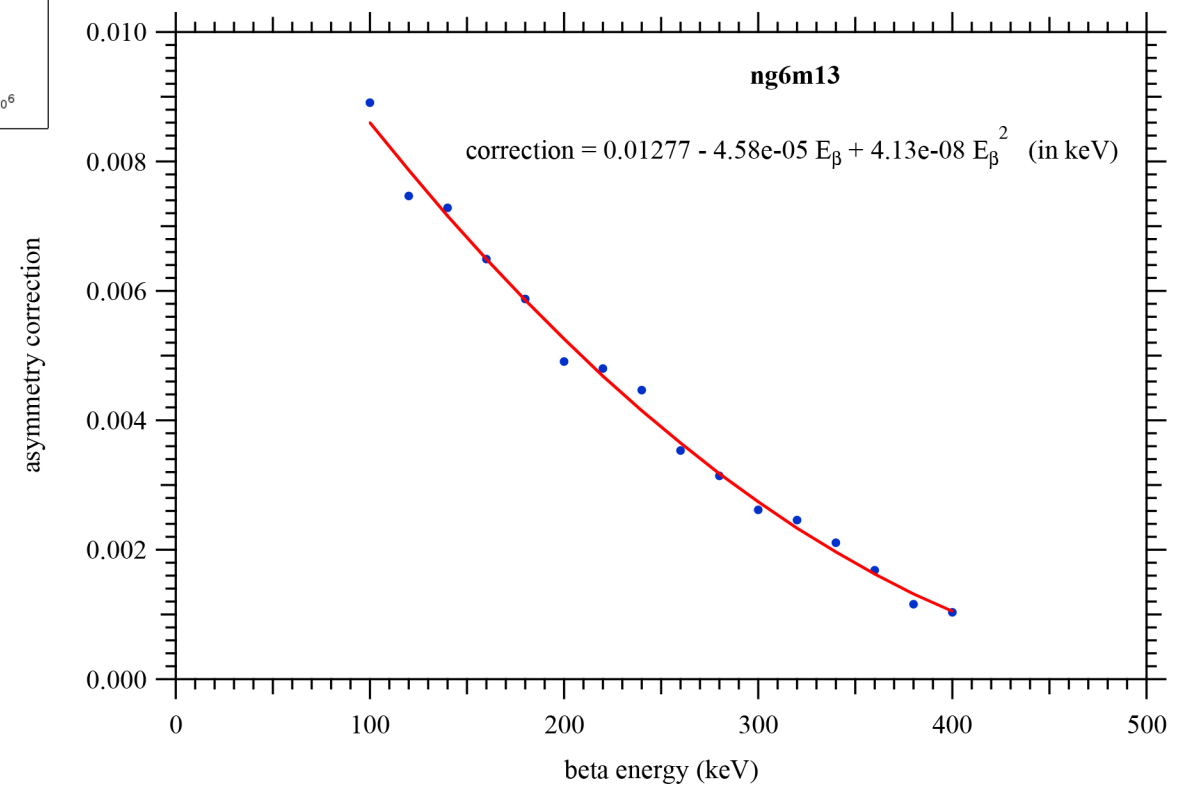


Figure 2: Axial (l) and Transverse (r) Electric fields in NG6Mirror13

4th order RK proton transport simulation

fully realistic 3D COMSOL model



Energy-dependent corrections

Electrostatic mirror

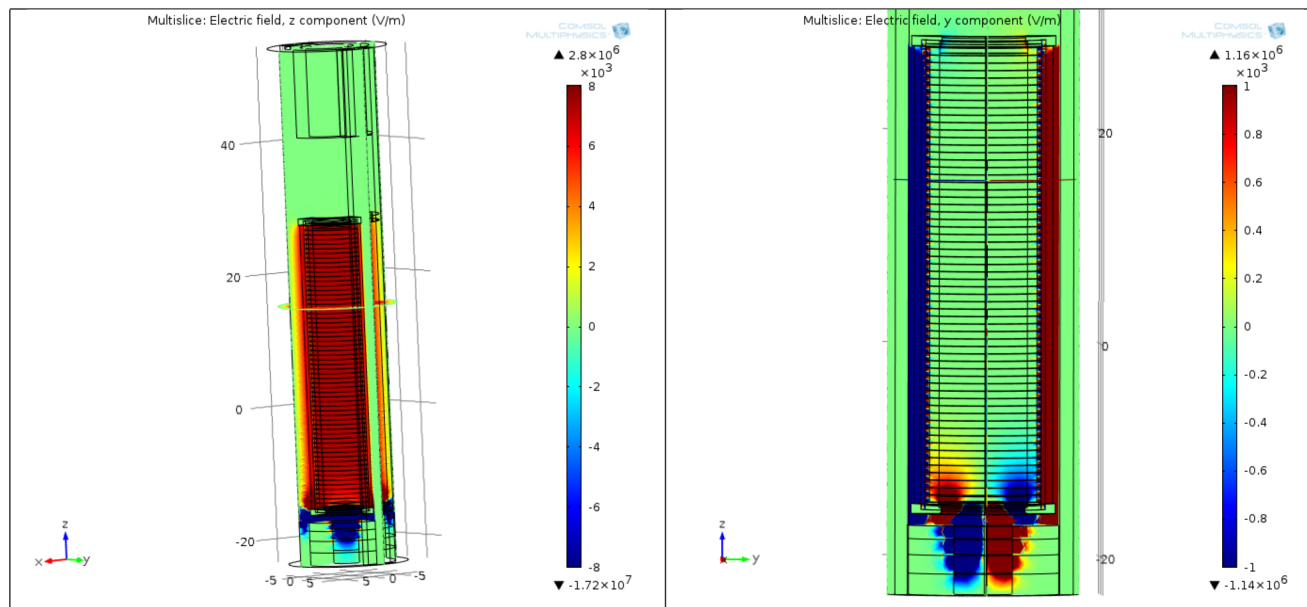
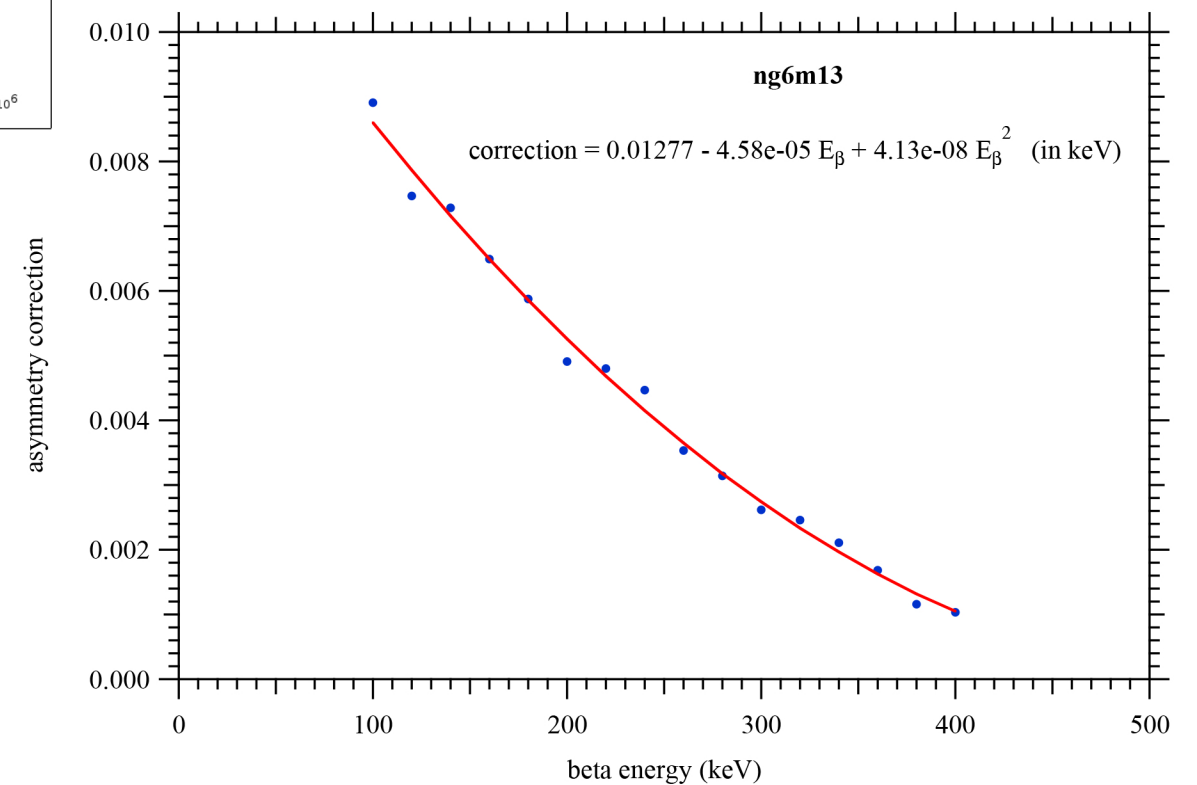


Figure 2: Axial (l) and Transverse (r) Electric fields in NG6Mirror13

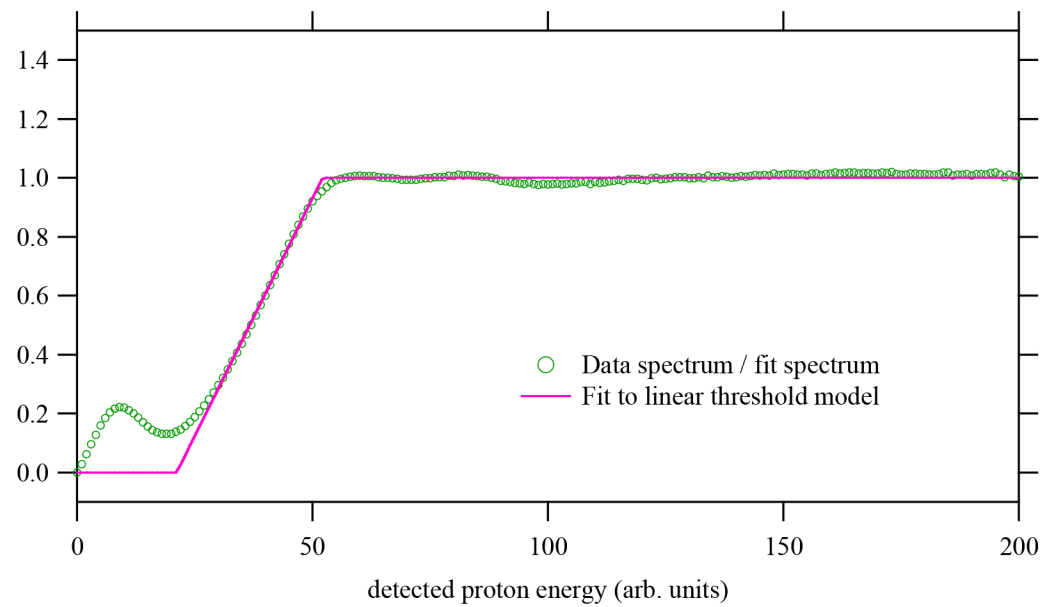
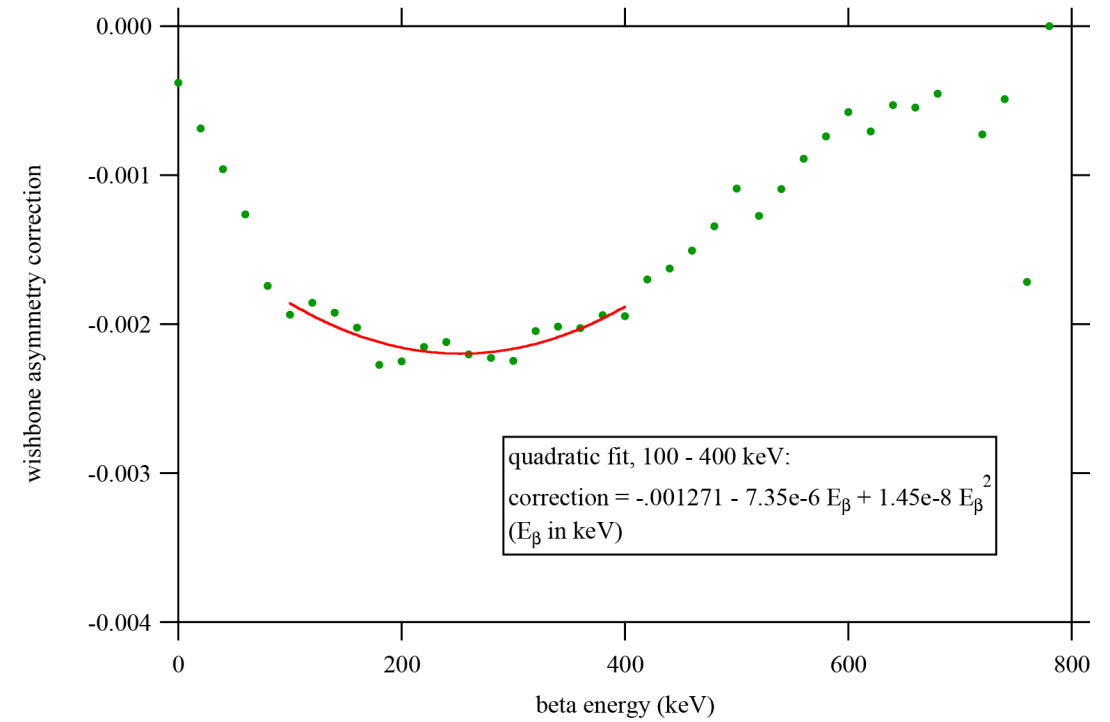
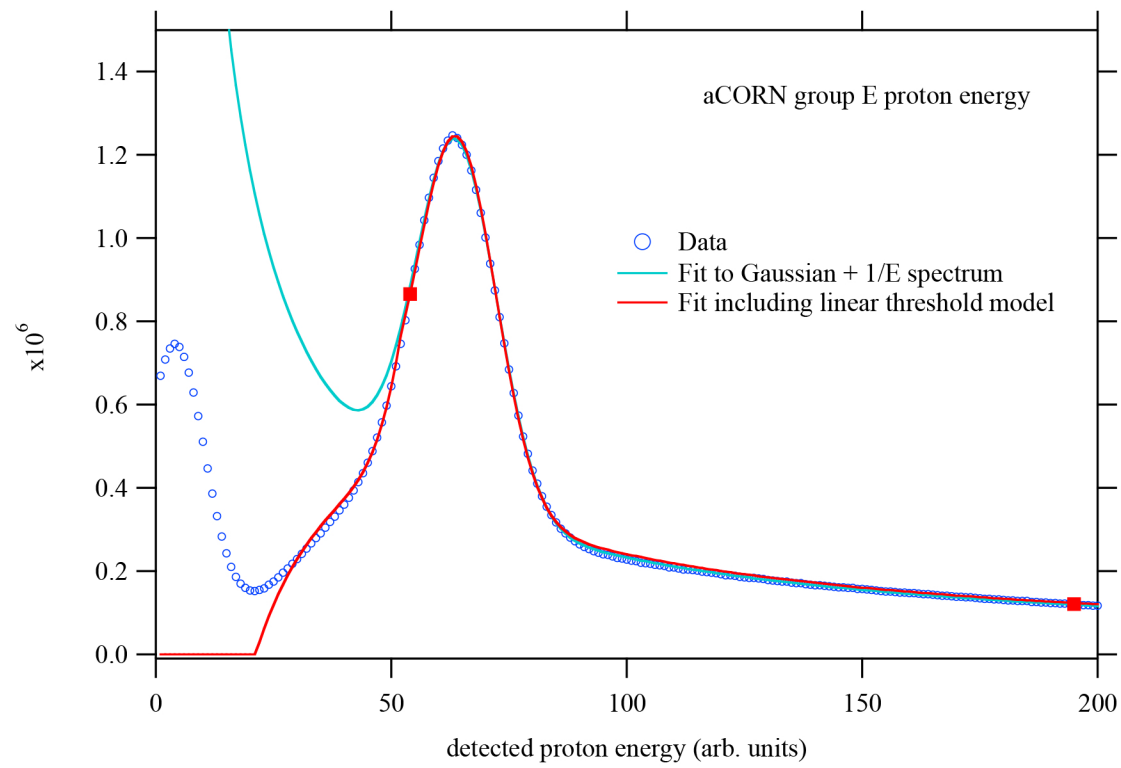
fully realistic 3D COMSOL model

4th order RK proton transport simulation

a +5.5% net correction to "a"

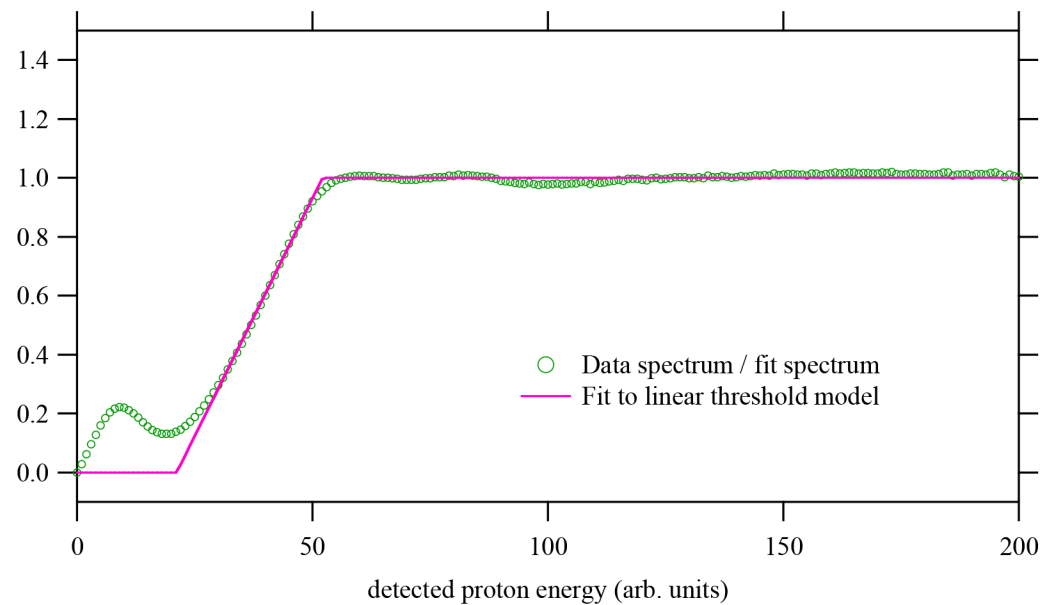
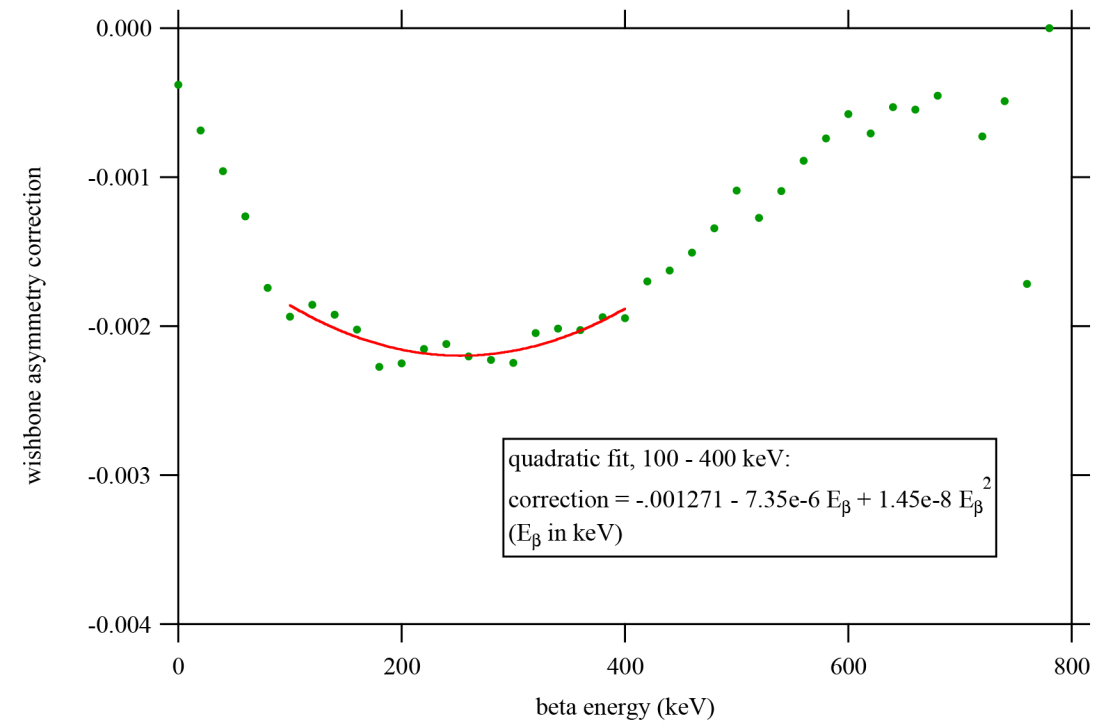
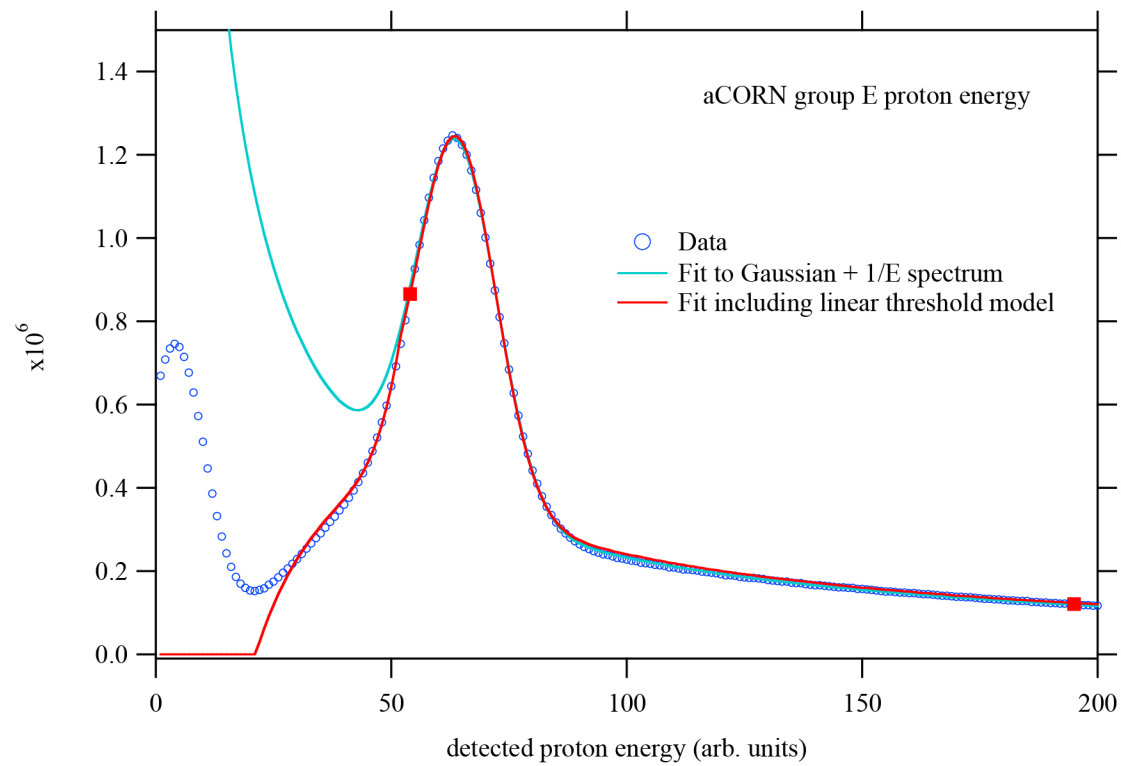


Soft proton threshold



aCORN Monte Carlo calculation
of the proton threshold effect

Soft proton threshold



aCORN Monte Carlo calculation
of the proton threshold effect

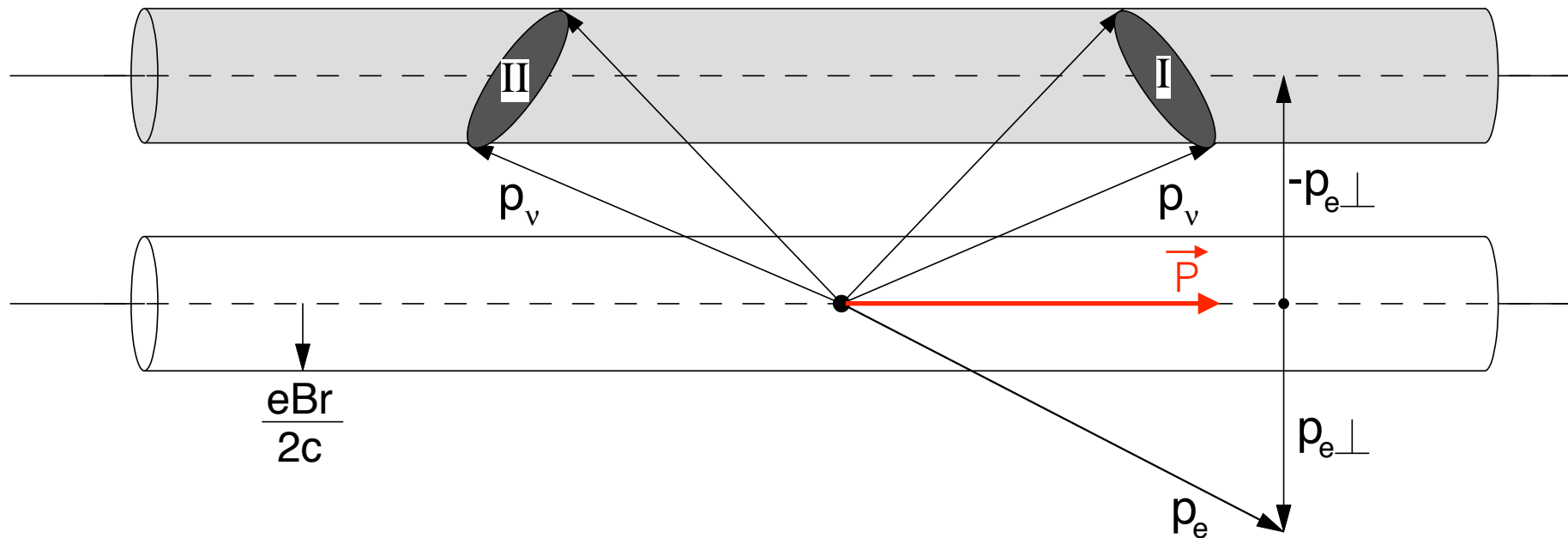
a -3.0% net correction to "a"

Beam Polarization

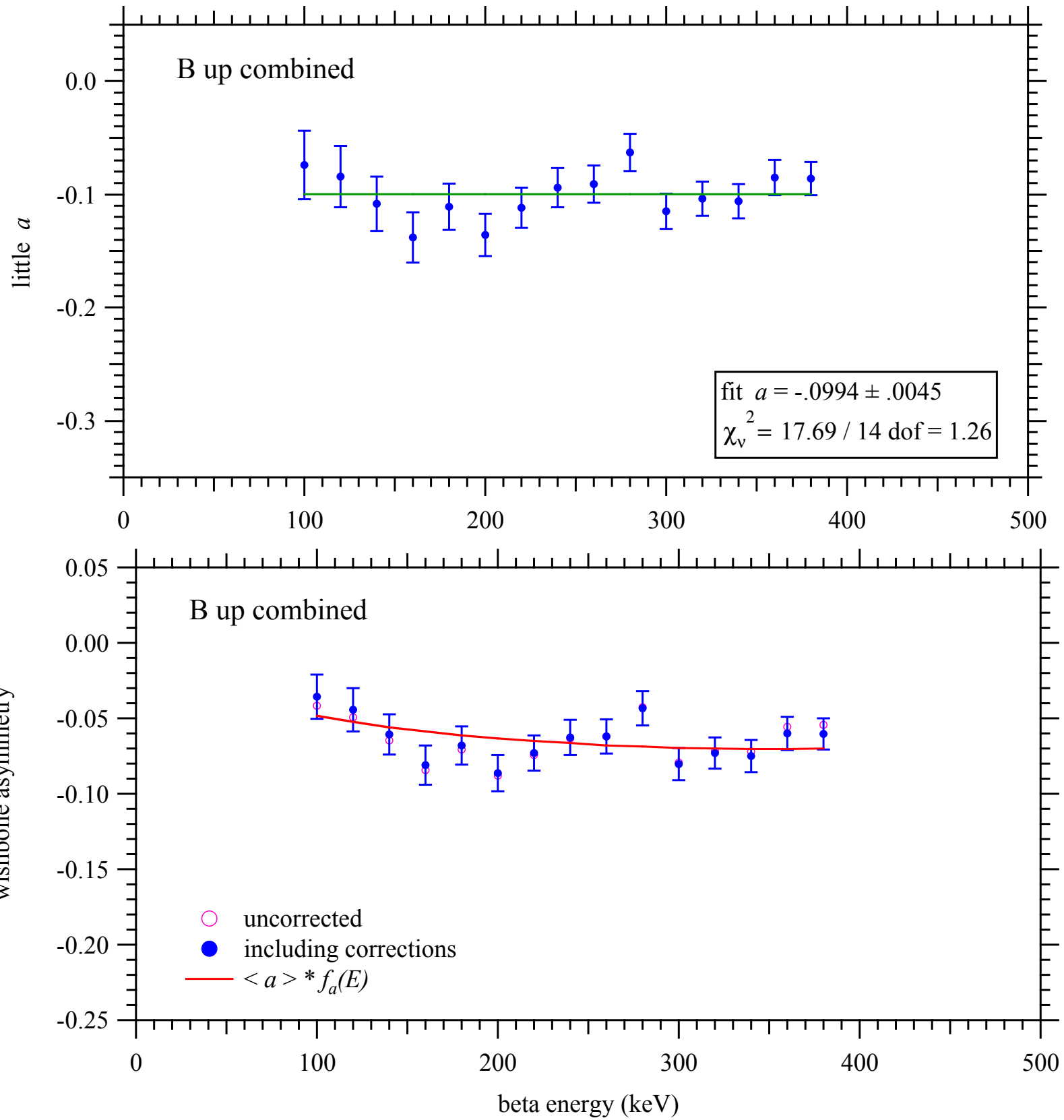
With a polarized neutron beam:

$$\text{wishbone asymmetry } A_{wb} = af_a(E_\beta) + PBf_B(E_\beta)$$

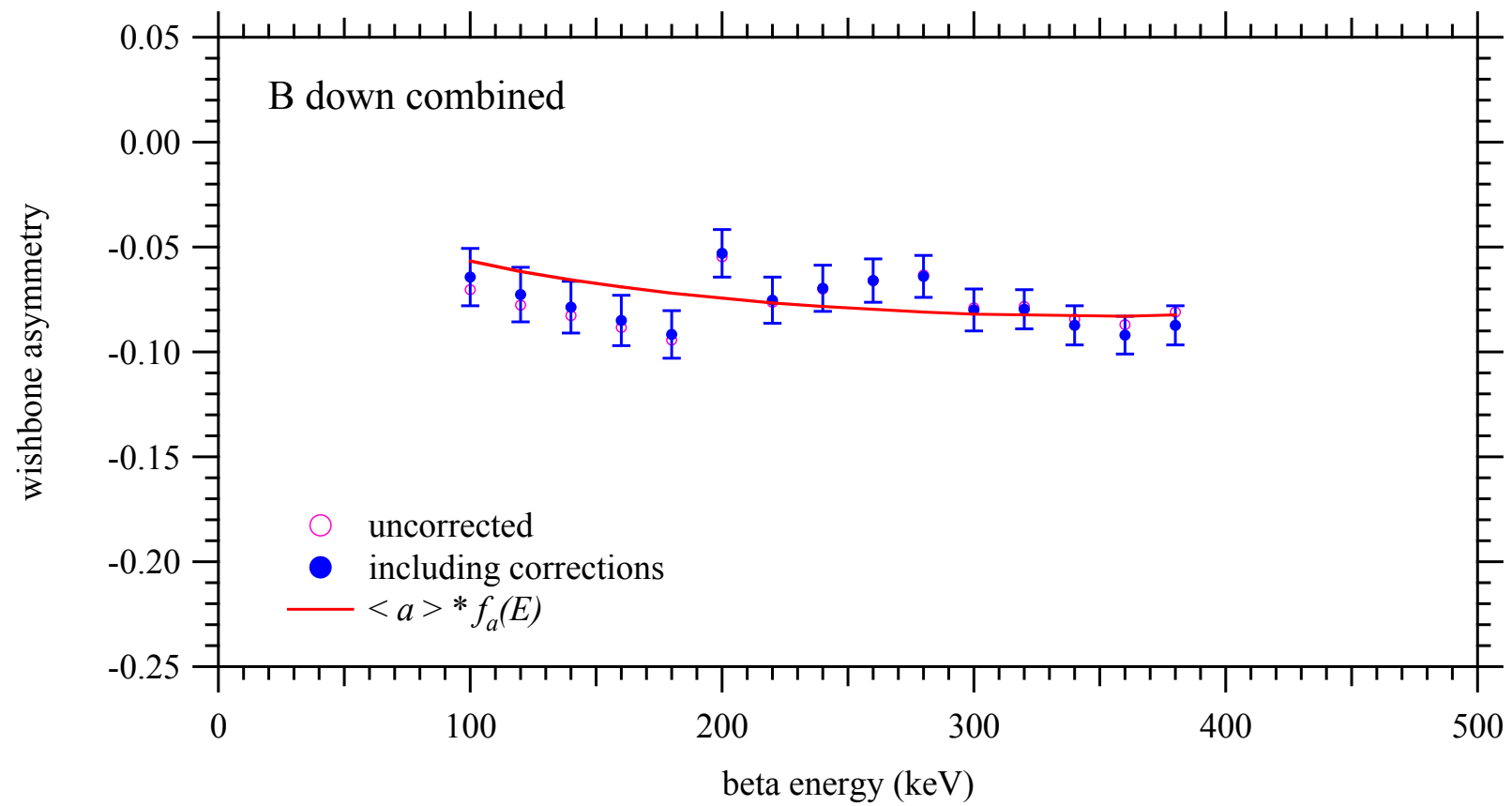
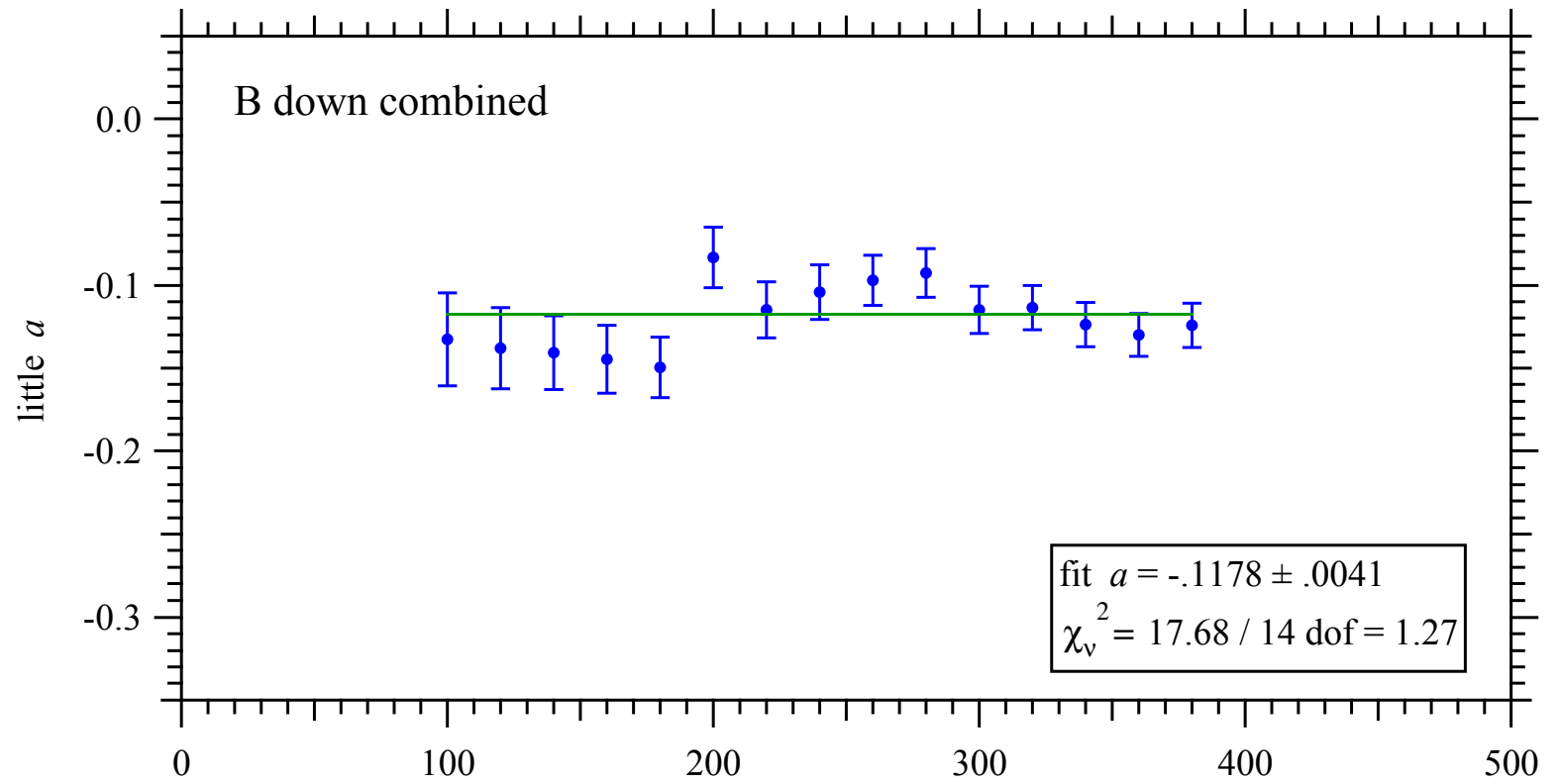
$$\frac{Bf_B(E_\beta)}{af_a(E_\beta)} \approx 14$$

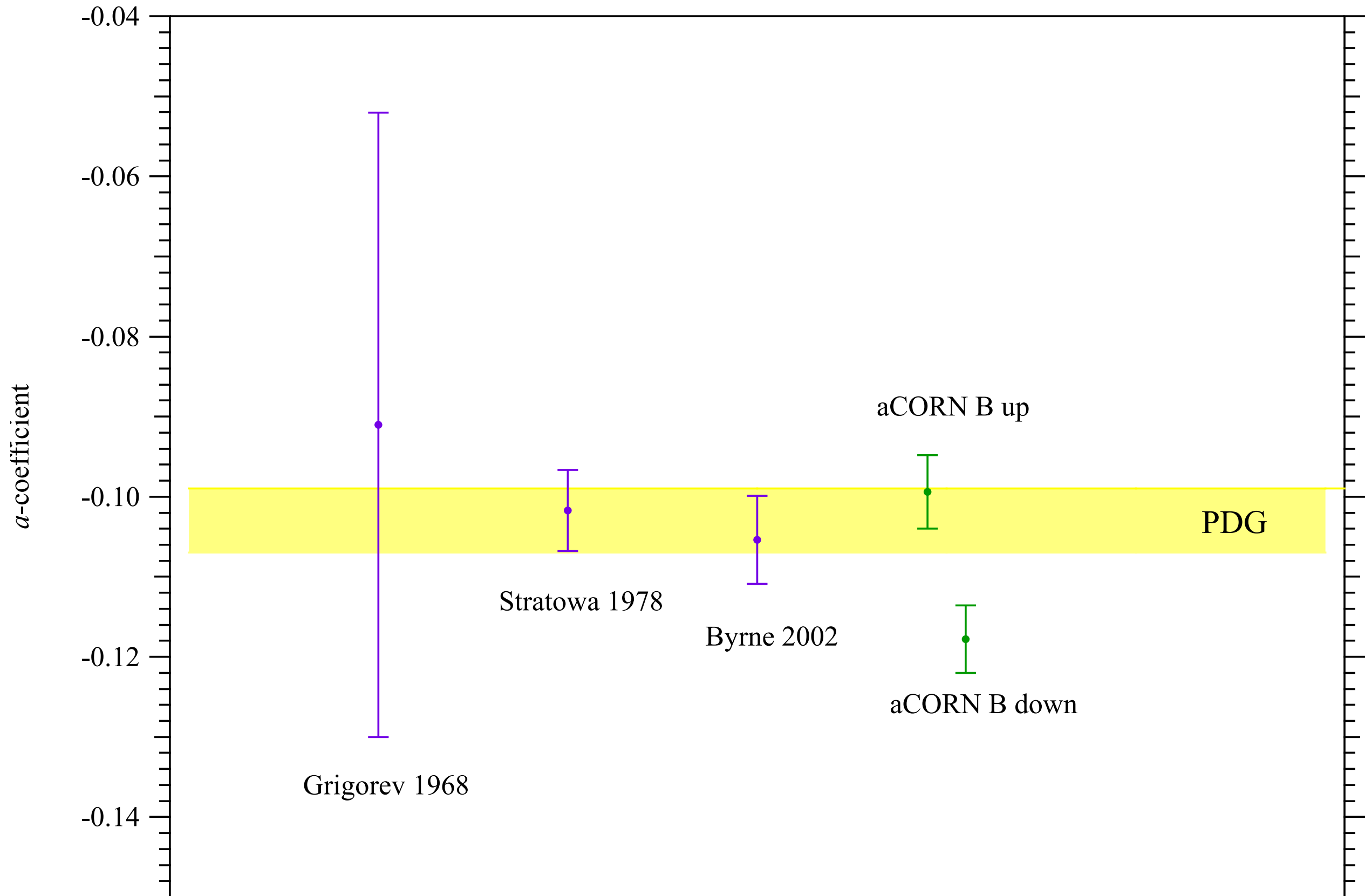


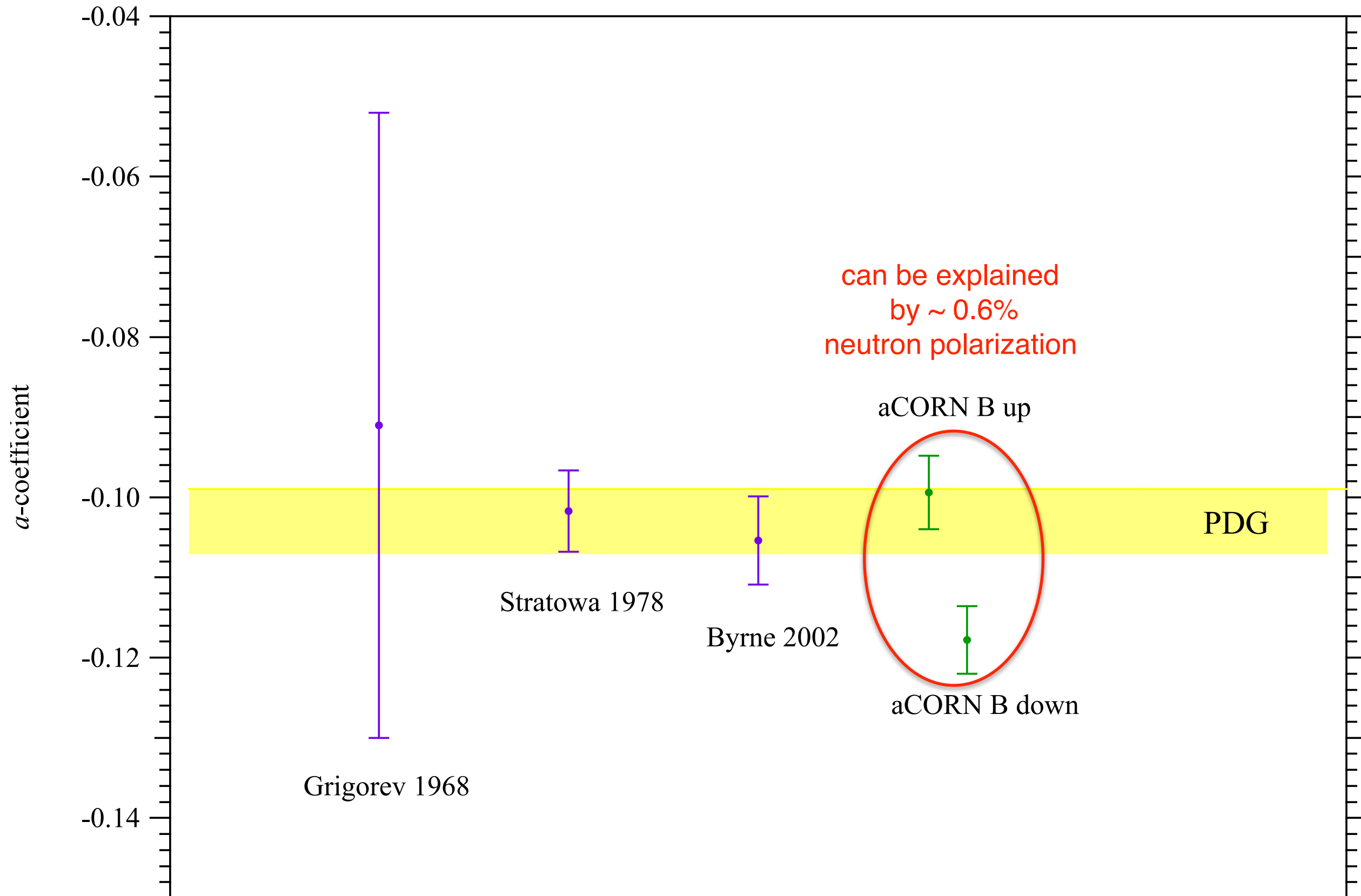
Ratio of $X(E) / f_a(E)$

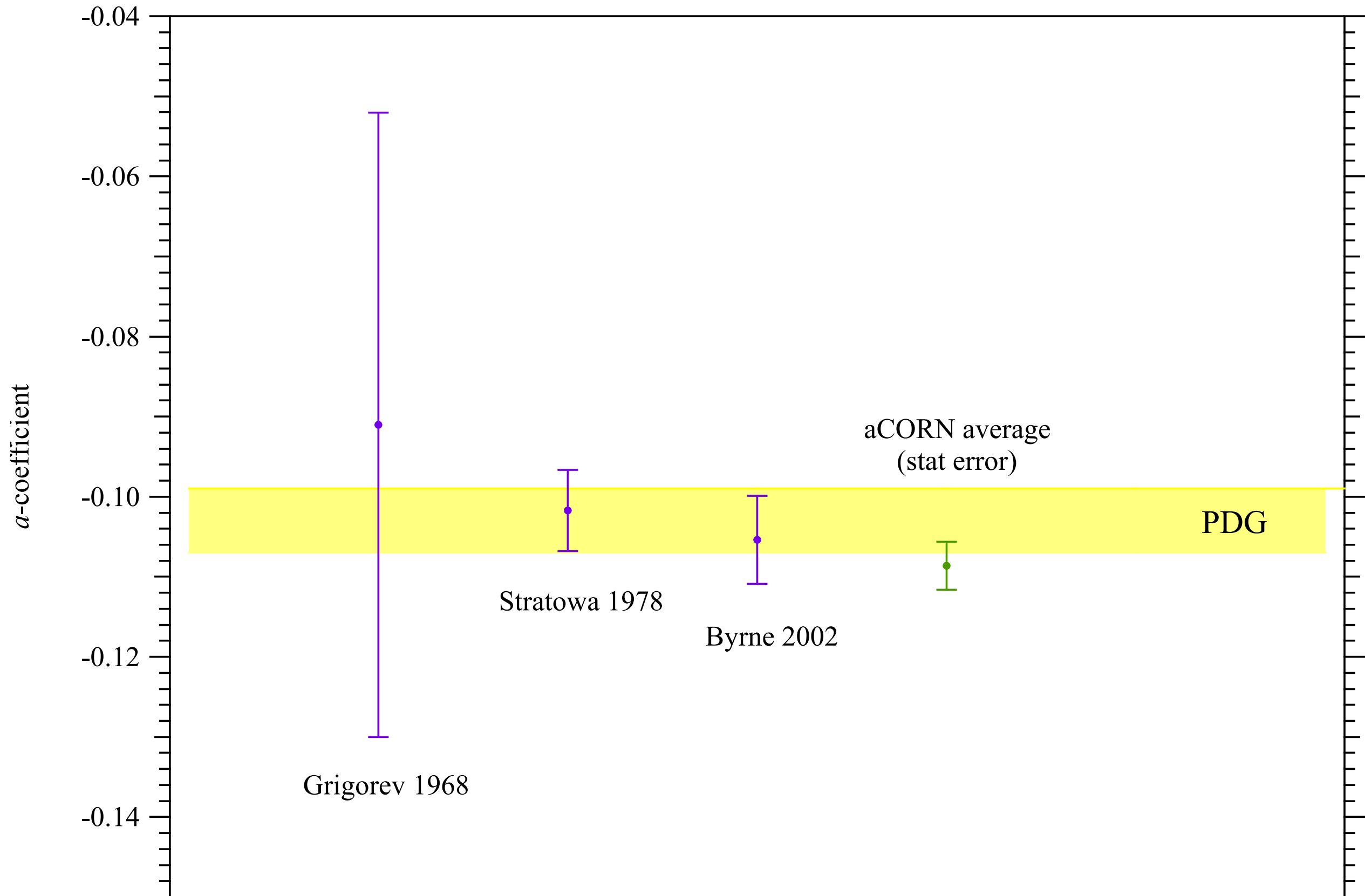


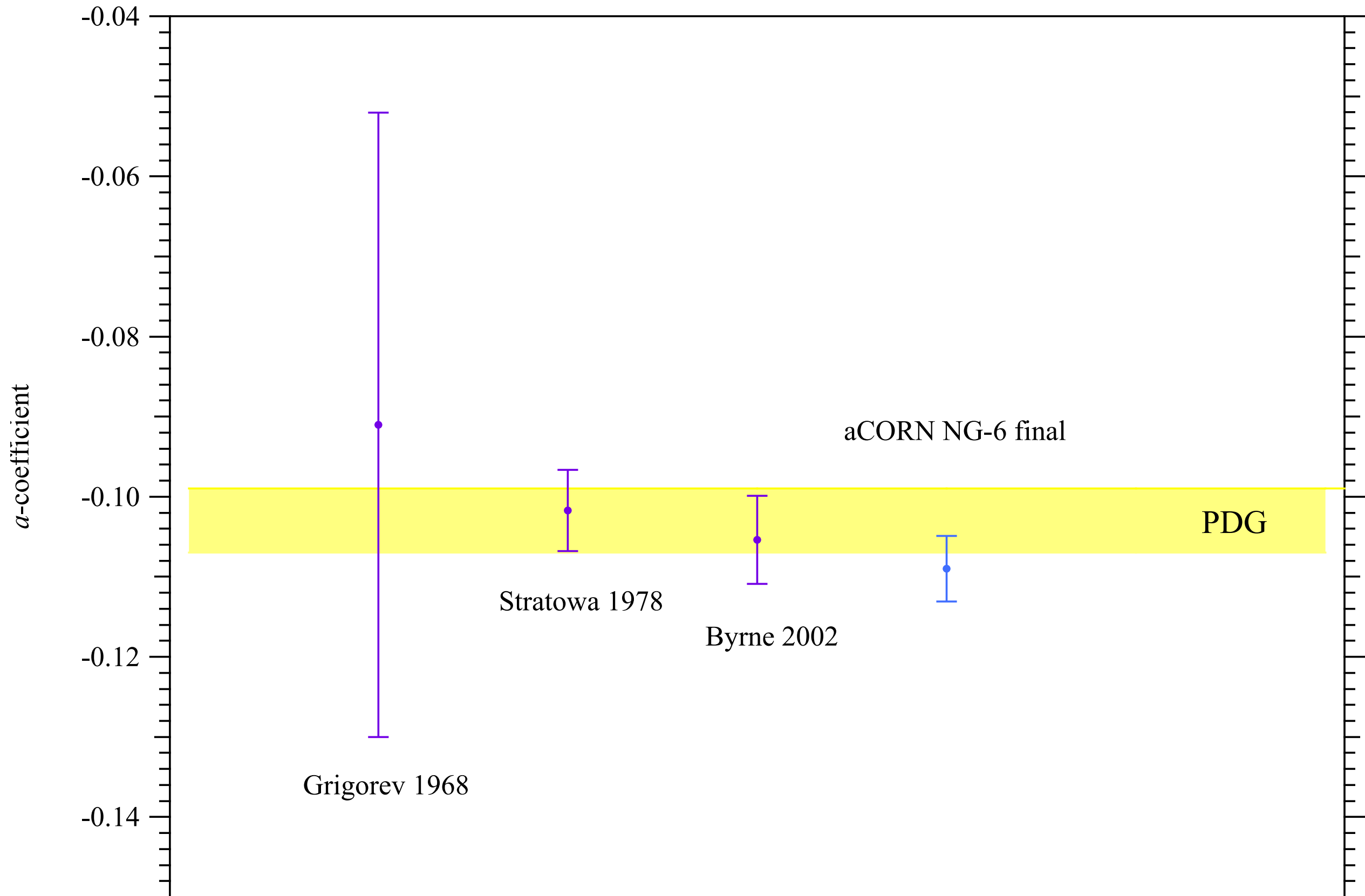
Ratio of $X(E) / f_a(E)$











aCORN NG-6 Result

	correction	1 σ uncert.	relative
electrostatic mirror	0.00571	0.00114	0.0105
proton threshold	-0.00318	0.00076	0.0070
energy loss in grid	-0.00111	0.00022	0.0020
absolute B field	-0.00010	0.00050	0.0046
B field shape	0.00031	0.00082	0.0075
residual gas	0.00046	0.00046	0.0042
e scattering	-0.00153	0.00153	0.0140
beta energy calibration		0.00031	0.0028
proton collimator align.		0.00050	0.0046
p scattering	0.00041	0.00050	0.0046
p focusing	0.00010	0.00010	0.0009
wishbone asymmetry		0.00100	0.0091
beam polarization		0.00102	0.0094
total systematic	0.00107	0.00283	0.0260
statistical		0.00302	0.0277
total uncertainty		0.00414	0.0380

aCORN NG-6 Result

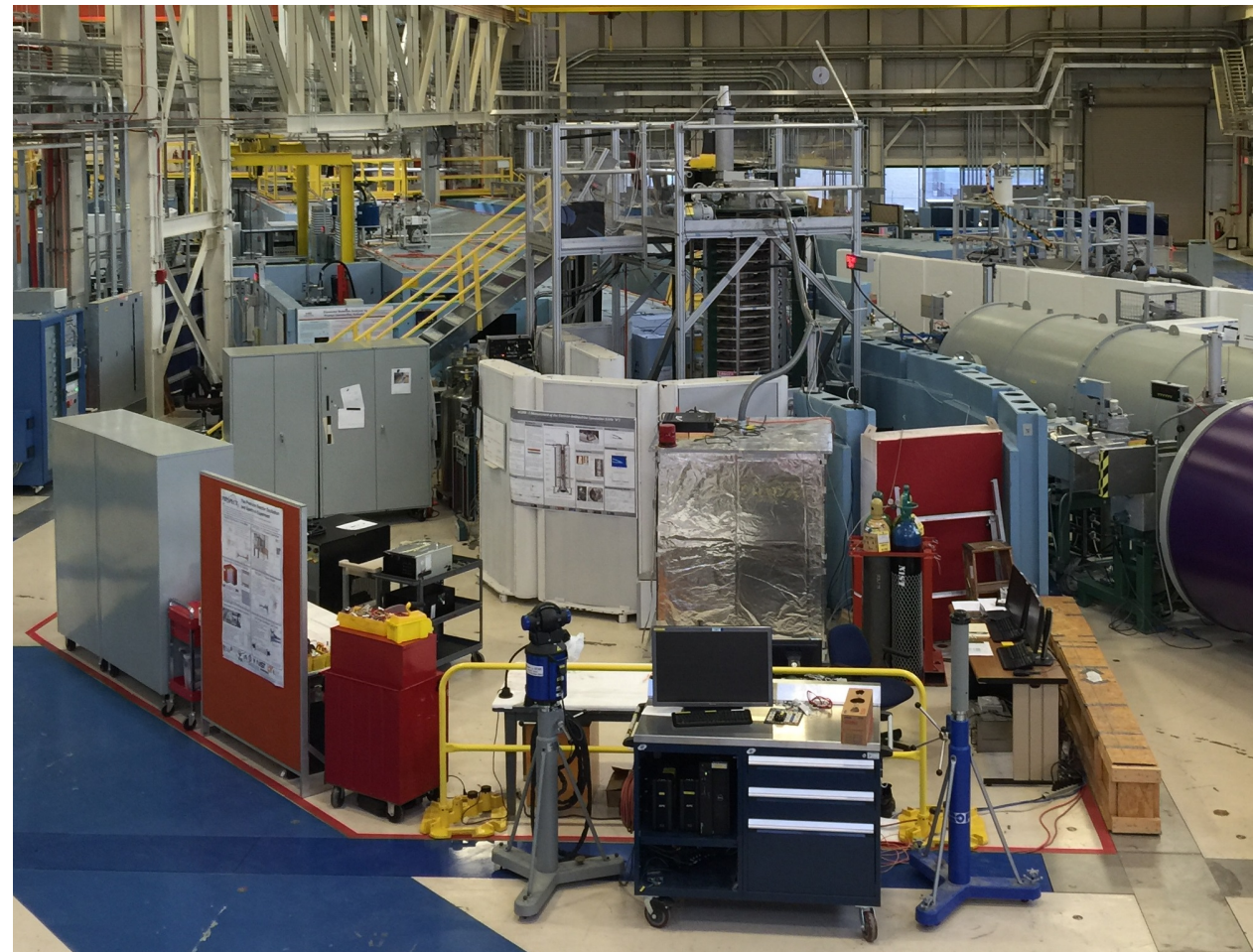
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beam polarization		0.00102	0.0094
total systematic	0.00107	0.00283	0.0260
statistical		0.00302	0.0277
total uncertainty		0.00414	0.0380

$$\alpha = -0.1090 \pm 0.0030 \text{ (stat)} \pm 0.0028 \text{ (sys)}$$

G. Darius, *et al.* Phys. Rev. Lett. **119**, 042502 (2017)

aCORN on new NG-C beamline

- aCORN moved to new NG-C end position at NIST in 2015
- Ran on NG-C from July 2015 - September 2016
- ~ 5x wishbone event rate, signal/bkgd similar to NG-6
- Collected a good data set ~10 times NG-6
- Improved systematics
- Analysis in progress
- **We expect a new result with relative uncertainty $< 2\%$**



aCORN B

Neutron Decay Parameters

Phenomenological ($J = 1/2 \rightarrow J = 1/2$) beta decay formula [Jackson, Treiman, Wyld, 1957] :

$$dW \propto \frac{1}{\tau} F(E_e) \left[1 + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + b \frac{m_e}{E_e} + A \frac{\vec{\sigma}_n \cdot \vec{p}_e}{E_e} + B \frac{\vec{\sigma}_n \cdot \vec{p}_\nu}{E_\nu} + D \frac{\vec{\sigma}_n \cdot (\vec{p}_e \times \vec{p}_\nu)}{E_e E_\nu} \right]$$

For allowed beta decay, neglecting recoil order terms, the standard electroweak model (Weinberg, Glashow, Salam, et al.) predicts:

$$a = \frac{1 - \lambda^2}{1 + 3\lambda^2} \quad b = 0 \quad A = -2 \frac{\lambda^2 + \text{Re}(\lambda)}{1 + 3\lambda^2} \quad B = 2 \frac{\lambda^2 - \text{Re}(\lambda)}{1 + 3\lambda^2}$$

$$D = 2 \frac{\text{Im}(\lambda)}{1 + 3\lambda^2} \approx 0 \quad \tau \propto \frac{1}{g_V^2 + 3g_A^2} \quad \text{where} \quad \lambda \equiv \frac{g_A}{g_V}$$

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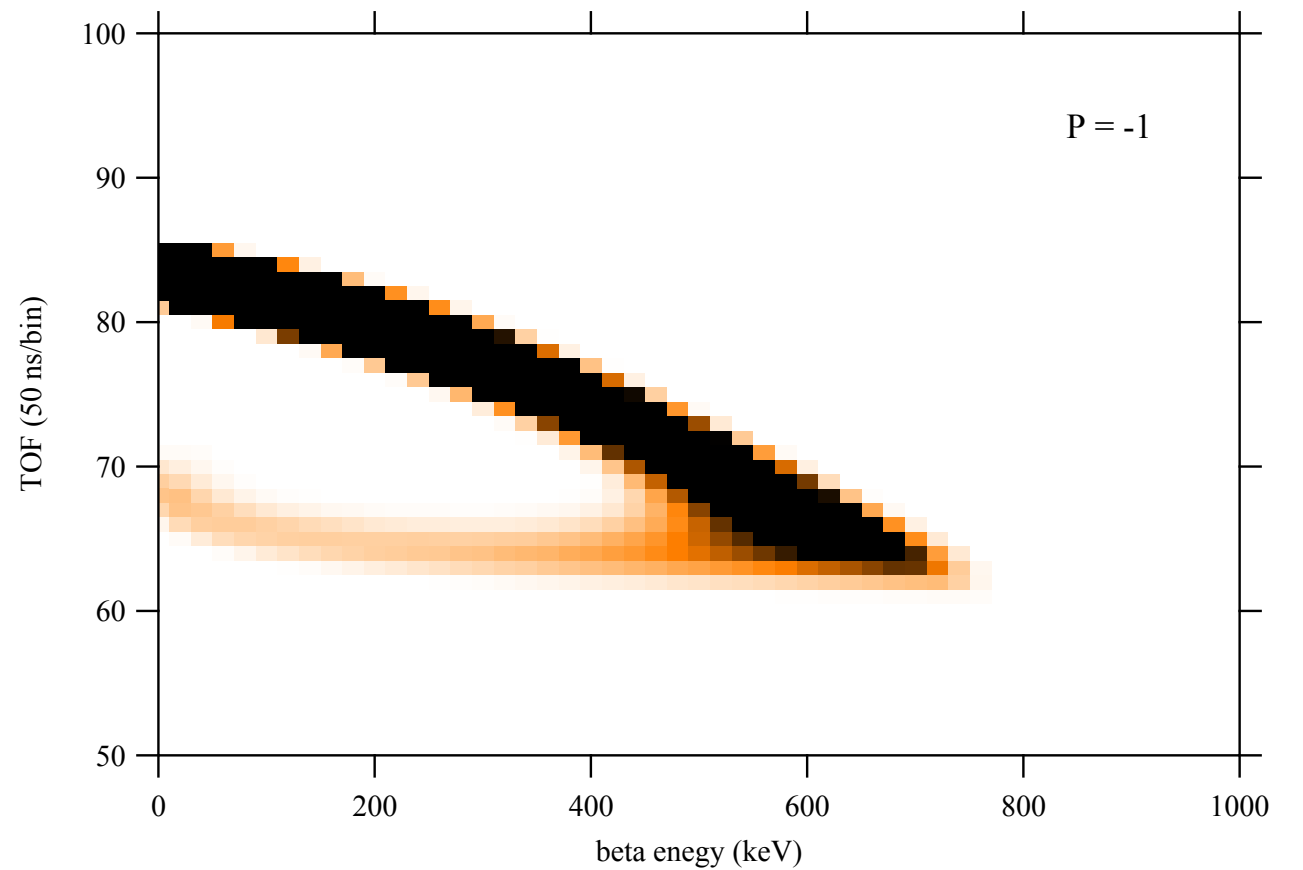
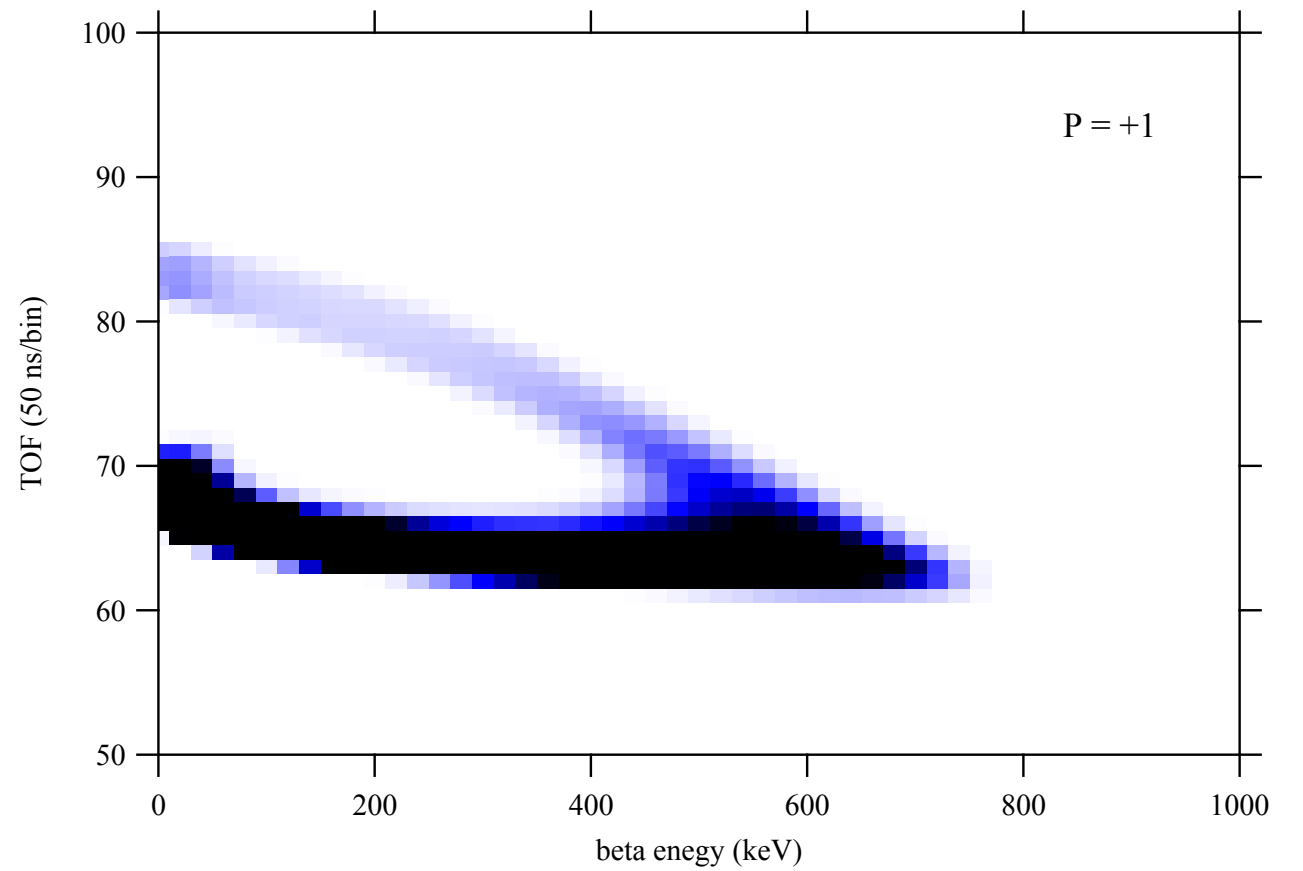
$$D = 2 \frac{\text{Im}(\lambda)}{1 + 3\lambda^2} \approx 0 \quad \tau \propto \frac{1}{g_v^2 + 3g_A^2} \quad \text{where} \quad \lambda \equiv \frac{g_A}{g_V}$$

of these correlation coefficients, B has the least sensitivity to λ
but the most sensitivity to possible right-handed currents

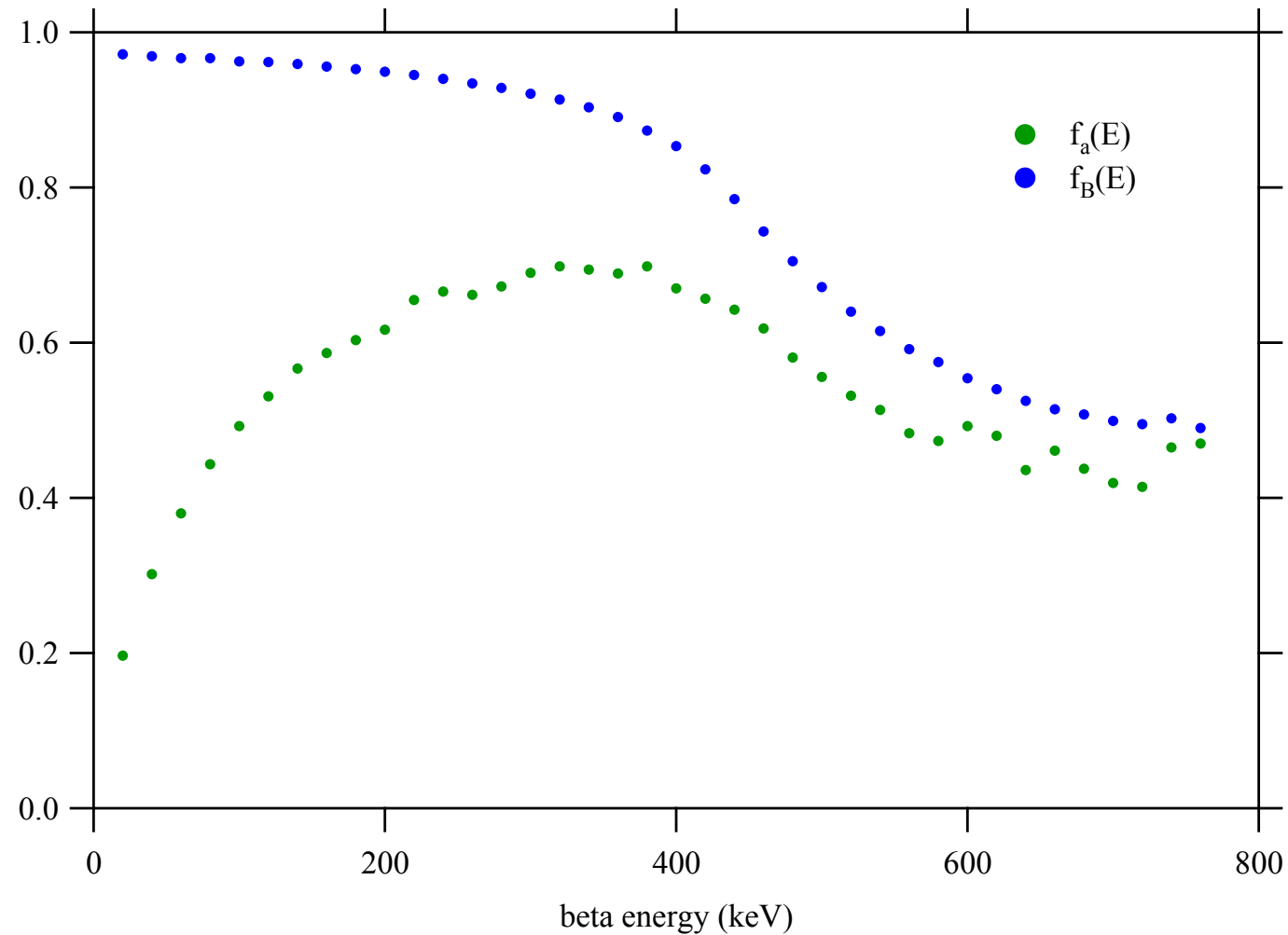
Monte Carlo:

measure both flip states:

$$\frac{A_{wb}^+ - A_{wb}^-}{2} = PBf_B(E_\beta)$$



Monte Carlo:



Statistics estimate:

- assume factor 10 lower neutron flux (XSM polarizer, collimation)
- 150 beam days (~1 year) → 1% little "a"
- 14x larger asymmetry signal
- assume S/B same as aCORN

$$\frac{\sigma_B}{B} \approx \frac{\sqrt{10}}{14} (1\%) = 0.0023 \text{ (stat)}$$

aCORN

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