

Stable Sexaquark as Dark Matter



Glennys R. Farrar
New York University

Stable Sexaquark as Dark Matter



How could we have missed a
stable particle made of quarks?

[Hints from Astrophysics]

Primordial Nucleosynthesis

Dark-Matter to Ordinary-Matter ratio

Detecting S dark matter

Unique among multi-quark states:

Fermi statistics is compatible with a totally symmetric spatial wave function AND

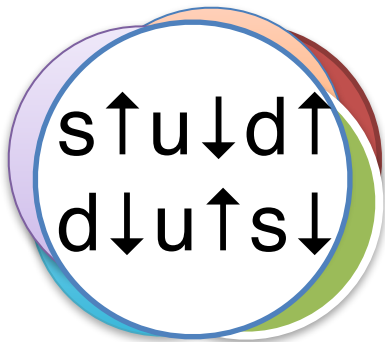
antisymmetric (singlet) in:

color

flavor

spin

totally symmetric in space



S

(Most-Attractive Channel)³ :

6-quark, Q=0, B=2

Spin-0, scalar

Flavor singlet

$m_S < 2 \text{ GeV}???$

Same quark content as H-dibaryon* (Jaffe 1977), but different physics: **not a loosely bound di- Λ !**

*mass $\sim 2150 \text{ MeV}$ in bag model — decays in 10^{-10} s

Why consider $m_S \sim 2 m_p$?

- Light quarks almost massless, i.e. relativistic
 - $m_{u,d} \approx 0$, $m_s = 91 \text{ MeV}$
- S has same QNs as ground state glueball
 - why not $m_S \approx m_{\text{glueball}} + 180 \text{ MeV} = (1.5-1.7) + 0.18 \text{ GeV} \approx 2 m_p$
- 3 x di-quark mass = 1.2 - 2-ish GeV
- $m_S < 2 (m_p + m_e)$: S is absolutely stable
- $m_S > 2 (m_p - 8 \text{ MeV})$: nuclei are stable Interesting DM candidate
- triple-singlet (color, flavor, spin): MAC, lattice, almost all models $\Rightarrow m_S < 2 m_\Lambda$
- *extensive experimental searches exclude weak-lifetime & $m > 2 \text{ GeV}$*
- bound state exists and mass $< 2 \text{ GeV}$ ($\tau > \tau_{\text{Univ}}$ or stable)

Stable Sexaquark Hypothesis

6	sexa- ^[19]	–	sen- ^[20]	sext- ^[21]	hex- ^[22]	hexakis- hexaplo- hexad- e.g. hexahedron	hect- ^[23] hectai-	shat-
---	-----------------------	---	----------------------	-----------------------	----------------------	---	----------------------------------	-------

^{a b} Sometimes Greek *hexa-* is used in Latin compounds, such as *hexadecimal*, due to **taboo avoidance** with the English word *sex*.
https://en.wikipedia.org/wiki/Numeral_prefix

Crucial fact:

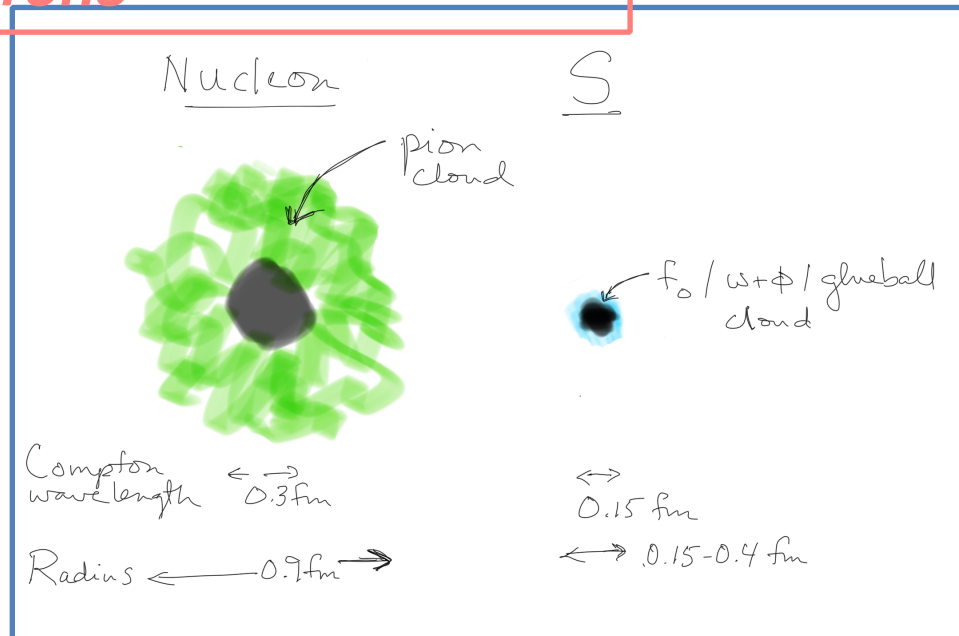
S does not couple to pions => much smaller than usual hadrons => hard to produce with hadrons

6-quark, $Q=0$, $B=2$

Spin-0, scalar

Flavor singlet

$m \sim 1.7-2 \text{ GeV}$



Stable S?

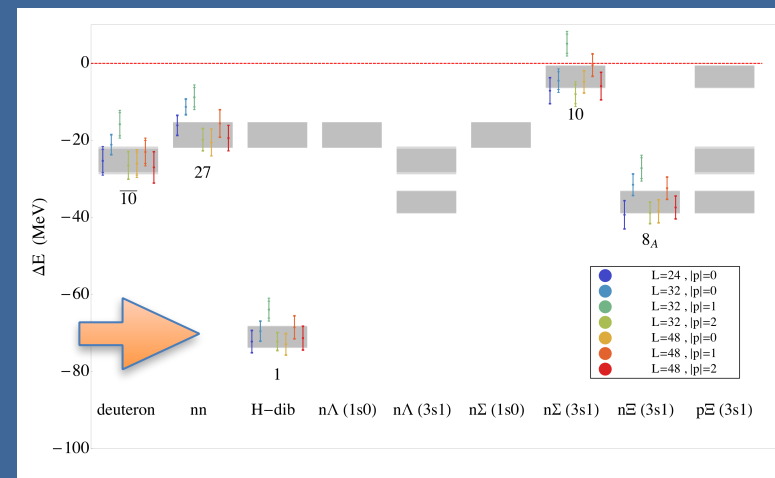
- $\tau > \tau_{\text{Univ}}$
 - $M_S < 2 m_p + 2 m_e = 1877.6 \text{ MeV} \rightarrow$ **absolutely stable**
 - $M_S > 2 m_p + 2 \text{ BE} = 1860 \text{ MeV} \rightarrow$ **nuclei absolutely stable**
 - higher and lower mass may also work $\Gamma \sim G_F^4 \times (\text{wave function overlap})^2$

- **Lattice predicts binding (Beane+13)**

- ($m_q = 850 \text{ MeV}$ so not realistic)
- 80 MeV binding

- **Experiments exclude decaying S**

=> it must be STABLE ! ;-)



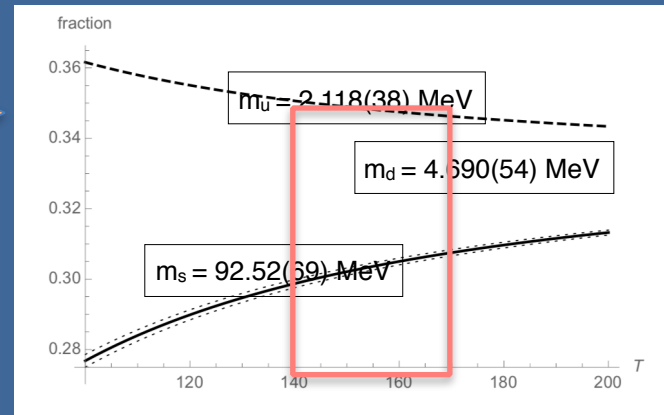
Conditions on QCD Dark Matter

- ✓ $\tau_{\text{DM}} > \tau_{\text{Univ}}$, cold, neutral
- ✓ primordial nucleosynthesis
- ✓ Particle must not be already excluded
 - accelerator searches
 - exotic isotopes
 - **DM searches**
 - indirect impacts (heating planets, helioseismology,...)
 - stability of nuclei
 - equation of state of neutron stars (and their stability)
- ✓ Correct relic density (for natural m_{DM} & $\sigma_{\text{f.o.}}$)

Ω_{DM} / Ω_b

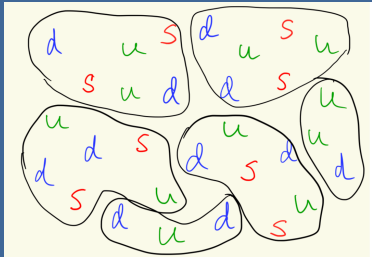
follows from **stat mech** , **quark masses** & temperature of **QGP-hadronization** transition

$$\frac{\Omega_{DM}}{\Omega_b} = \frac{m_S / (2m_p) y_b \kappa_S 3f_S}{1 - \kappa_S 3f_S}$$



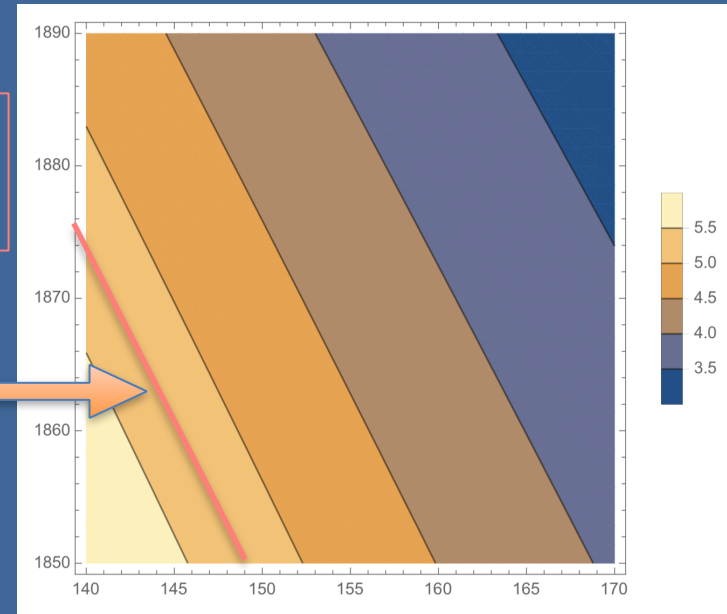
$$\kappa_S(m_S, T) = \frac{1}{1 + (r_{\Lambda, \Lambda} + r_{\Lambda, \Sigma} + 2r_{\Sigma, \Sigma} + 2r_{N, \Xi})}$$

$$r_{1,2} \equiv \exp[-(m_1 + m_2 - m_S)/T]$$



Prediction is both correct AND accurate to ~20% for entire range (uncertainties cancel)

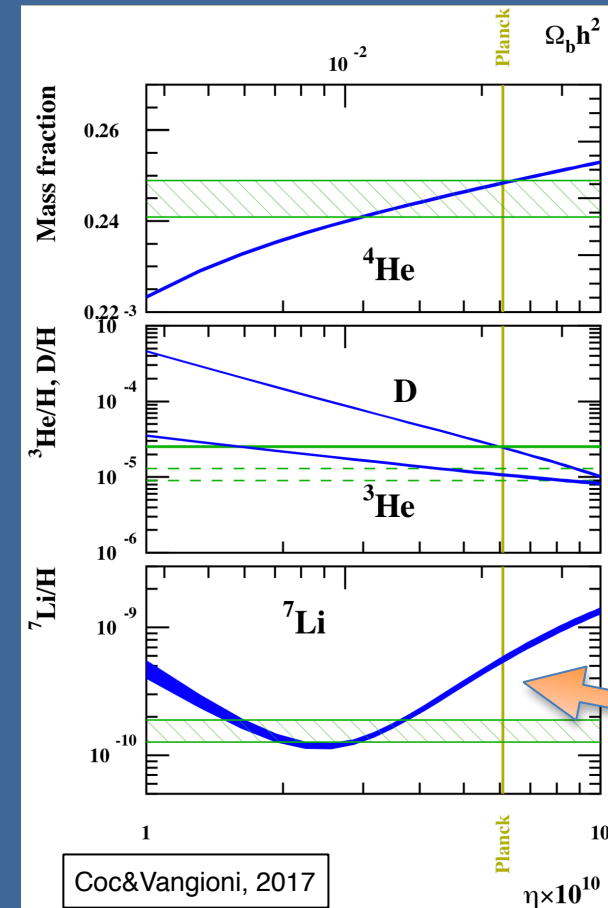
$$\Omega_{DM} / \Omega_b = 5.3 \pm 0.1$$



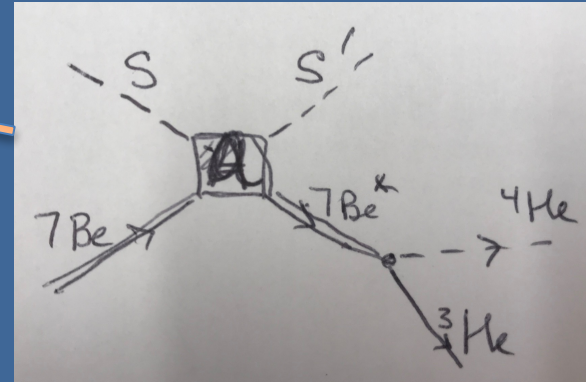
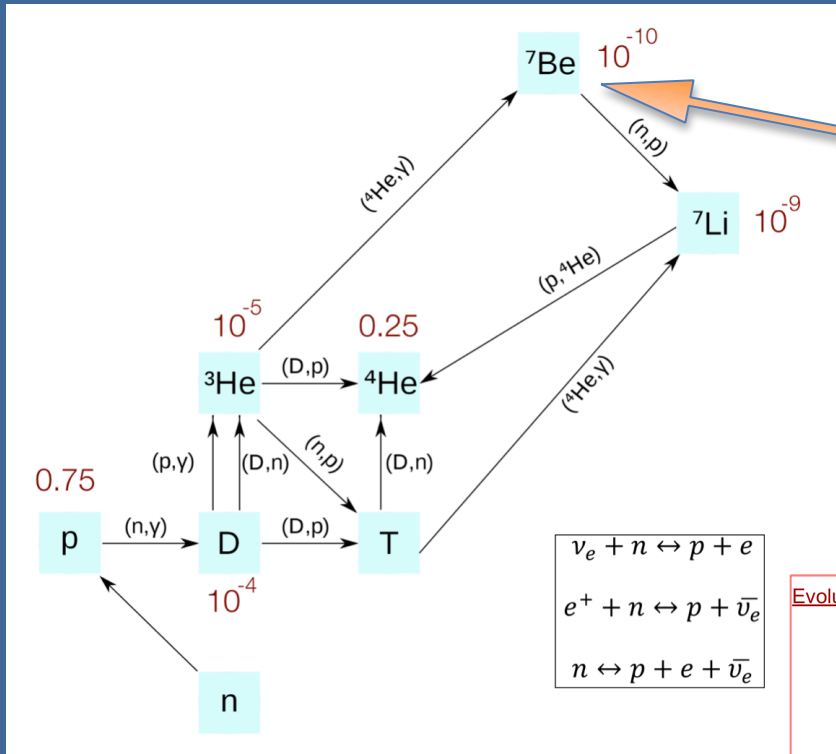
Prediction also applies to strange quark nuggets...

BBN's problem with primordial ${}^7\text{Li}$

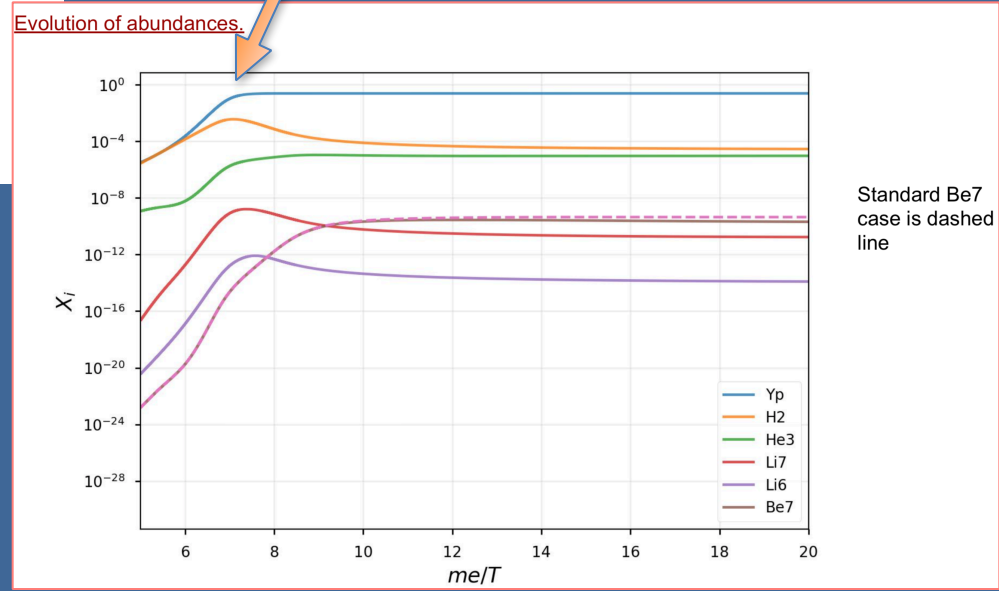
- Big Bang Nucleosynthesis works brilliantly *except 10σ problem*
 - Predicted abundance of ${}^7\text{Li} = (5.61 \pm 0.26) 10^{-10}$
 - Observed abundance of ${}^7\text{Li} = (1.58 \pm 0.31) 10^{-10}$
- Discrepancy is now very serious:
 - Nuclear rates all well-measured
 - $\eta = n_b/n_\gamma = (6.58 \pm 0.02) 10^{-10}$ from CMB
 - Astrophysics now secure (Spite plateau):
 - small scatter
 - ${}^7\text{Li}$ constant over > 3 decades of low metallicity
- **S solves the puzzle** (GRF + Richard Galvez, in preparation)
 - No other (reasonable) solution known



S dark matter breaks up ${}^7\text{Li}$ & ${}^7\text{Be}$



The "action" is at $T \sim 100$ keV so S only affects weakly bound nuclei



- Solves the puzzle

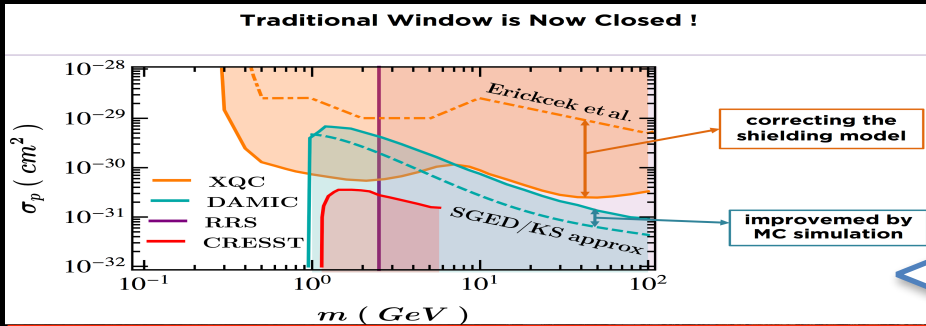
- Doesn't affect He or d

KE threshold for breakup =

$1.58, 2.46, 4.47, 5.75, 19.3$ [2.2] MeV
 ${}^7\text{Be}$ ${}^7\text{Li}$ ${}^3\text{He}$ T ${}^4\text{He}$ [d]

Stable S as Dark Matter

$10^{-26} - 10^{-25} \text{ cm}^2$



If Born Approximation good & XQC efficiency is perfect

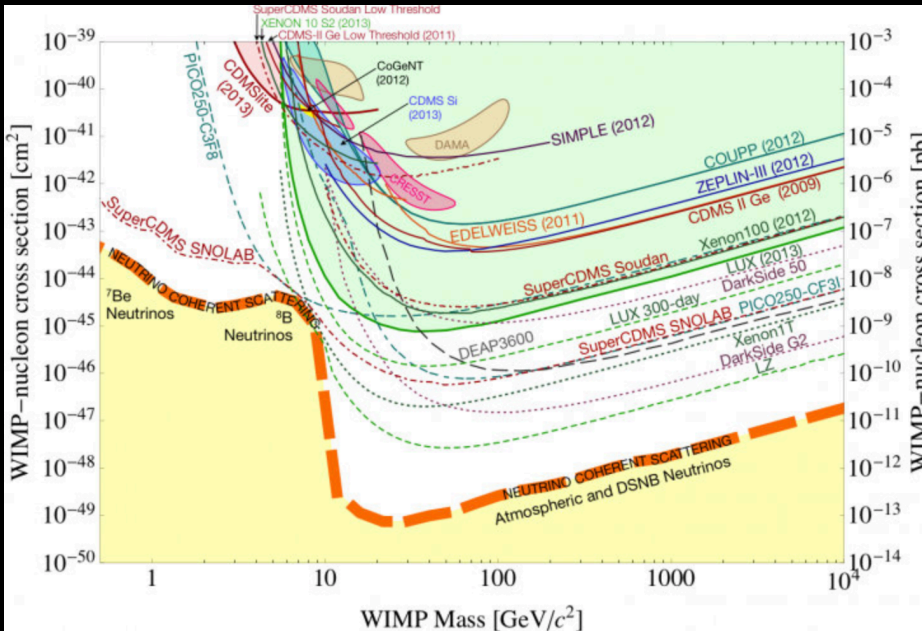
Closing the window on $\sim \text{GeV}$ Dark Matter with moderate ($\sim \mu\text{b}$) interaction with nucleons

M. Shafi Mahdawi and Glennys R. Farrar

Center for Cosmology and Particle Physics, Department of Physics, New York University,
4 Washington Place, New York, NY 10003, USA
E-mail: shafi.mahdawi@nyu.edu, g25@nyu.edu

Abstract. We improve limits on the spin-independent scattering cross section of Dark Matter on nucleons, for DM in the 300 MeV – 100 GeV mass range, based on the DAMIC and XQC experiments. Our results close the window which previously existed in this mass range, for a DM-nucleon cross section of order $\sim \mu\text{b}$, assuming the standard velocity distribution.

Shielded (e.g. underground) detectors are not sensitive (energy loss)



Dark Matter with Hadronic Interactions

(GRF + Xingchen Xu, to appear shortly)

$$V(r) = \frac{\alpha}{r} e^{-r m_\phi}$$

$m_\phi = 1$ GeV (flavor-singlet ω - ϕ combo), sourced by p or A

- v/c (DM) $\sim 10^{-3}$ 10^3 km/s (galaxy clusters) down to 1 km/s (atm & $z = 17$)
 - must solve Schroedinger Eqn. **Born approximation generically fails badly**
 - cross section depends only on combos $a = \frac{v}{2\alpha}$ and $b = \frac{2\alpha\mu}{m_\phi}$

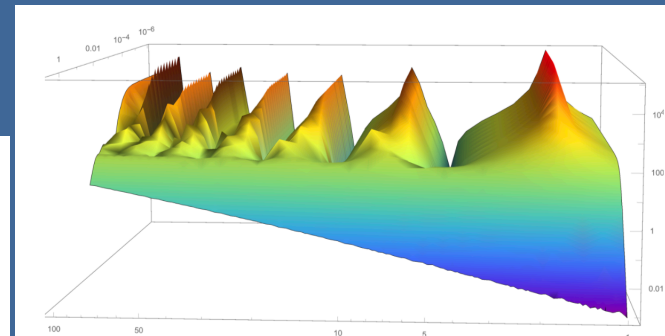


FIG. 2: 3D plot of σm_ϕ^2 in the a, b plane; b increases to the left and a decreases toward the back.

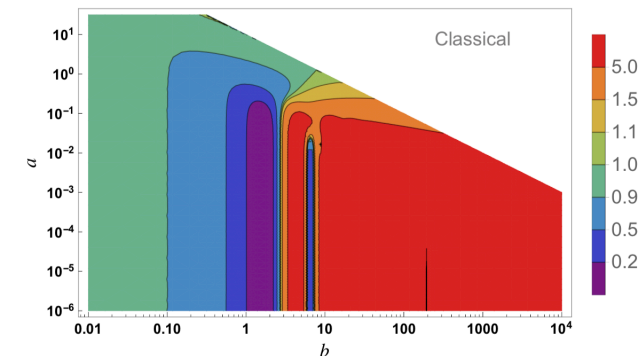


FIG. 3: Ratio of Born Approximation and Schroedinger Equation

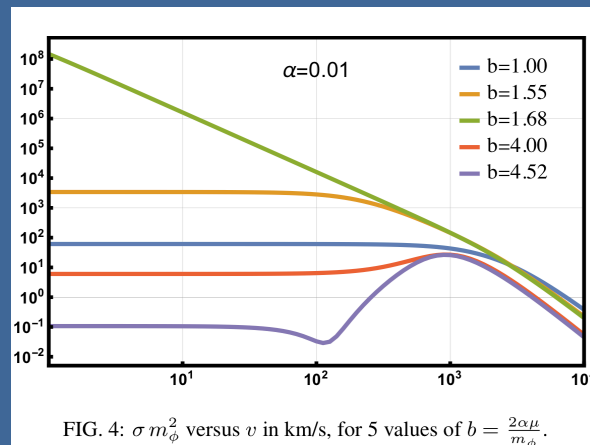


FIG. 4: σm_ϕ^2 versus v in km/s, for 5 values of $b = \frac{2\alpha\mu}{m_\phi}$.

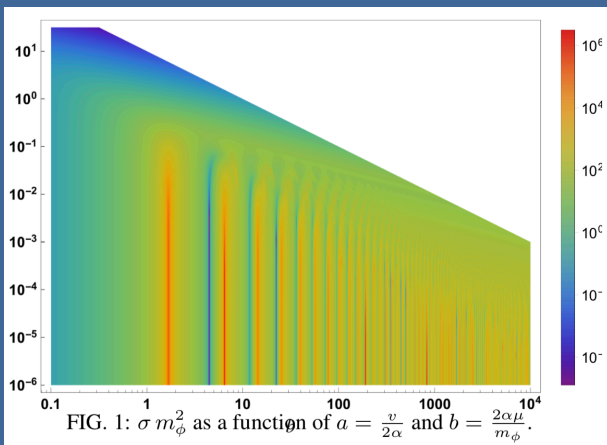
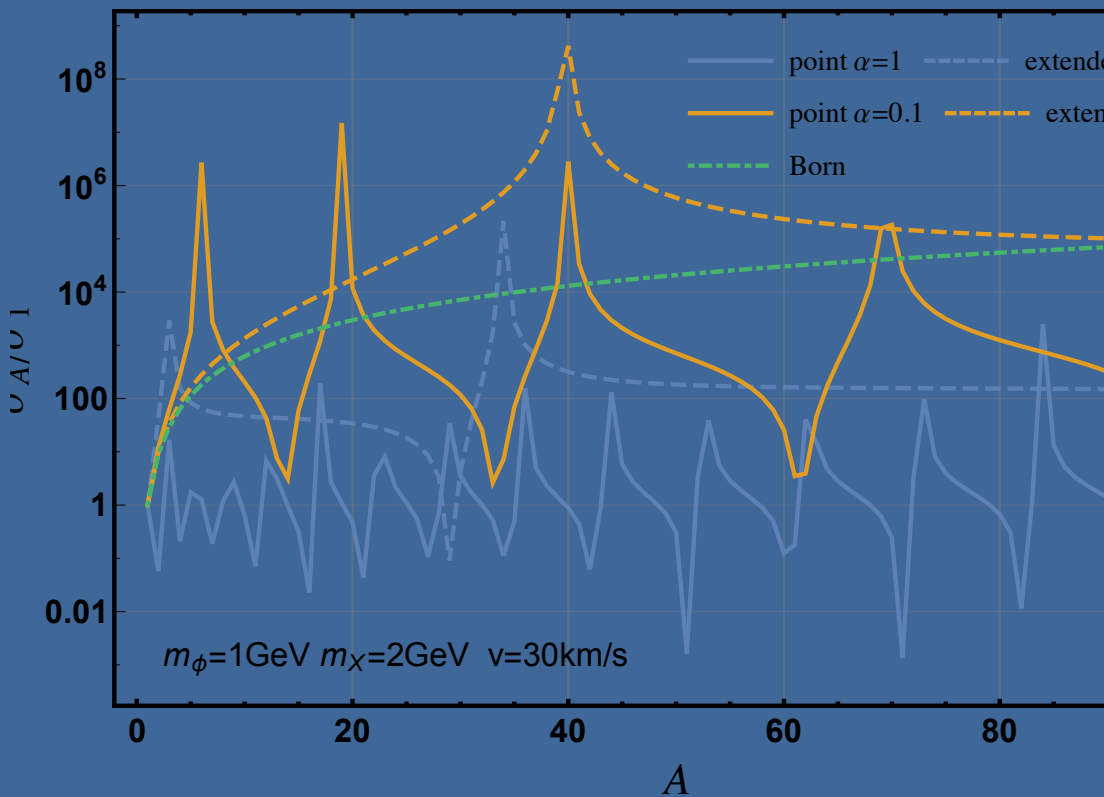


FIG. 1: σm_ϕ^2 as a function of $a = \frac{v}{2\alpha}$ and $b = \frac{2\alpha\mu}{m_\phi}$.

Plenty of Room for HDM, for now...

(GRF + Xingchen Xu, to appear shortly)

Caution: A -dependence very sensitive to nuclear form factor. Born approximation often misleading, by orders of magnitude.



Allowed regions of coupling from XQC (best Direct Detection)

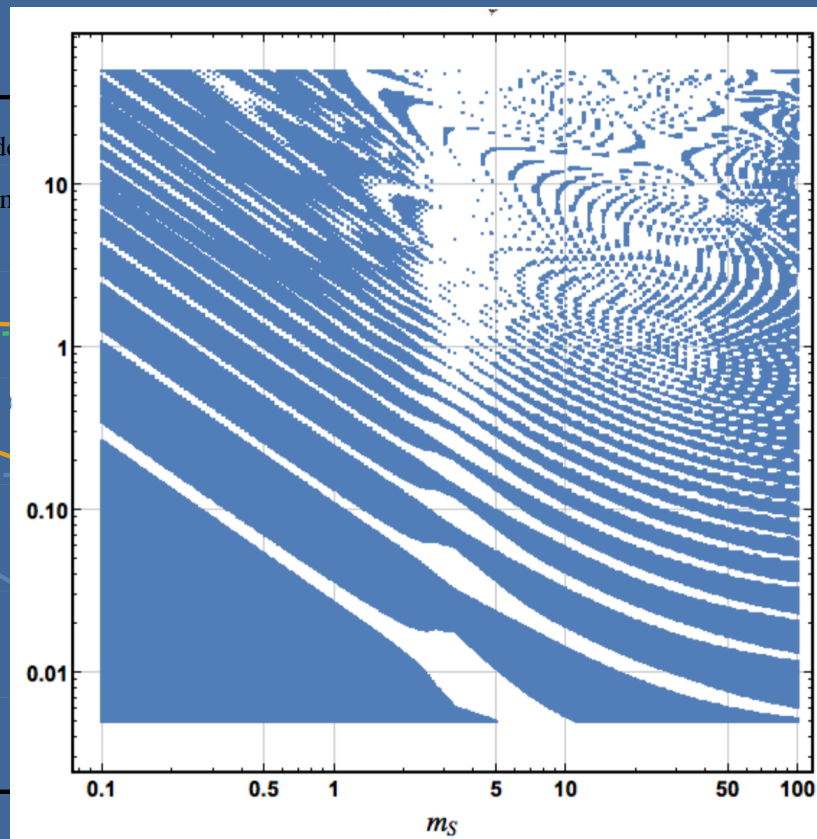


FIG. 7: Allowed regions (blue) in the coupling-DM mass plane α (vertical axis) and m_{DM} in GeV (horizontal axis) from XQC using

S has not been discovered at accelerators because it is elusive

- **Many negative searches, but all are inapplicable.** They either:
 - looked for H-dibaryon through decays (but S is stable)
 - restricted to mass > 2 GeV (but $m_S < 2$ GeV)
 - required fast production in $S=-2$ hypernuclei (but small overlap with baryons)

- **Wavefunction overlap with baryons is very small.** Extremely rare fluctuation required for $S \leftrightarrow \Lambda\Lambda$; $S \leftrightarrow NN$ is G_F^4 smaller \Rightarrow
 - nuclei can be stable ($\tau > 10^{29}$ yr) even for $m_S > 2 m_p$
 - hard to produce in fixed target experiments

- **S is similar to (much more copious) neutrons**

- **Promising accelerator detection strategies**

- **Apparent lack of baryon number and strangeness conservation:**

- $\Delta B = \pm 2$ with $\Delta S = \mp 2$

- **Reconstruct missing mass, e.g.:**

- $\Upsilon \rightarrow \Lambda \Lambda \bar{S}$ (+ pions) $M_S^2 = (p_\Upsilon - p_{\Lambda 1} - p_{\Lambda 2} - \sum p_{\pi_i})^2$

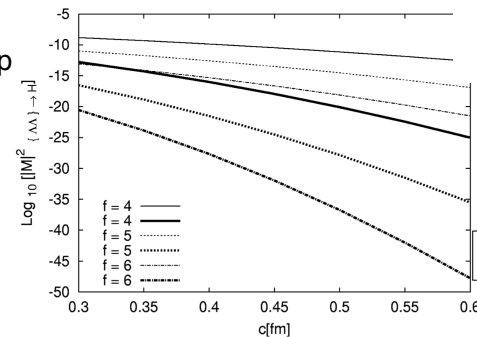
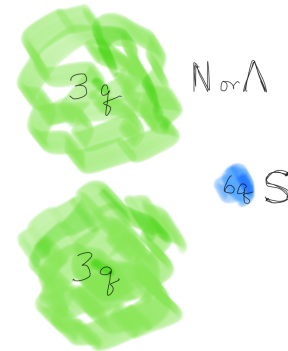


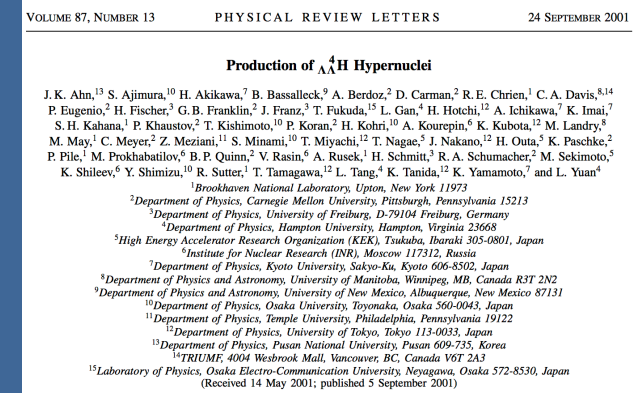
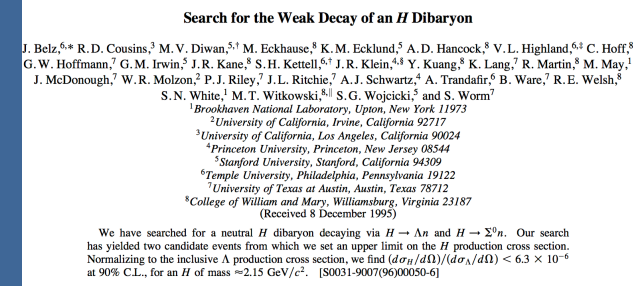
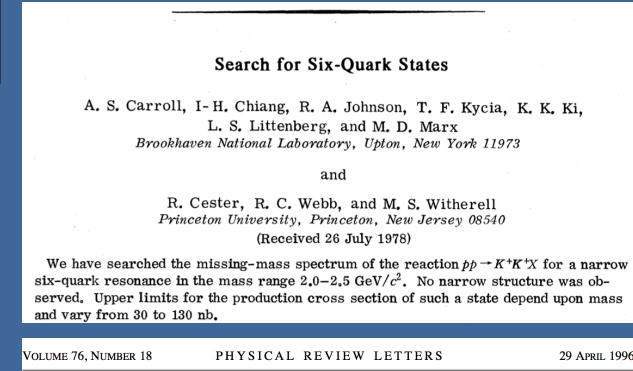
FIG. 1. Log_{10} of $|M|_{\Lambda\Lambda\text{-H}}^2$ versus hard core radius in femtometers, for ratio $f=R_H/R_H$ and two values of the Isgur-Karl oscillator parameter: $\alpha_B=0.406$ GeV (thick lines) and $\alpha_B=0.221$ GeV (thin lines).

GRF+G.Zaharijas 2004

Experimental searches so far

Looking for Jaffe's H-dibaryon (same QN but assumed to be unstable and $r \sim 1$ fm)

- Require $M > 2$ GeV:
 - Gufstafson+ FNAL1976 : Beam-dump + tof *Limit on production of neutral stable strongly interacting particle with mass > 2 GeV.*
 - Carroll+ BNL 1978: No narrow missing mass peak above 2 GeV in $pp \rightarrow K K X$
- Require H-dibaryon decay:
 - Badier+ NA3 1986
 - Bernstein+ FNAL 1988: Limit on production of neutral with $10^{-8} < \tau < 2 \times 10^{-6}$ s
 - Belz+ BNL 1996: $H \rightarrow \Lambda n$ or Σn [c.f., issue raised by L. Littenberg]
 - Kim+ Belle 2013: no narrow resonance in $\Upsilon \rightarrow \Lambda p K$
- Limits from production in doubly-strange hypernuclei:
 - Ahn+ BNL 2001
 - Takahashi+ KEK 2001



Experimental Searches

- Require $M > 2$ GeV:
 - Gufstafson+ FNAL 1976 : Beam-dump + tof *Limit on production of neutral stable strongly interacting particle with mass > 2 GeV.*
 - Carroll+ BNL 1978: No narrow missing mass peak above 2 GeV in $pp \rightarrow K K X$
- Require H-dibaryon decay:
 - Badier+ NA3 1986
 - Bernstein+ FNAL 1988: Limit on production of neutral with $10^{-8} < \tau < 2 \times 10^{-6}$ s
 - Belz+ BNL 1996: $H \rightarrow \Lambda \pi$ or $\Sigma \pi$ [c.f., issue raised by L. Littenberg]
 - Kim+ Belle 2013: no narrow resonance in $\Upsilon \rightarrow \Lambda \pi K$
- Limits from production in doubly-strange hypernuclei:
 - Ahn+ BNL 2001
 - Takanashi+ KEK 2001

Search for Six-Quark States

A. S. Barro, H. C. Cheng, J. R. Johnson, T. F. Kycia, K. K. Ki, L. S. Littenberg, and M. D. Marx
Brookhaven National Laboratory, Upton, New York 11973

and

R. C. Ester, R. C. Webb, and M. S. Witherell
Princeton University, Princeton, New Jersey 08540
 (Received 26 July 1978)

We have searched the missing-mass spectrum of the reaction $pp \rightarrow K^* K^* X$ for a narrow six-quark resonance in the mass range 2.0–2.5 GeV/c². No narrow structure was observed. Upper limits for the production cross section of such a state depend upon mass and vary from 30 to 130 nb.

VOLUME 76, NUMBER 18 PHYSICAL REVIEW LETTERS 29 APRIL 1996

Search for the Weak Decay of an H Dibaryon

J. Belz,^{6,*} R. D. Cousins,³ M. V. Diwan,^{5,1} M. Eckhause,⁸ K. M. Ecklund,⁵ A. D. Hancock,⁵ V. L. Highland,^{6,4} C. Hoff,⁵ G. W. Hoffmann,² G. M. Irwin,³ J. R. Kane,⁵ S. H. Kettel,^{4,1} J. R. Klein,^{4,4} Y. Kuang,⁸ K. Lang,⁷ R. Martin,⁸ M. May,¹ J. McDonough,⁷ W. R. Molzon,³ P. J. Riley,⁷ J. L. Ritchie,⁷ A. J. Schwartz,⁴ A. Trandafir,⁹ B. Ware,⁷ R. E. Welsh,³ S. N. White,¹ M. T. Witkowski,¹⁰ S. G. Wojcicki,² and S. Worn¹

¹Brookhaven National Laboratory, Upton, New York 11973
²University of California, Irvine, California 92717
³University of California, Los Angeles, California 90024
⁴Princeton University, Princeton, New Jersey 08544
⁵Stanford University, Stanford, California 94309
⁶Temple University, Philadelphia, Pennsylvania 19122
⁷University of Texas at Austin, Austin, Texas 78712
⁸College of William and Mary, Williamsburg, Virginia 23187
 (Received 8 December 1995)

We have searched for a neutral H dibaryon decaying via $H \rightarrow \Lambda n$ and $H \rightarrow \Sigma^0 n$. Our search has yielded two candidate events from which we set an upper limit on the H production cross section. Normalizing to the inclusive Λ production cross section, we find $(d\sigma_H/d\Omega)/(d\sigma_\Lambda/d\Omega) < 6.3 \times 10^{-6}$ at 90% C.L., for an H of mass ≈ 2.15 GeV/c². [S0031-9007(96)0050-6]

VOLUME 87, NUMBER 13 PHYSICAL REVIEW LETTERS 24 SEPTEMBER 2001

Production of $\Lambda\Lambda^4$ Hypernuclei

J. K. Ahn,¹³ S. Ajimura,¹⁰ H. Akiyawa,⁷ B. Bassalleck,⁹ A. Berdoz,⁷ D. Carman,² R. E. Chrien,¹ C. A. Davis,^{8,14} P. Eugenio,² H. Fischer,³ G. B. Franklin,² J. Franz,² T. Fukuda,¹² L. Gan,¹ H. Hotchi,¹² A. Ichikawa,⁷ K. Imai,⁷ S. H. Kahana,¹ P. Khastov,² T. Kishimoto,¹⁰ P. Koran,² H. Kohri,¹⁰ A. Kourepin,⁵ K. Kubota,¹² M. Landry,⁸ M. May,¹ C. Meyer,² Z. Meziani,¹¹ S. Minami,¹⁰ T. Miyachi,¹² T. Nagae,² J. Nakano,¹² H. Ota,⁵ K. Paschke,² P. Pile,¹ M. Prokhabatilo,⁶ B. P. Quinn,² V. Rasin,⁶ A. Rusek,¹ H. Schmitt,³ R. A. Schumacher,² M. Sekimoto,³ K. Shileev,⁶ Y. Shimizu,¹⁰ R. Sutter,¹ T. Tamagawa,¹² L. Tang,⁴ K. Tanida,¹² K. Yamamoto,² and L. Yuan¹

¹Brookhaven National Laboratory, Upton, New York 11973
²Department of Physics, Carnegie Mellon University, Pittsburgh, Pennsylvania 15213
³Department of Physics, University of Freiburg, D-79104 Freiburg, Germany
⁴Department of Physics, Hampton University, Hampton, Virginia 23668
⁵High Energy Accelerator Research Organization (KEK), Tsukuba, Ibaraki 305-0801, Japan
⁶Institute for Nuclear Research (INR), Moscow 117312, Russia
⁷Department of Physics, Kyoto University, Sakyo-Ku, Kyoto 606-8502, Japan
⁸Department of Physics and Astronomy, University of Manitoba, Winnipeg, MB, Canada R3T 2N2
⁹Department of Physics and Astronomy, University of New Mexico, Albuquerque, New Mexico 87131
¹⁰Department of Physics, Osaka University, Toyonaka, Osaka 560-0043, Japan
¹¹Department of Physics, Temple University, Philadelphia, Pennsylvania 19122
¹²Department of Physics, University of Tokyo, Tokyo 113-0033, Japan
¹³Department of Physics, Pusan National University, Pusan 609-735, Korea
¹⁴TRIUMF, 4004 Wesbrook Mall, Vancouver, BC, Canada V6T 2A3
 (Received 14 May 2001; published 5 September 2001)

An experiment demonstrating the production of double- Λ hypernuclei in (K^-, K^+) reactions on ^9Be was carried out at the D6 line in the BNL alternating-gradient synchrotron. The technique was the observation of pions produced in sequential mesonic weak decay, each pion associated with one unit of strangeness change. The results indicate the production of a significant number of the double hypernucleus $\Lambda\Lambda^4\text{H}$ and the twin hypernuclei $\Lambda^4\text{H}$ and $\Sigma^4\text{H}$. The relevant decay chains are discussed and a simple model of the production mechanism is presented. An implication of this experiment is that the existence of an $S = -2$ dibaryon more than a few MeV below the $\Lambda\Lambda$ mass is unlikely.

Cosmology & structure formation

- **DM-baryon interaction: momentum transfer => *slight drag on DM during structure formation***
 - Dvorkin, Blum, Kamionkowski (2014):
 - **Ly-alpha forest: $\sigma < \sim 10$ mb if v-indept — no problem for S**
 - Buen-Abad, Marques-Tavares, Schmaltz (2015):
 - **momentum transfer helps reconcile H_0 & σ_8**
 - Boring or an opportunity? To be determined...
- **S-S self interactions + S-baryon interactions:**
 - may have similar benefits as Self Interacting DM
 - core-cusp, “too-big-to-fail” & missing sub-halos problems.

Galaxies & Clusters

DM-gas scattering provides a source of heating, needed for

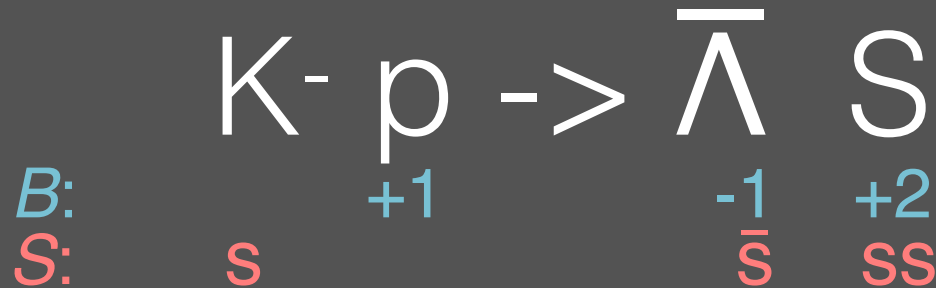
- Milky Way's extended hot gas halo — 2×10^6 K
- Quenching star formation
- Avoiding “cooling flow catastrophe” in X-ray clusters

Key points to take home

- ***There may a tightly bound 6-quark state $S = uuddss$***
 - **Unique, symmetric structure \Rightarrow other hadrons don't provide guidance**
 - mass is not driven by chiral symmetry breaking (unlike baryons)
 - constituent quark model probably completely misleading
 - ***If $M_S < 2 m_p + 2 m_e$, S is absolutely stable***
- ***If S is stable, its an excellent Dark Matter candidate***
 - Relic abundance is natural. **EXPLAINS 7Li Discrepancy in BBN and Dark Matter to baryon ratio**
 - Usual WIMP detection strategy isn't applicable.
 - May reconcile tension in H_0 & σ_8 and explain astrophysics puzzles ("quenching", core-cusp, DM rotation curves...)
- ***S may be waiting to be discovered in existing Υ -decays or LHC experiments... mass can be accurately measured in Υ -decay exclusive final states.***
- ***SDM will be challenging to detect, but not impossible. Astrophysical and cosmological effects may allow it to be constrained, excluded or confirmed.***

Backup Slides

Classic Approach: would be great, but *very* low rate due to low overlap



- $\bar{\Lambda}$ is a gold-plated signature : $\bar{\Lambda} \rightarrow \pi^+ \bar{p}$
 - Easy to ID & reconstruct 4-momentum
 - $c\tau = 8 \text{ cm}$ all $\bar{\Lambda}$ are ID'd
- S : undetected, but 4 momentum determined
 - $p_S = p_K + p_p - p_{\bar{\Lambda}}$
 - NA61: est. $\sim 20 \text{ MeV}$ accuracy on “missing-mass” of S
 - For $p_{\text{beam}} < 5.35 \text{ GeV}/c$, no conventional source of $\bar{\Lambda}$'s
- NA61: $9 \text{ GeV}/c$ K^- beam

Sexaquark Discovery Strategy

- Apparent lack of B and S conservation:

- $\Delta B = \pm 2 + \Delta S = \mp 2$

- Reconstruct missing mass, e.g.:

- $\Upsilon \rightarrow \Lambda \Lambda \bar{S}$ (+ pions) $M_S^2 = (p_\Upsilon - p_{\Lambda 1} - p_{\Lambda 2} - \sum p_{\pi_i})^2$

- LHC: $\bar{S} + N \rightarrow \bar{\Lambda} K^+$ $M_S^2 = (p_{\bar{\Lambda}} + p_K - p_N)^2$

- Snolab nuclei: $pn \rightarrow S e^+ \nu$ $G_F^4 \tau > 10^{+29}$ yr

$$\Upsilon \rightarrow \Lambda \Lambda \bar{S} \quad \& \quad \bar{\Lambda} \bar{\Lambda} S$$

(+ pions)

- Υ is *localized* source of ggg

\Rightarrow production of S is (relatively) enhanced

- Many $\times 10^8$ events collected (CLEO, Babar, Belle)

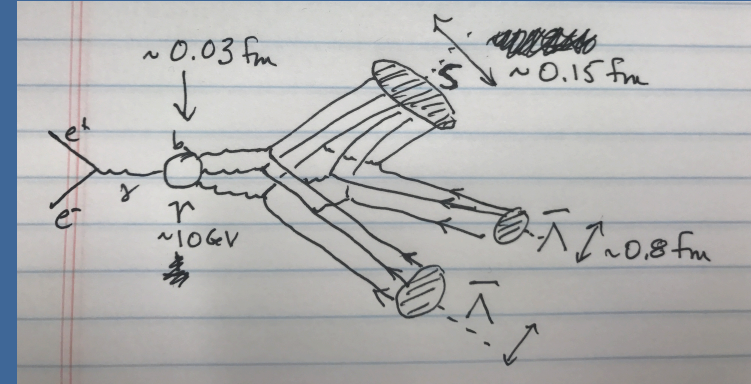
- detectors pretty hermetic, have good mass resolution, $O(10 \text{ MeV})$
- Λ decays quickly to $p\pi$ so easy to ID. $c\tau = 8 \text{ cm}$

- Can MEASURE m_s via missing mass

- *Very clean*

- Main bkg is $K_S K_S K_L K_L$ (+ pions)
 - K_S 's mis-ID'd as Λ 's and K_L 's escaping before decay : *negligible for Belle*
 - rare and can model accurately
 - $K_S K_S K_L K_L$ (+ pions) *is measurable*, from $K^+ K^+ K^- K^-$ (+ pions)
- “Conspiracy” of missed particles producing $\Delta B = \pm 2$, $\Delta S = \mp 2$ very hard

Background does not have narrow peak in missing mass!



LHC

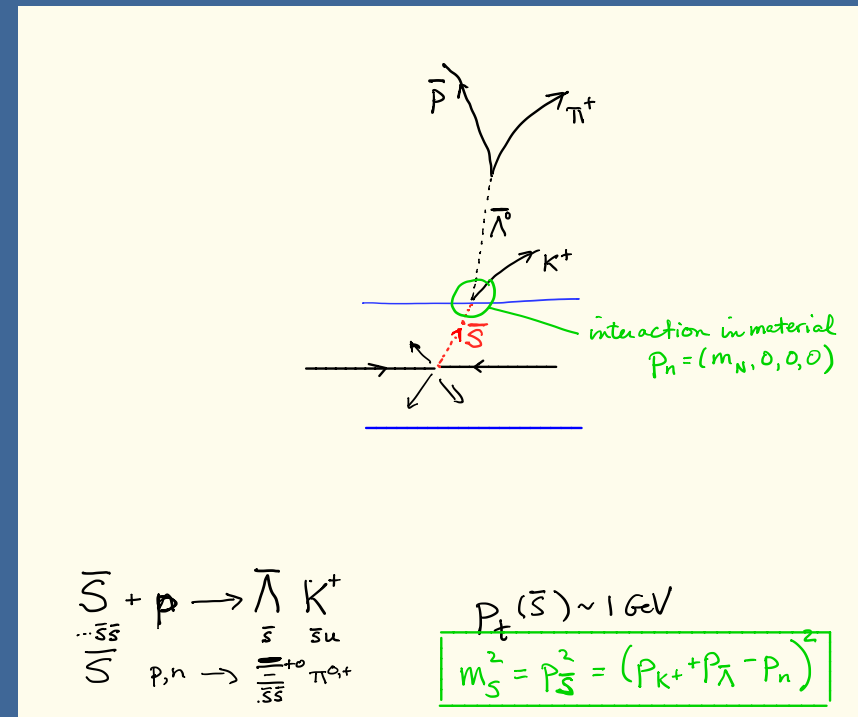
- Hadronic collisions: low production rate due to small wfn overlap
- Find a needle in a haystack (10^{11} recorded events; potential for trigger $\gg \gg$ more)
- Statistical examination of correlation between

$$\Delta B = \pm 2, \Delta S = \mp 2$$

- **\bar{S} annihilation in tracker, tag by**

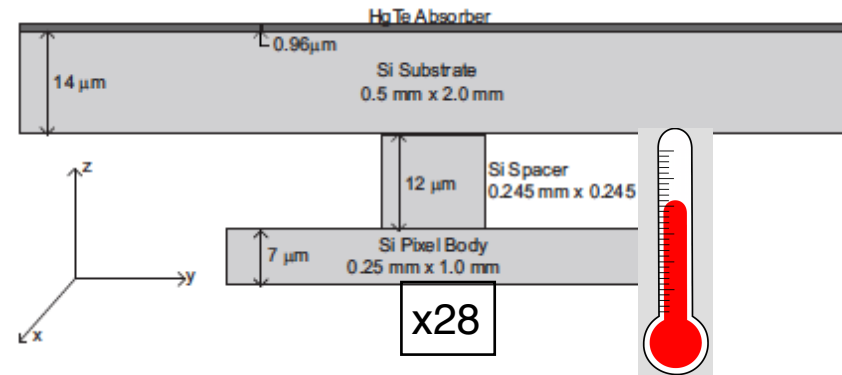
$$\bar{E}^{+,0} \rightarrow \bar{\Lambda} \pi^{+,0} \quad (c\tau = 5\gamma \text{ cm}) \quad \bar{\Lambda} \rightarrow \bar{p} \pi^+$$

or $\bar{\Lambda} K^+$



- 2nd exponential in scattering-length distribution of n-like interactions, due to S

X-Ray Quantum Calorimeter (XQC)



- On sounding rocket, 200 km above earth
- Best limit for high x-secn (McCammon+02, Wandelt+02, GF+Zaharijas05, Erickcek+07, Mahdawi & GF 17)
 - sensitive to X-rays with $E \geq 29$ eV
 - 100 sec flight, ~ 100 events
 - nuclear recoil \Rightarrow X-rays, which thermalize (*assumption*)

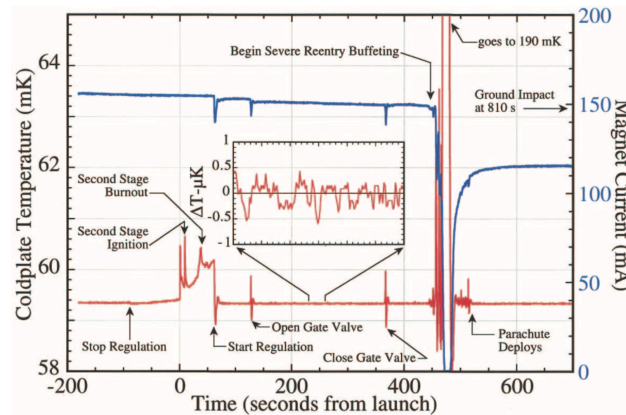
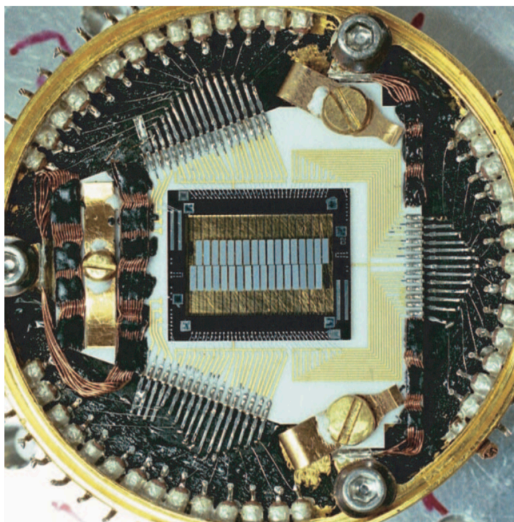
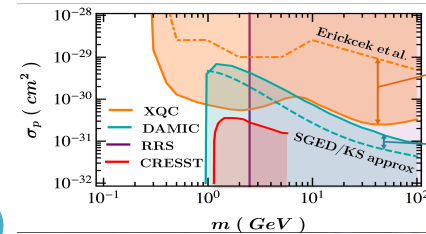


Fig. 7.—In-flight performance of the temperature control system, showing the coldplate temperature and magnet current. Temperature fluctuations during data taking are about 210 nK rms. The gate-valve motor is located on the vacuum jacket and caused the most serious thermal disturbance up to reentry. Accelerations during reentry exceeded 20 g with tumbling at ~ 1 Hz, introducing heat to the cold stage faster than it could be removed. Temperature regulation is recovered once tumbling stops, allowing calibration data to be obtained.

Calibrate with X-rays

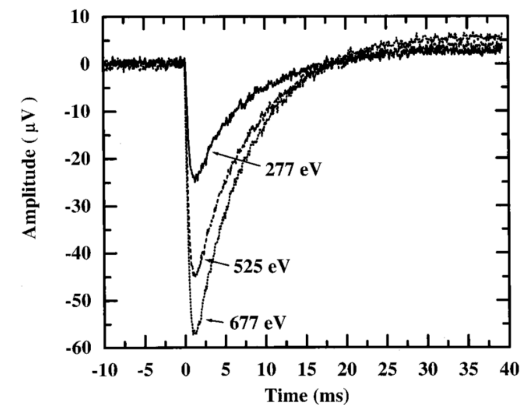


Fig. 8.—Unfiltered X-ray pulses from the gate-valve calibration source.

Closer look at XQC sensitivity

Silicon nucleus recoil: $KE_{\max} \sim 500 \text{ eV} \Rightarrow v_{\max} \sim 20 \text{ km/s} \ll v_e \Rightarrow$

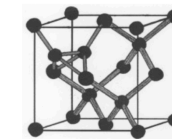
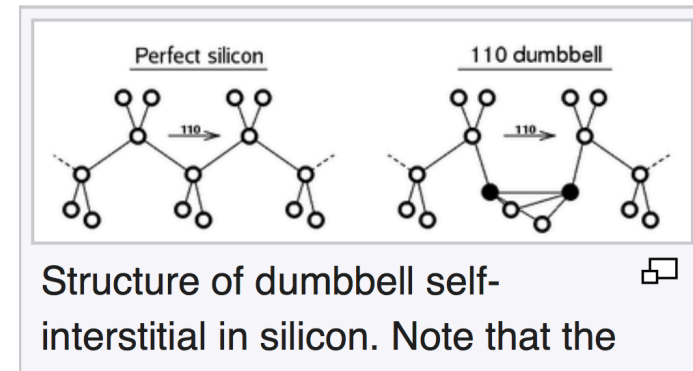
- atomic interaction is adiabatic \Rightarrow
- negligible ionization.

Si atom moving in semiconductor crystal:

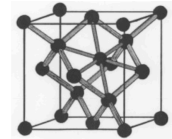
- rearranges covalent bonds
- produces interstitial defects
- 500eV atom produces Frenkel pairs (V+I)
 - $E_{\text{Fp_min}} = 5 \text{ eV}$
 - $E_{\text{migration}} \sim 0.1 \text{ eV}$
- Cascade energy loss producing
 - $N \sim (KE_{\text{rec}} / 5 \text{ eV})$ Frenkel pairs,

$\leftarrow \sim 2\%$ of KE_{rec} goes to thermalization

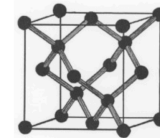
**$\Rightarrow KE_{\text{rec,min}} > 1.5 \text{ keV} \Rightarrow KE_{\text{DM,min}} > 6 \text{ keV} \Rightarrow$
 $v_{\text{DM,min}} > 300 \text{ km/s}$**



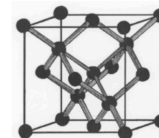
(a) Split-(110)



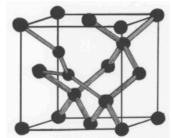
(b) Hexagonal



(c) Perfect crystal



(d) Tetrahedral



(e) Concerted exchange

Figure 1. Illustrations of the split-(110), hexagonal, and tetrahedral self-interstitial defects, together with the perfect crystal and the saddle point of Pandey's concerted exchange.