

CIPANP2018 Palm Spring

29 May 2018

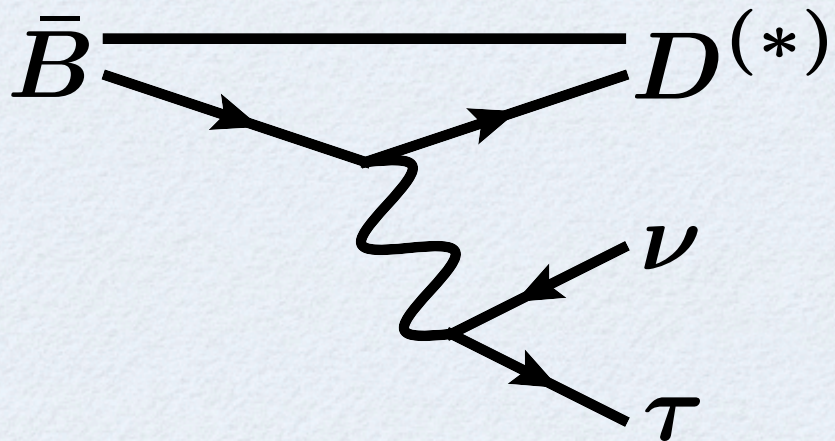
RD and RD*

Theoretical Developments

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are tree level processes

- Precise prediction on the **Ratio** has been done

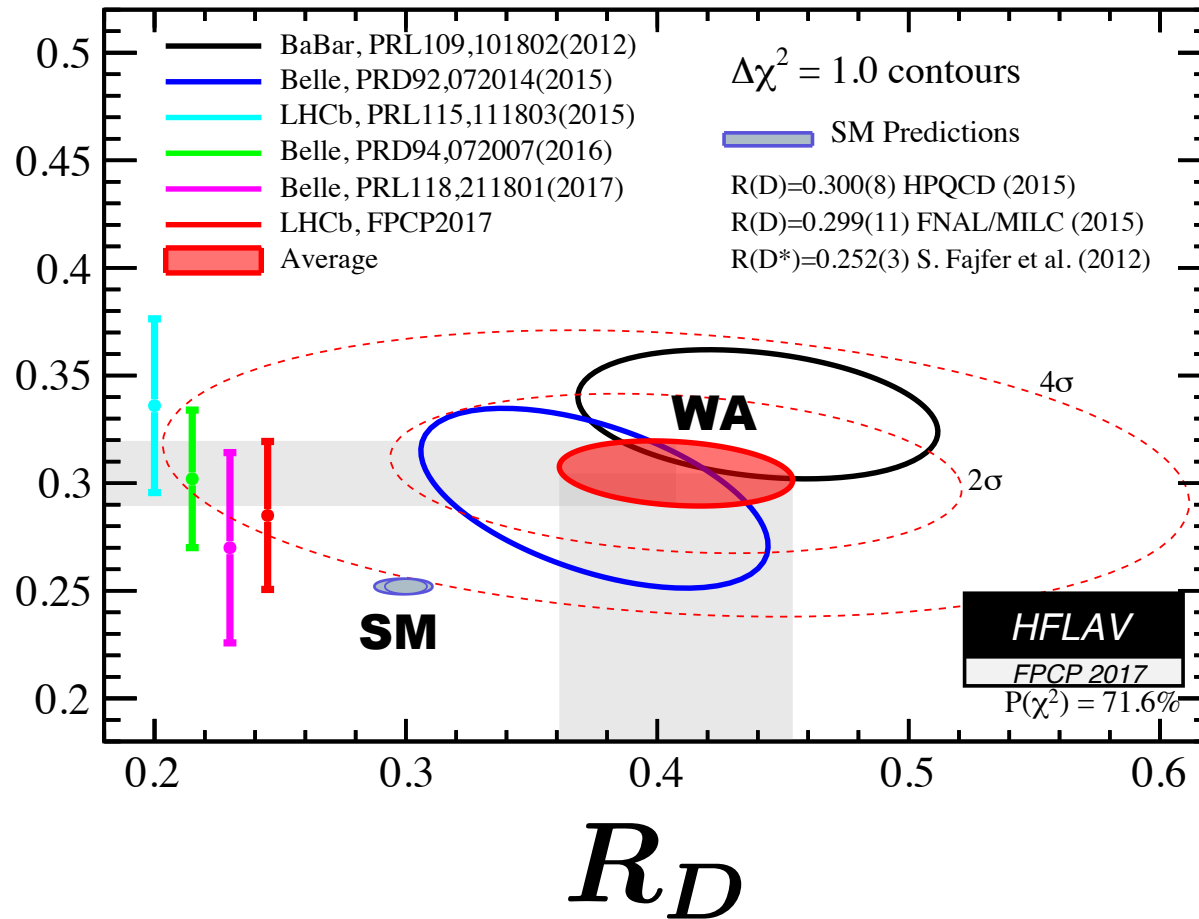
$$R_D = \frac{\Gamma(\bar{B} \rightarrow D\tau\bar{\nu})}{\Gamma(\bar{B} \rightarrow D\ell\bar{\nu})} = 0.299 \pm 0.003$$

$$R_{D^*} = \frac{\Gamma(\bar{B} \rightarrow D^*\tau\bar{\nu})}{\Gamma(\bar{B} \rightarrow D^*\ell\bar{\nu})} = 0.257 \pm 0.003$$

$\sim 1\%$

Nevertheless,

- the SM values are **NOT in agreement with data**

R_{D^*}  $\sim 4.1\sigma$

BaBar : PRL 109, 101802 (2012), PRD 88, 072012 (2013)

Belle : PRD 92, 072014 (2015), PRD 94, 072007 (2016), arXiv 1608.06391

LHCb : PRL 115, 111803 (2015), arXiv 1708.08856

Topics

- **SM predictions**
- **NP explanations**
- **Relevant observables**

SM predictions

[1] Form Factor

Main uncertainty in $RD(^*)$ comes from **Form Factors**

$$\langle D(v') | \bar{c} \gamma_\mu b | B(v) \rangle = \sqrt{m_B m_D} \left[h_+(q^2) (v + v')_\mu + h_-(q^2) (v - v')_\mu \right]$$

$$\frac{d\Gamma(B \rightarrow D \ell \nu)}{dw} = \frac{G_F^2 |V_{cb}|^2 \eta^2 m_B^5}{48\pi^3} (w(q^2)^2 - 1)^{3/2} r_D^3 (1 + r_D)^2 \mathcal{G}(q^2)^2$$

Using Heavy Quark Effective Theory, q^2 dependence can be described

$$\mathcal{G} = h_+ - \frac{1 - r_D}{1 + r_D} h_- = \xi_{IW}(q^2) + \frac{\alpha_s}{\pi} \chi_1(q^2) + \frac{\Lambda_{\text{QCD}}}{m_{c,b}} \chi_2(q^2) + \dots$$

The functions are then determined with QCD sum-rule / lattice

+ fit to data of the light lepton mode.

QCDSR + lattice QCD + Fit to Belle data of $B \rightarrow D(*)\ell\nu$ ($\ell = e, \mu$)

up to the NLO, i.e. $O(\alpha_s), O(1/m_Q)$

Ligeti et al., 1703.05330

Scenario	$R(D)$	$R(D^*)$	Correlation
$L_{w=1}$	0.292 ± 0.005	0.255 ± 0.005	41%
$L_{w=1}+SR$	0.291 ± 0.005	0.255 ± 0.003	57%
NoL	0.273 ± 0.016	0.250 ± 0.006	49%
NoL+SR	0.295 ± 0.007	0.255 ± 0.004	43%
$L_{w \geq 1}$	0.298 ± 0.003	0.261 ± 0.004	19%
$L_{w \geq 1}+SR$	<u>0.299 ± 0.003</u>	<u>0.257 ± 0.003</u>	44%
th: $L_{w \geq 1}+SR$	0.306 ± 0.005	0.256 ± 0.004	33%
Data [9]	<u>0.403 ± 0.047</u>	<u>0.310 ± 0.017</u>	-23%
Refs. [53, 57, 59]	0.300 ± 0.008	—	—
Ref. [58]	0.299 ± 0.003	—	—
Ref. [34]	—	0.252 ± 0.003	—

This study

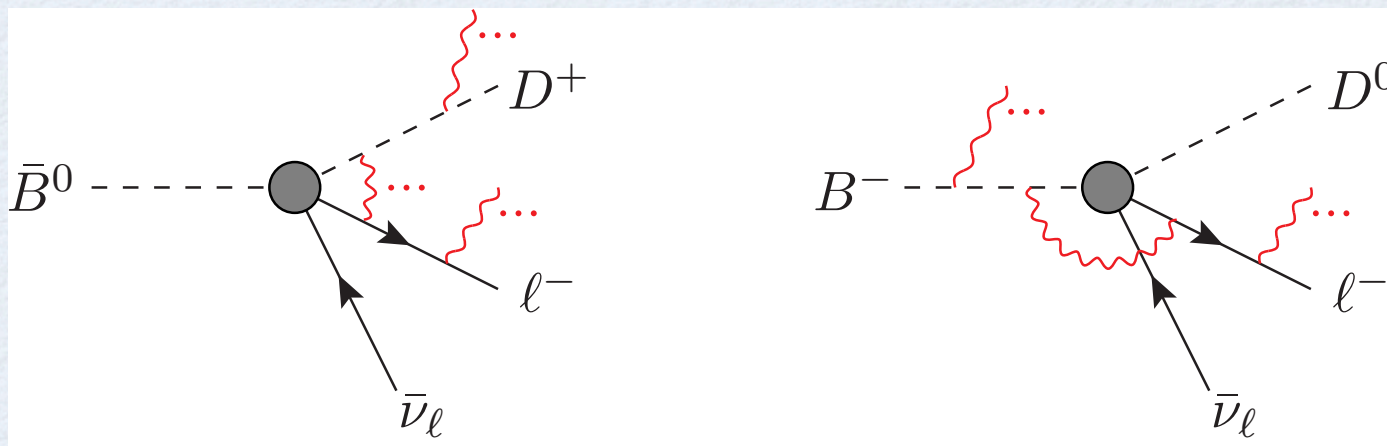
Measurements

Previous studies

[2] Radiative correction

Kitahara et al., 1803.05881

Another development. Soft-photon effects depend on lepton mass, which leads to **corrections** even in $RD(^*)$



Soft-photon corrections to RD result in

(1) leading to $RD^{*+} \neq RD^{*0}$

(2) depending on photon energy cut

(3) non-negligible **constructive contribution to RD** ,
at most 4~6%

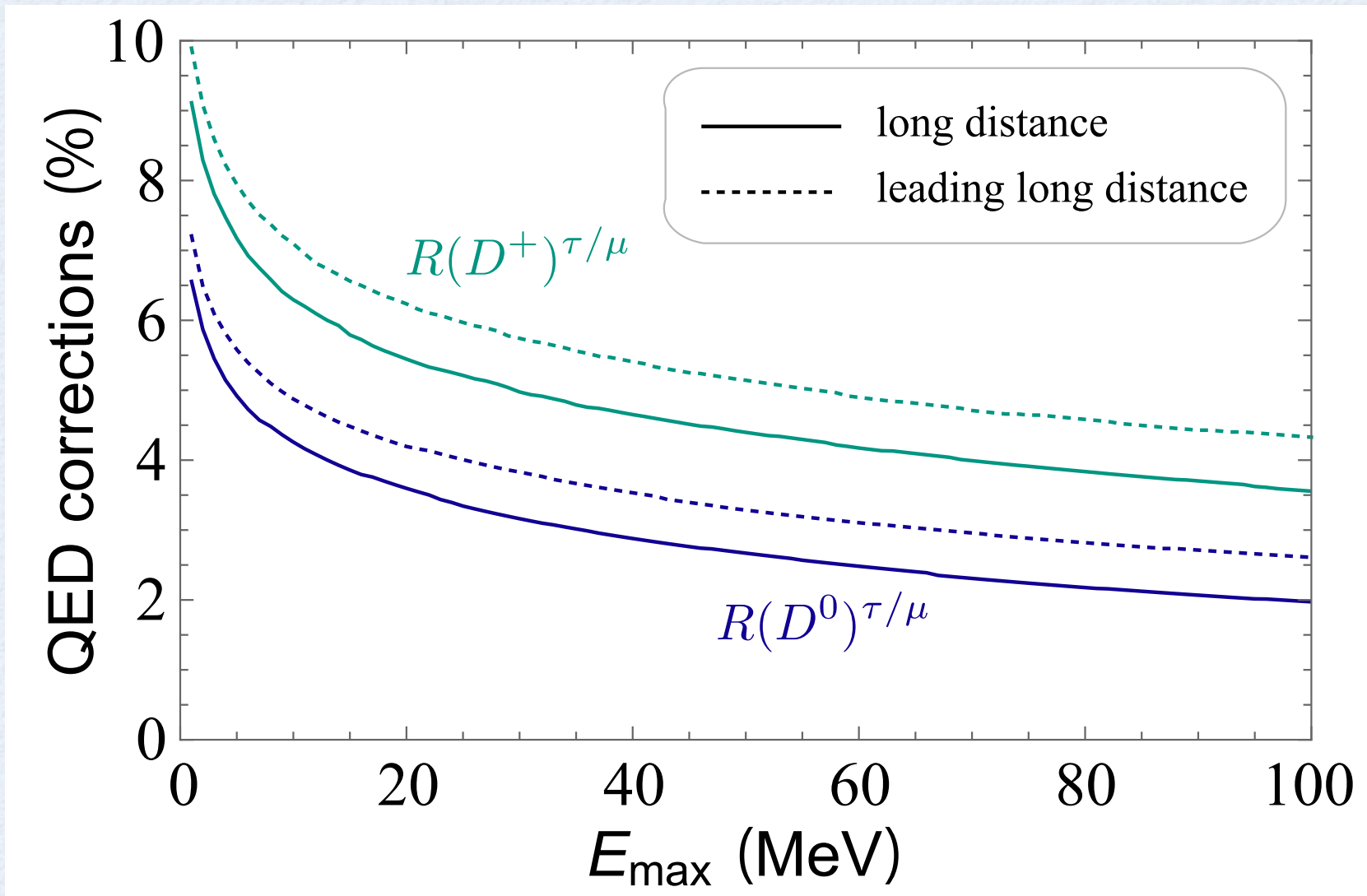
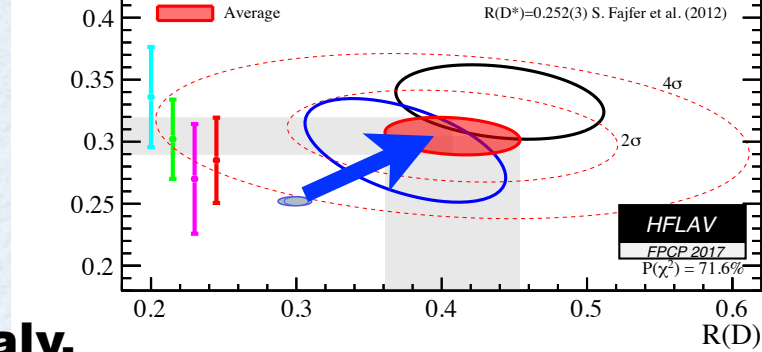


FIG. 3. The (leading) long-distance QED corrections to $R(D^+)_{\tau/\mu}$ and $R(D^0)_{\tau/\mu}$ as a function of E_{\max} .

NP explanations



There exist several solutions to the $R(D^*)$ anomaly.

In terms of effective operators, possible NPs are given as follows

$$\mathcal{L}_{\text{eff}}^{\text{NP}} \equiv -2\sqrt{2}G_F V_{cb} C_{\text{NP}} \mathcal{O}_{\text{NP}}$$

$$\mathbf{V-A} : \mathcal{O}_{V_1} = (\bar{c}\gamma^\mu P_L b)(\bar{\tau}\gamma_\mu P_L \nu)$$

Models : (SM), W' boson, Vector Leptoquark, ...

Fit to data : $C_{V_1} \sim +0.17$

- NP with 17% contribution of the SM value is required
- Assuming NP coupling =1, it implies $\sim 2\text{TeV}$ NP scale

V+A (quark sector) : $\mathcal{O}_{V_2} = (\bar{c}\gamma^\mu P_R b)(\bar{\tau}\gamma_\mu P_L \nu)$

Models : **W' boson, ...?**

Fit to data: $C_{V_2} \sim 0.01 + 0.6i$

- **Complex coupling is necessary**
-

Scalar types : $\mathcal{O}_{S_{1(2)}} = (\bar{c} P_{R(L)} b)(\bar{\tau} P_L \nu)$

Models : **Charged Higgs, Scalar Leptoquark, ...**

Fit to data: $C_{S_1} = \text{no solution}, \quad C_{S_2} \sim -1.5$

- **2HDM of type I is disfavored**
 - **Large scalar contribution is needed**
-

Tensor type : $\mathcal{O}_T = (\bar{c}\sigma^{\mu\nu} P_L b)(\bar{\tau}\sigma_{\mu\nu} P_L \nu)$

Models : **Doublet vector/scalar Leptoquark (in part)**

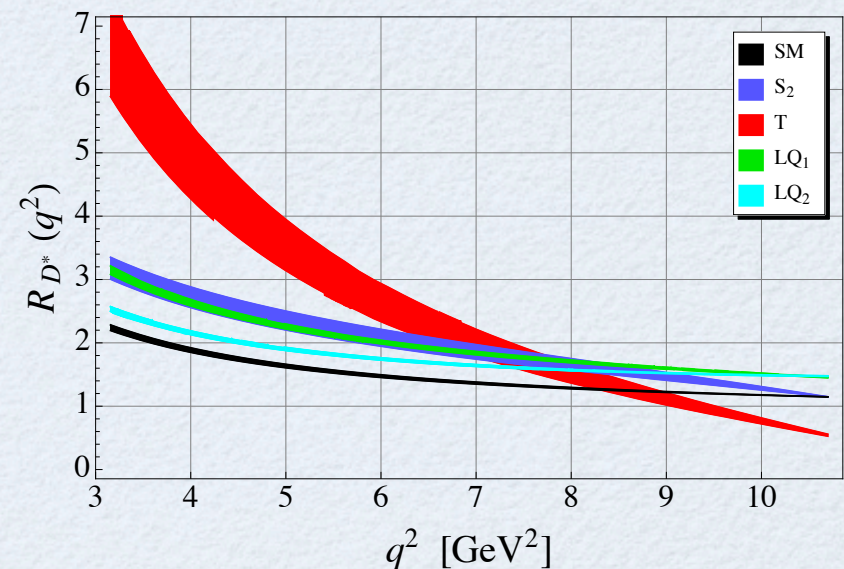
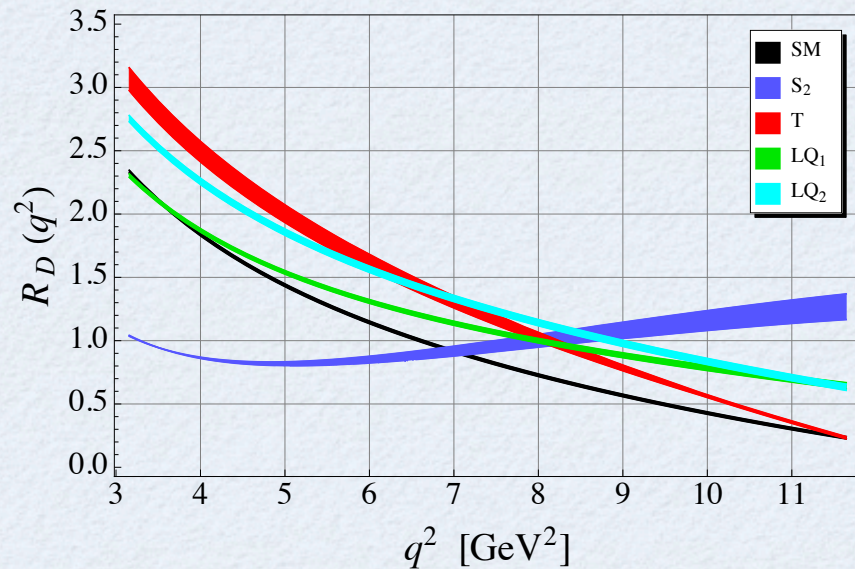
Fit to data: $C_T \sim 0.3$

Relevant observables

[1] q^2 distribution

RW et al. 1412.3761, B2TiP report

Distributions for the case that C_{NP} = best fit to the current results of $R_{D^{(*)}}$

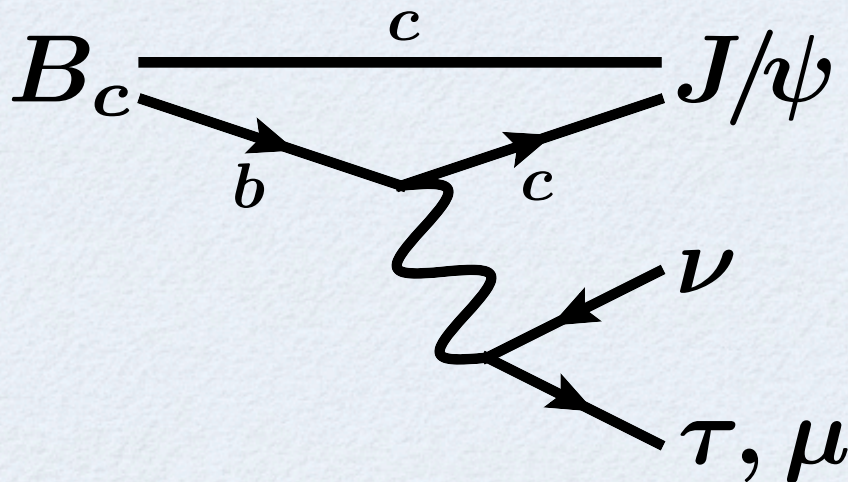


Some simple test with statistics was done and it turns out that

“**5ab⁻¹ data for q^2 distributions** enable us to distinguish the NP scenarios in case that the present anomalies still remain in the future”

[2] $B_c \rightarrow J/\psi \tau \nu$

RW, 1709.08644



has been observed at LHCb

$$R_{J/\psi} = \frac{\Gamma(B_c \rightarrow J/\psi \tau \bar{\nu})}{\Gamma(B_c \rightarrow J/\psi \mu \bar{\nu})}$$

data : $R_{J/\psi} = 0.71 \pm 0.17 \pm 0.18$

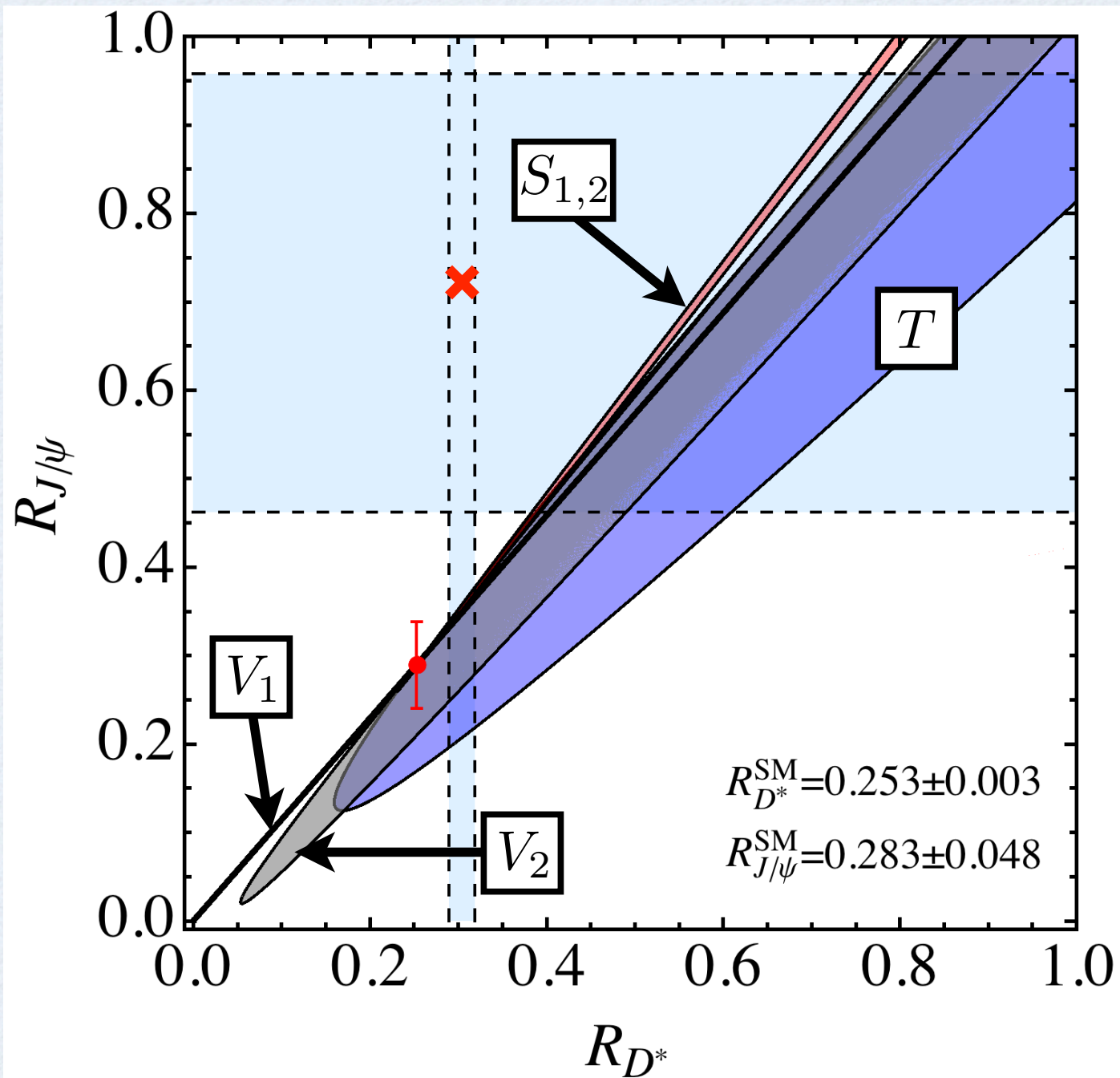
LHCb, 1711.05623

SM : $R_{J/\psi} = 0.283 \pm 0.048$

RW, 1709.08644

- **Perturbative QCD analysis provides the form factor.**
- **Still large errors both in data (35%) and SM (17%)**
- **Deviation $\sim 1.7\sigma$**

NP : **Central value of the data cannot be reproduced**



[3] $B_c \rightarrow \tau \nu$

Grinstein et al, 1611.06676
Akeroyd, 1708.04072

Not directly measured,
but some limits have been estimated from Bc decay

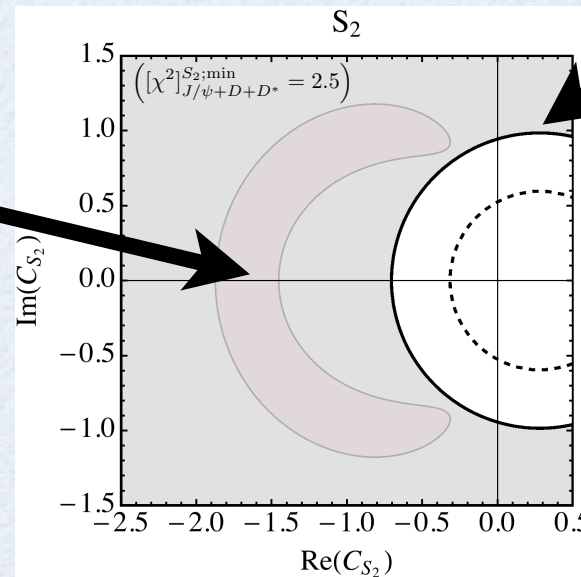
From Bc life time : uncertainty of theoretical evaluation implies

From LEP data : extracted from data at the Z boson peak

- Indirect bound is then given as $\mathcal{B}(B_c \rightarrow \tau \nu) < 10 - 30\%$
- This **kills** the Scalar NP explanation

[Best fit point]

$$C_{S_2} \sim -1.5$$



Thank you!

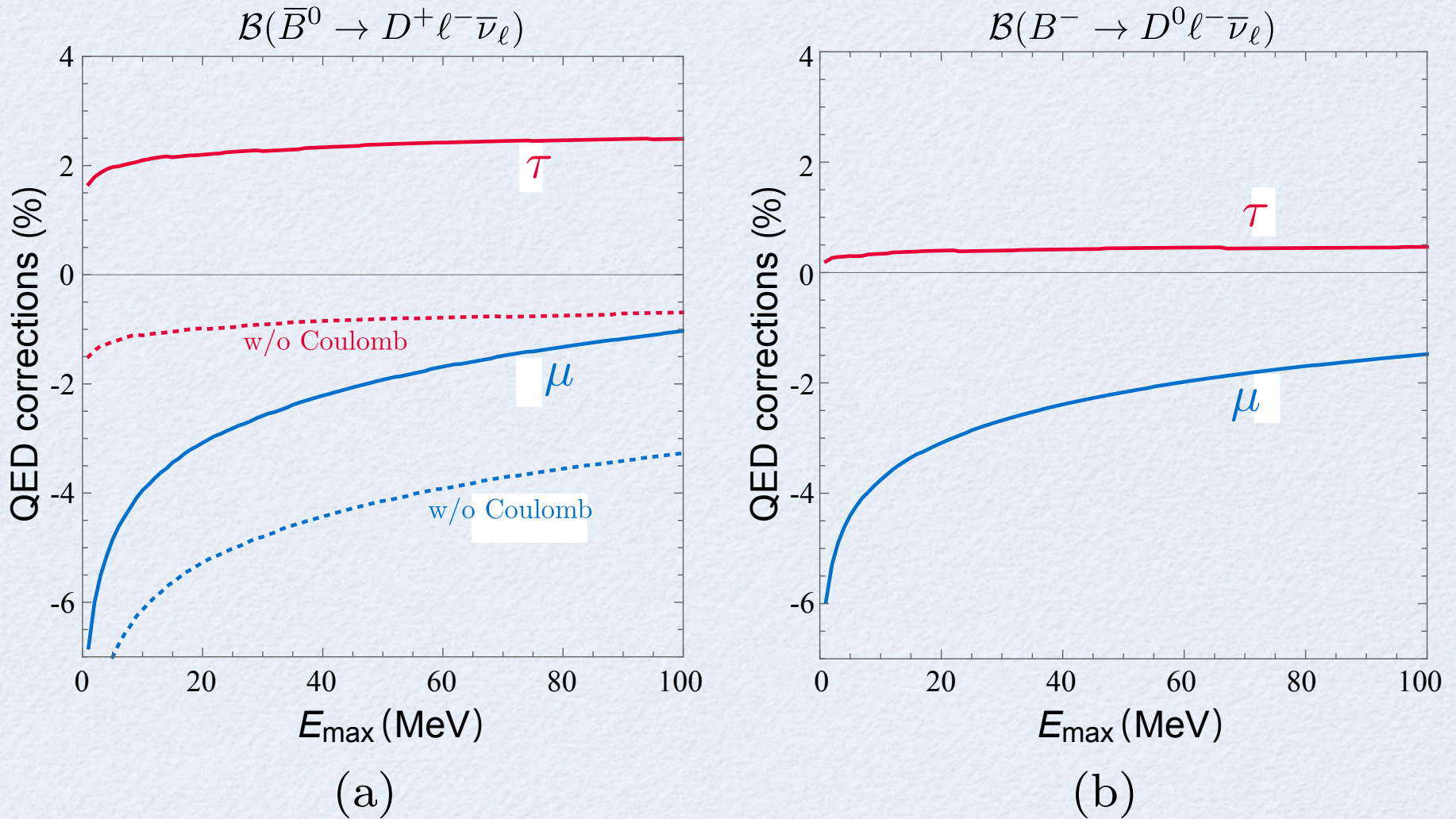


FIG. 2. (a) The long-distance QED corrections to the branching ratios of $\bar{B}^0 \rightarrow D^+ \ell^- \bar{\nu}_\ell$ and (b) $B^- \rightarrow D^0 \ell^- \bar{\nu}_\ell$, where $\ell = \mu, \tau$, as a function of E_{\max} . The dotted lines show the corrections to $\bar{B}^0 \rightarrow D^+ \ell^- \bar{\nu}_\ell$ without the Coulomb contributions, for the purpose of illustration.